

# Резонансная нанофотоника



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Российский  
научный фонд

22-12-00351

**ЛЕТНЯЯ ШКОЛА  
ФОНДА «БАЗИС»**

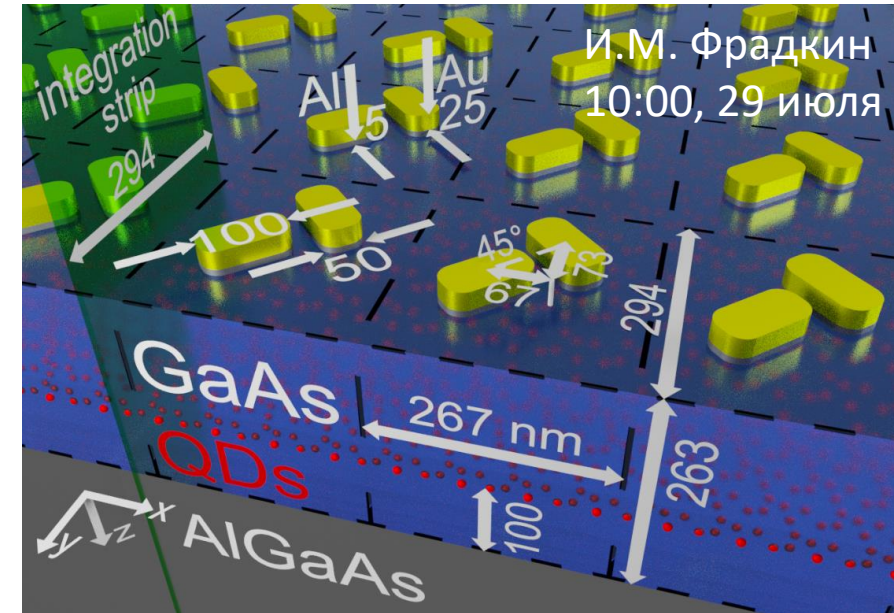
СОВРЕМЕННЫЕ ПРОБЛЕМЫ  
ФИЗИКИ КОНДЕНСИРОВАННОГО  
СОСТОЯНИЯ

21 ИЮЛЯ — 1 АВГУСТА

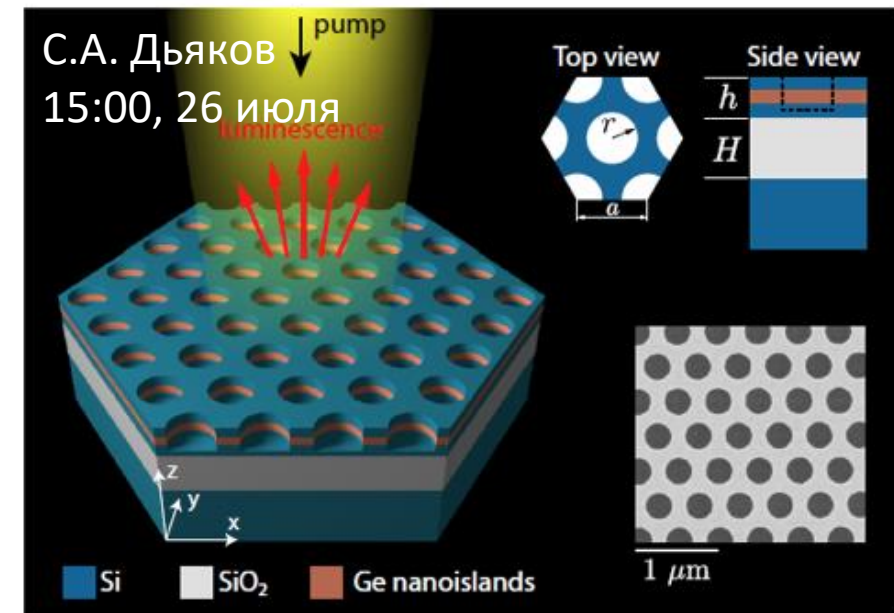
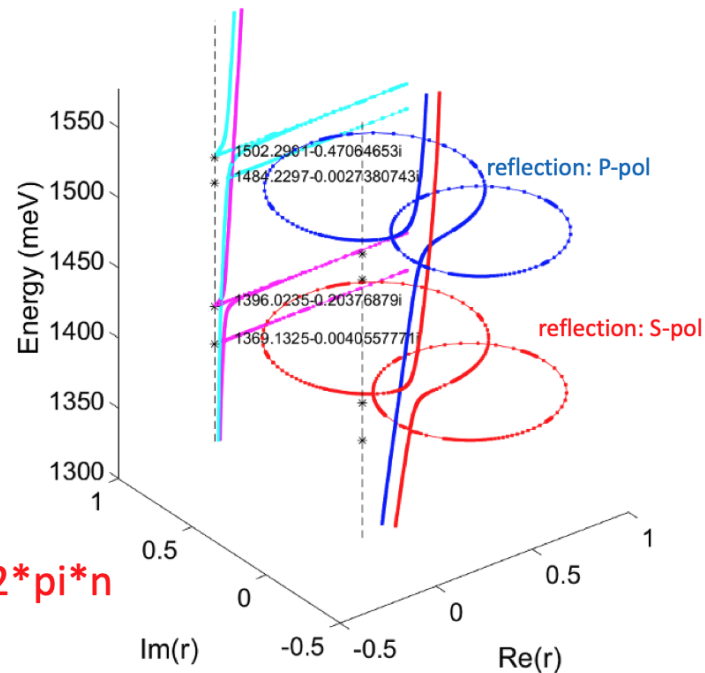
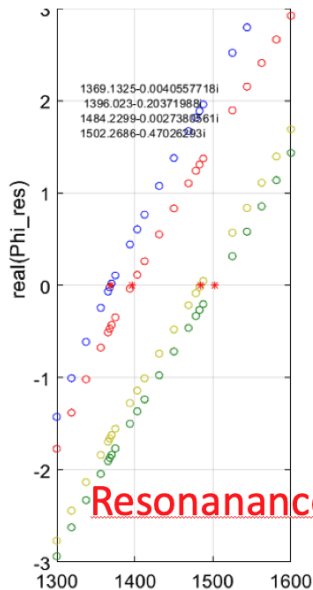
2024

# План лекций

1. Резонансная нанофотоника
2. Резонансы в фотонно-кристаллических слоях: фундаментальные основы и применения
3. Гибридные резонансы в решетках плазмонных наночастиц

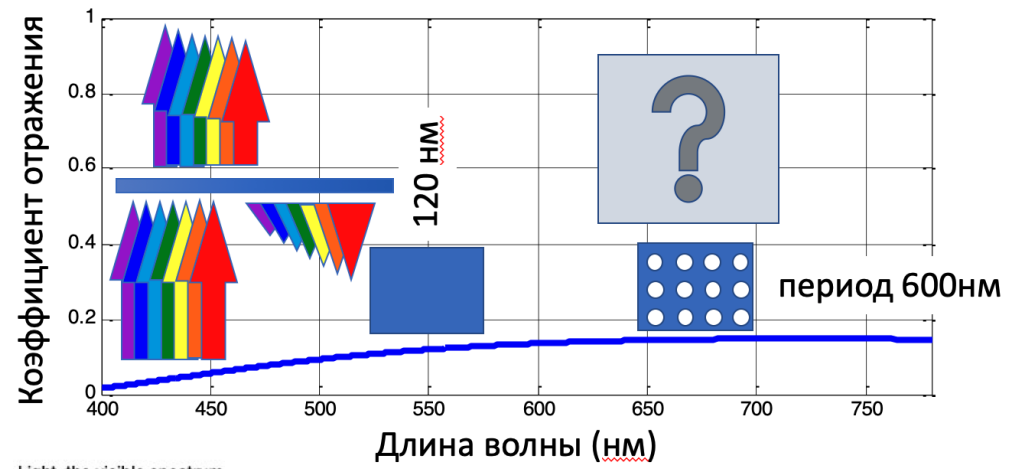


Н.А. Гиппиус, 10:00, 23 июля

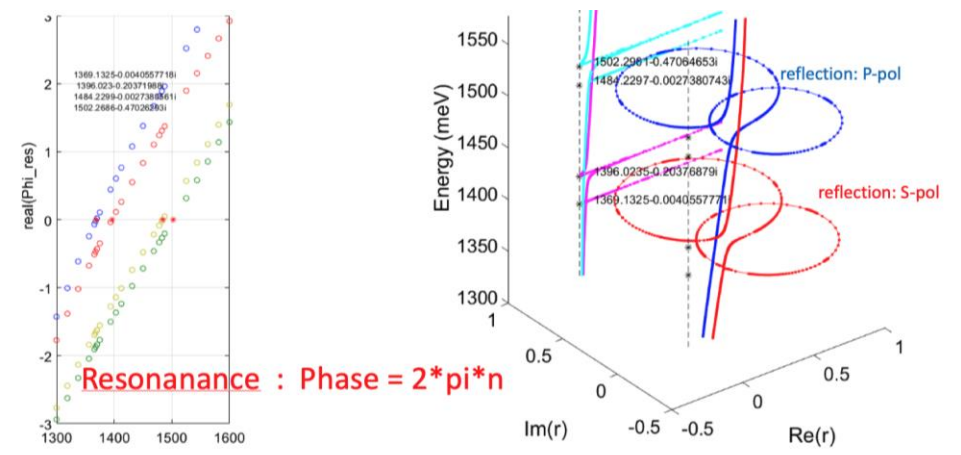


# План этой лекции

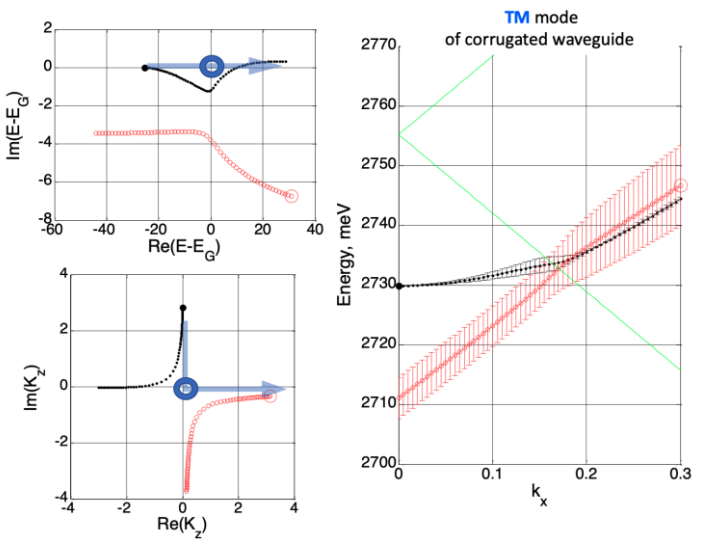
## 1. Как поймать свет в решетке?



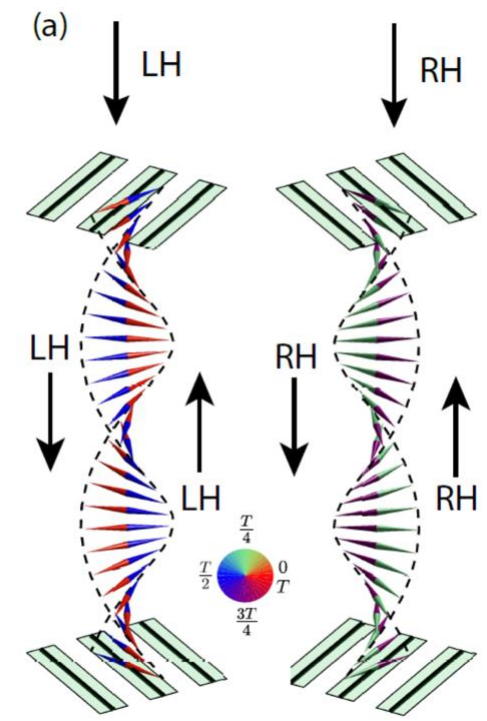
## 2. Сколько можно считать одно и тоже?!

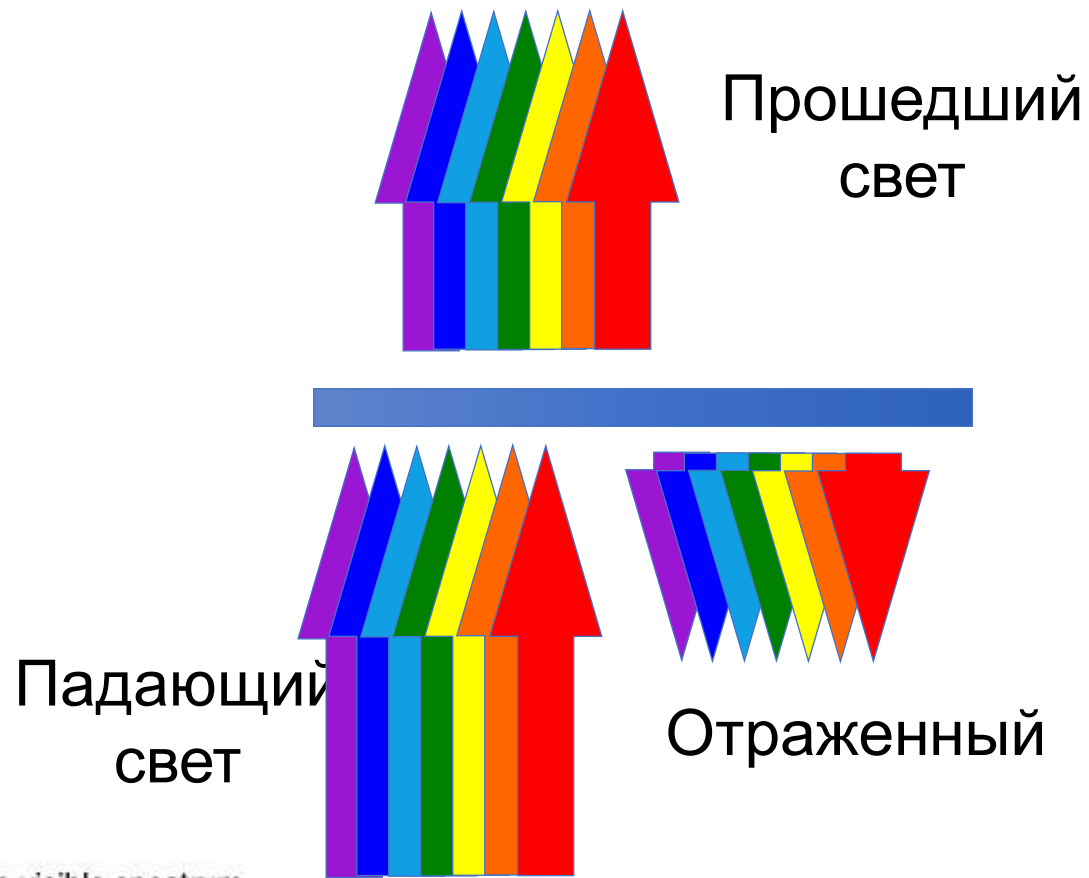


## 3. Как раскалываются резонансы?

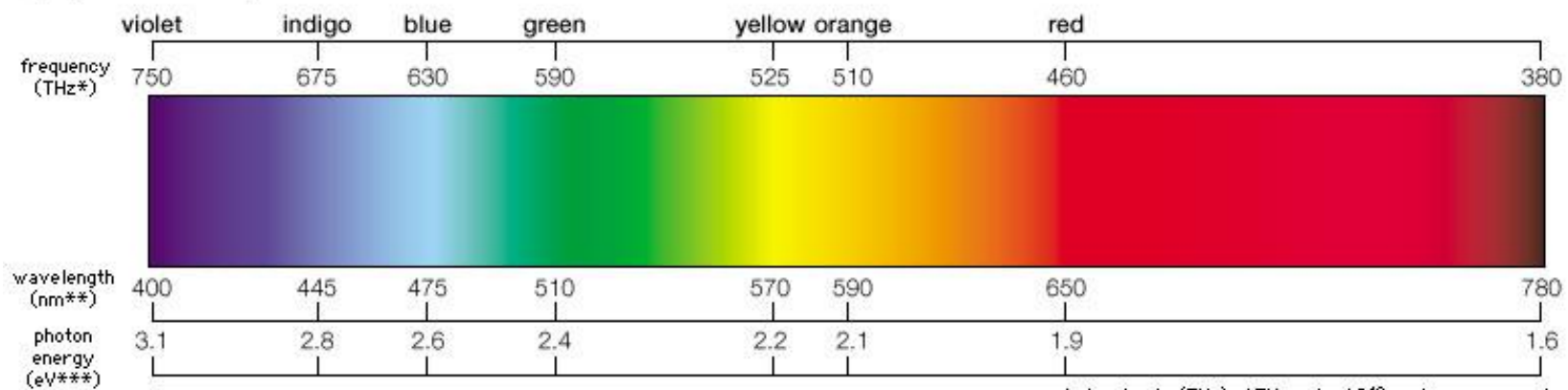


## 4. Как скрутить свет винтом?



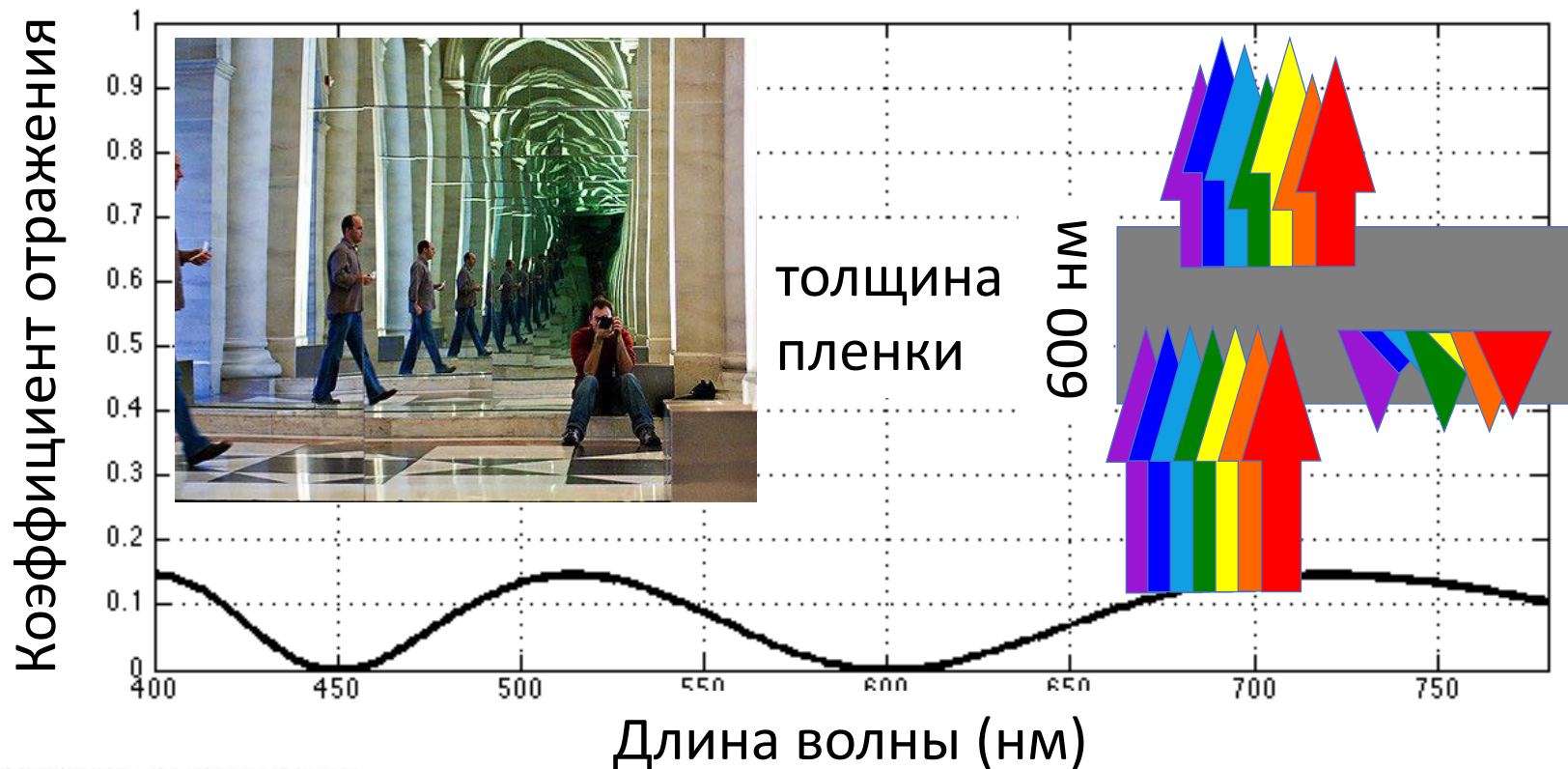


Light, the visible spectrum

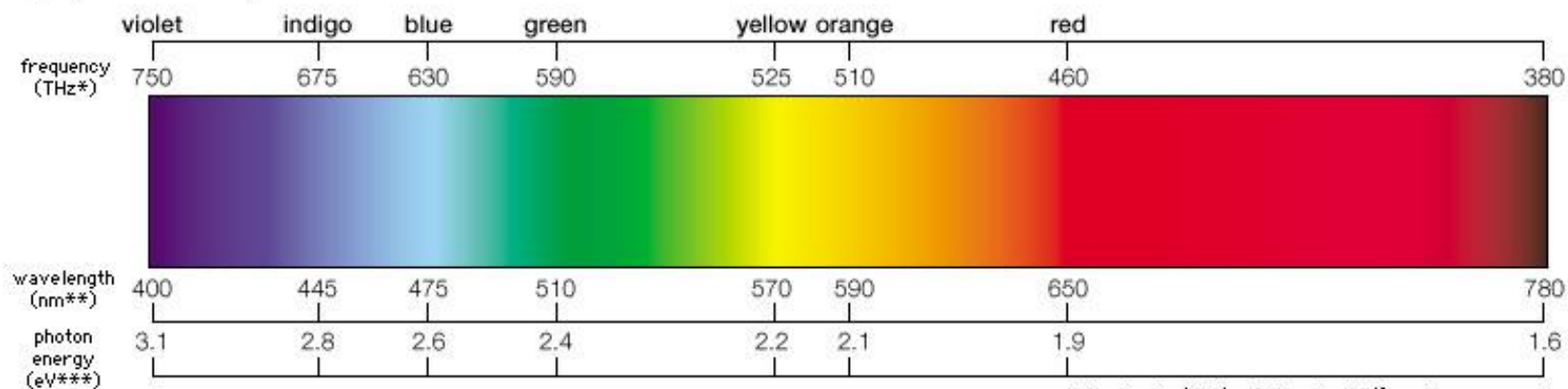


\* In terahertz (THz); 1 THz =  $1 \times 10^{12}$  cycles per second.  
 \*\* In nanometres (nm); 1 nm =  $1 \times 10^{-9}$  metre.  
 \*\*\* In electron volts (eV).

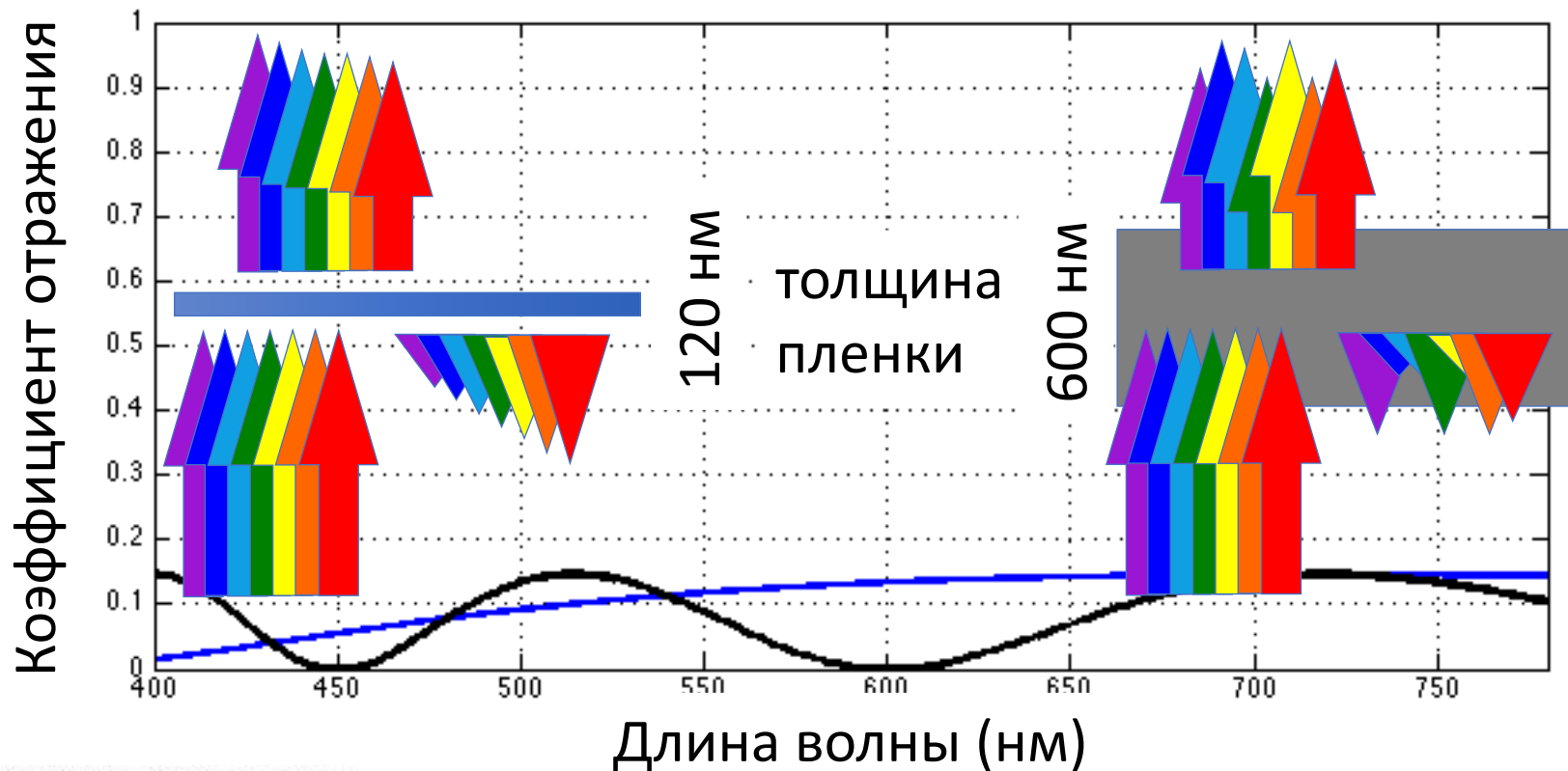




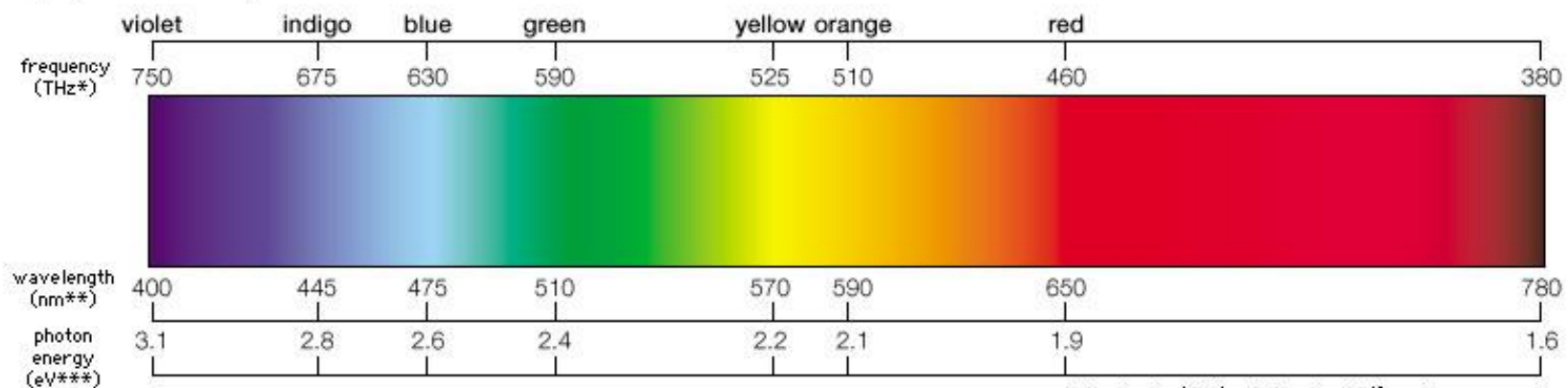
Light, the visible spectrum



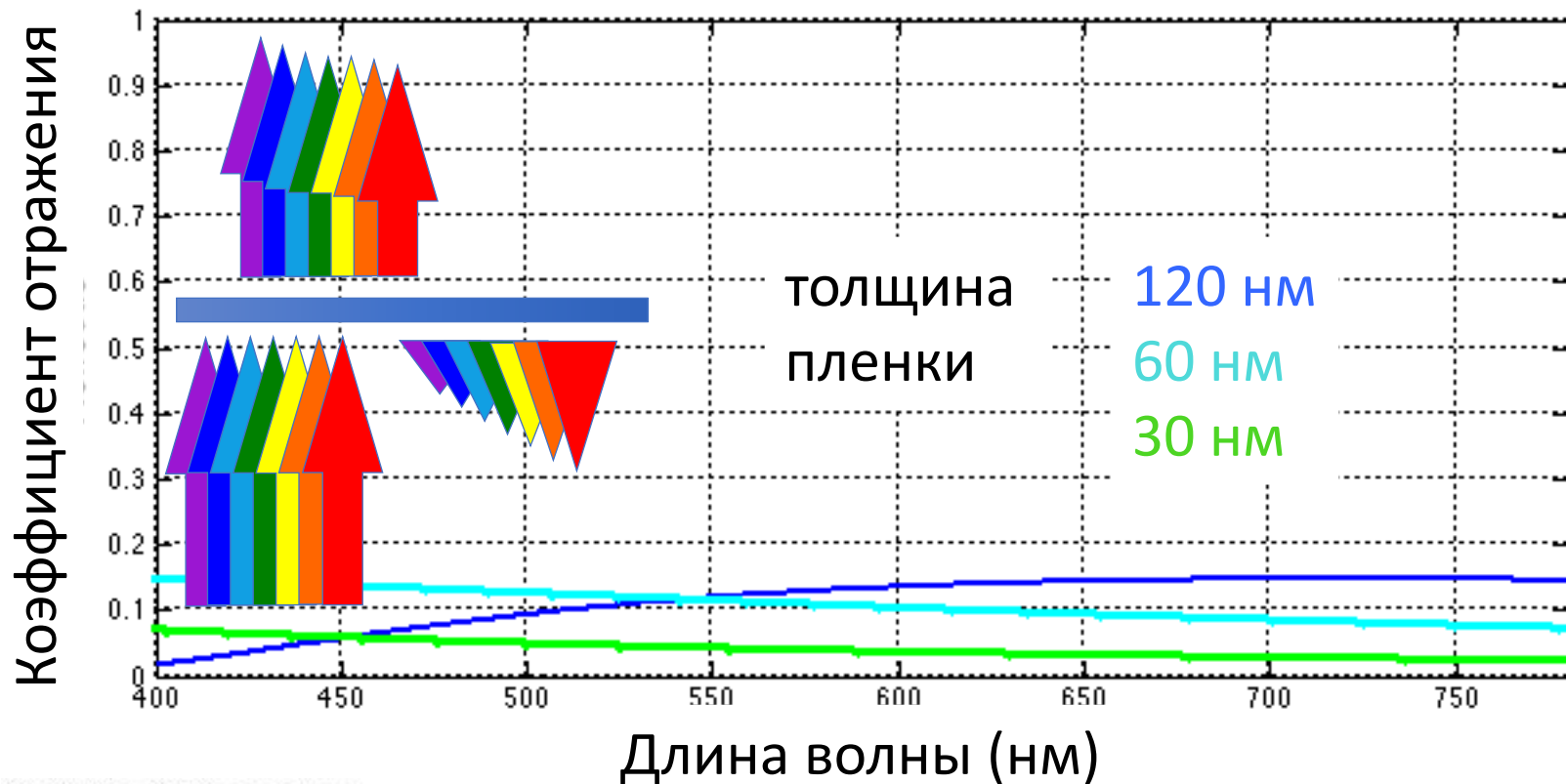
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 \*\* In nanometres (nm); 1 nm =  $1 \times 10^{-9}$  metre.  
 \*\*\* In electron volts (eV).



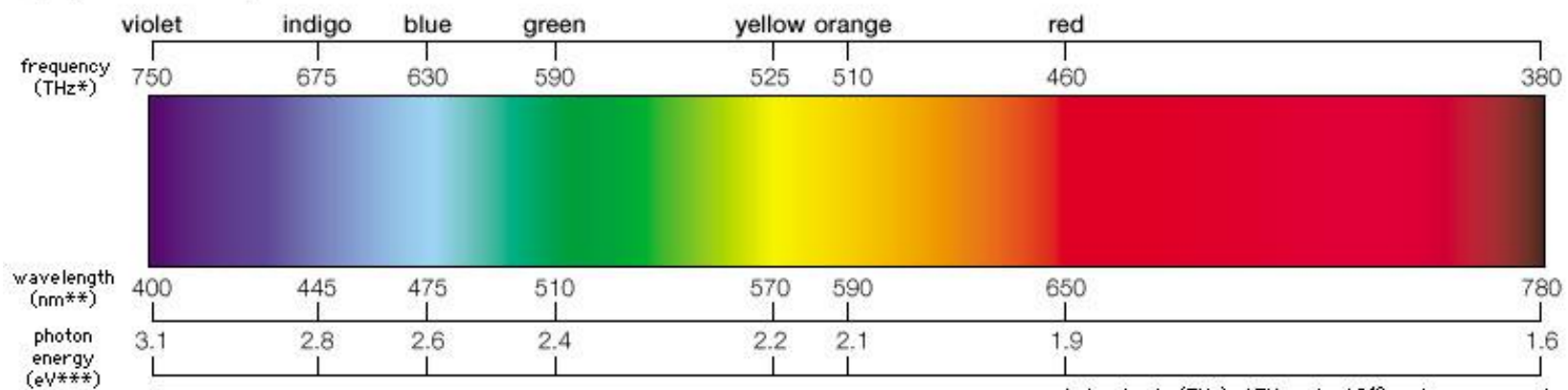
Light, the visible spectrum



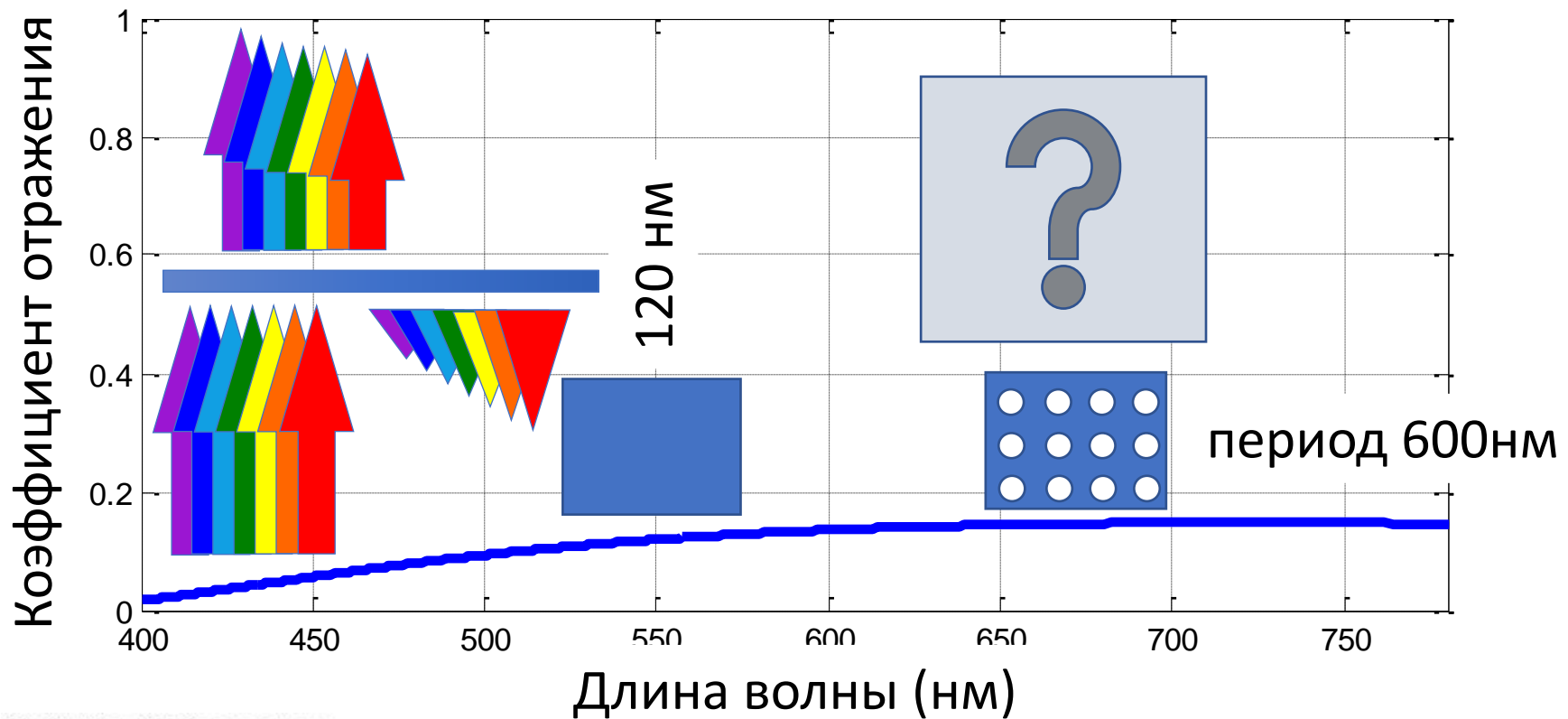
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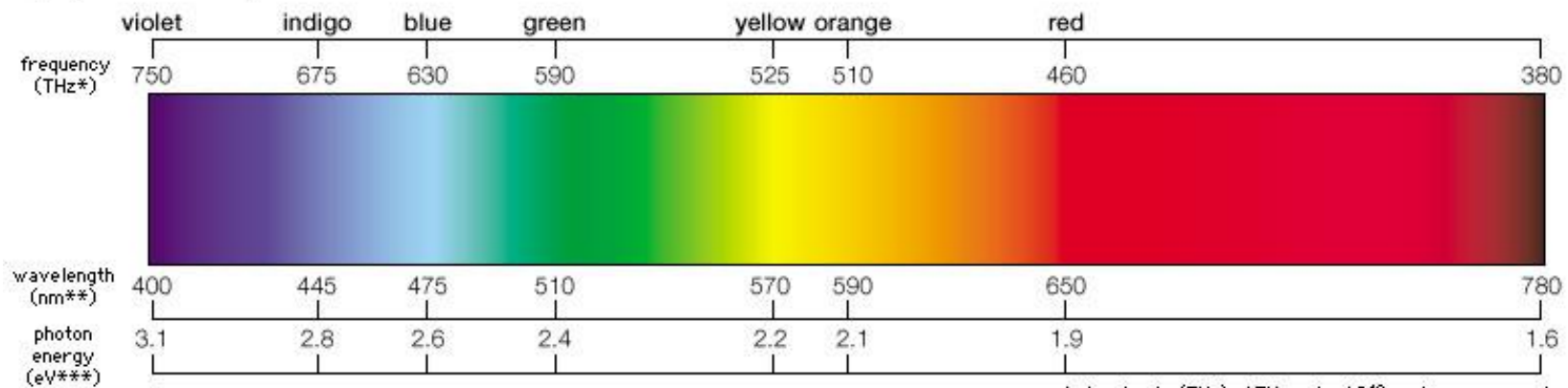
Light, the visible spectrum



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 \*\*\* In electron volts (eV).

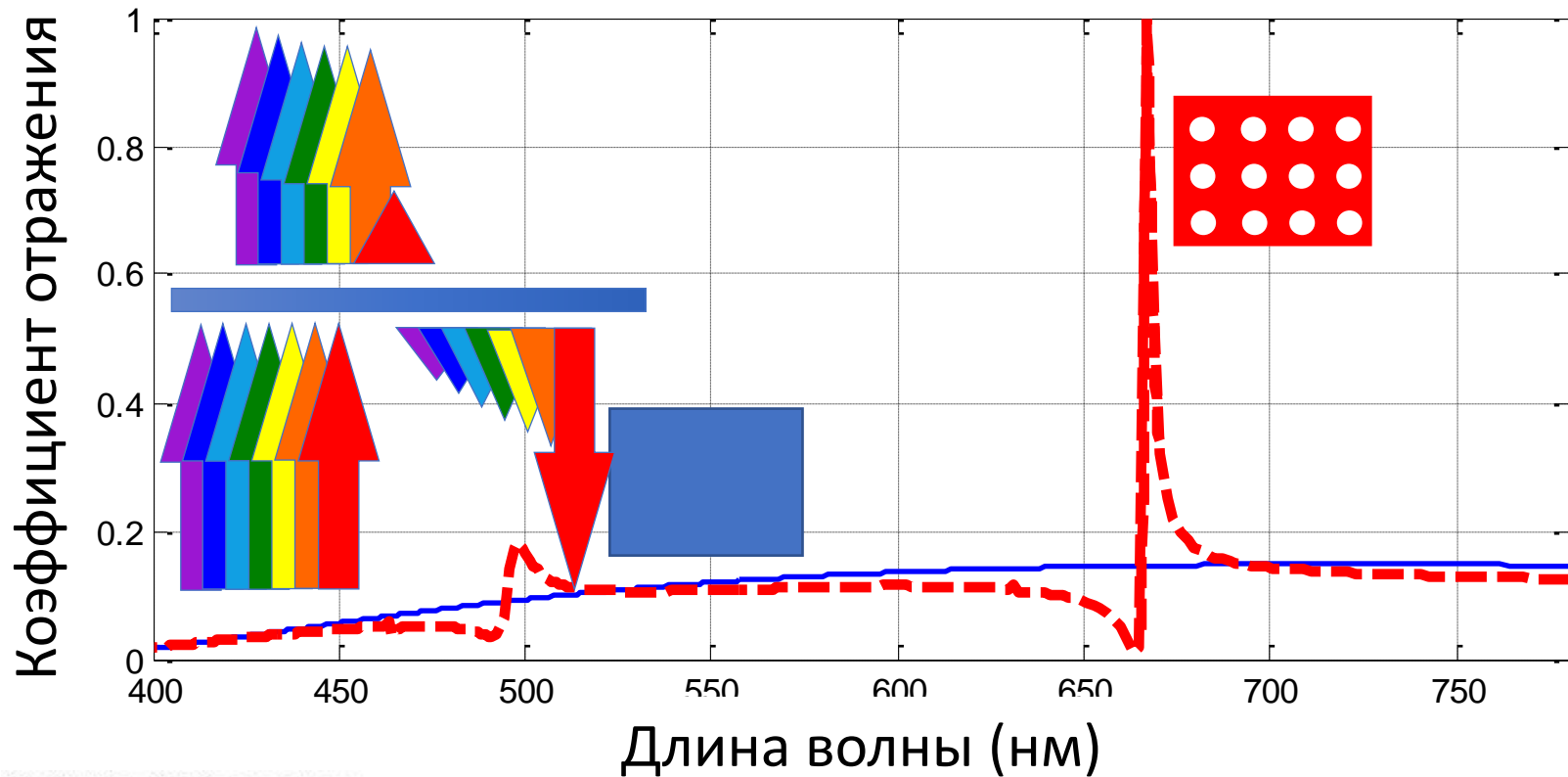


Light, the visible spectrum

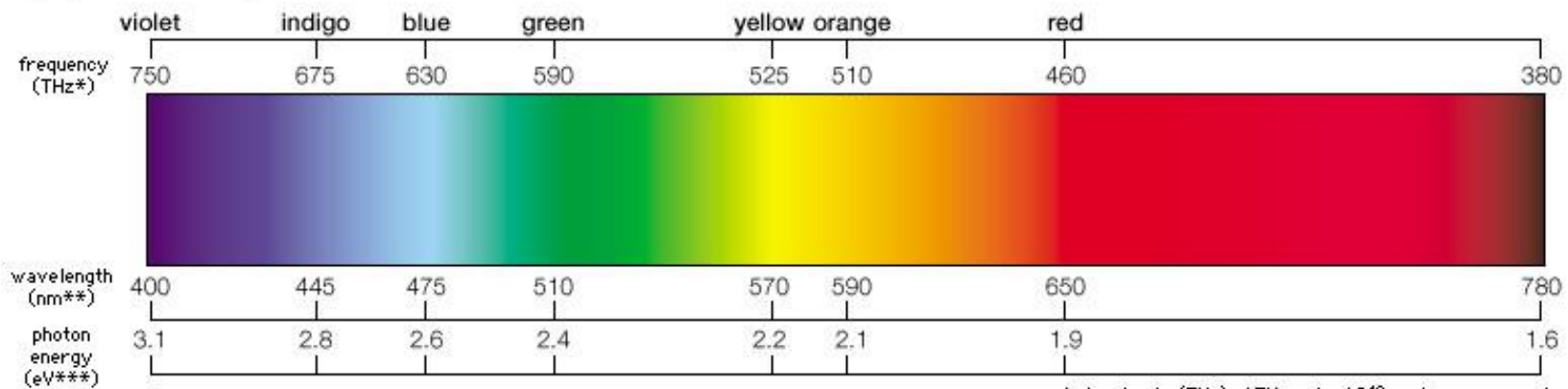


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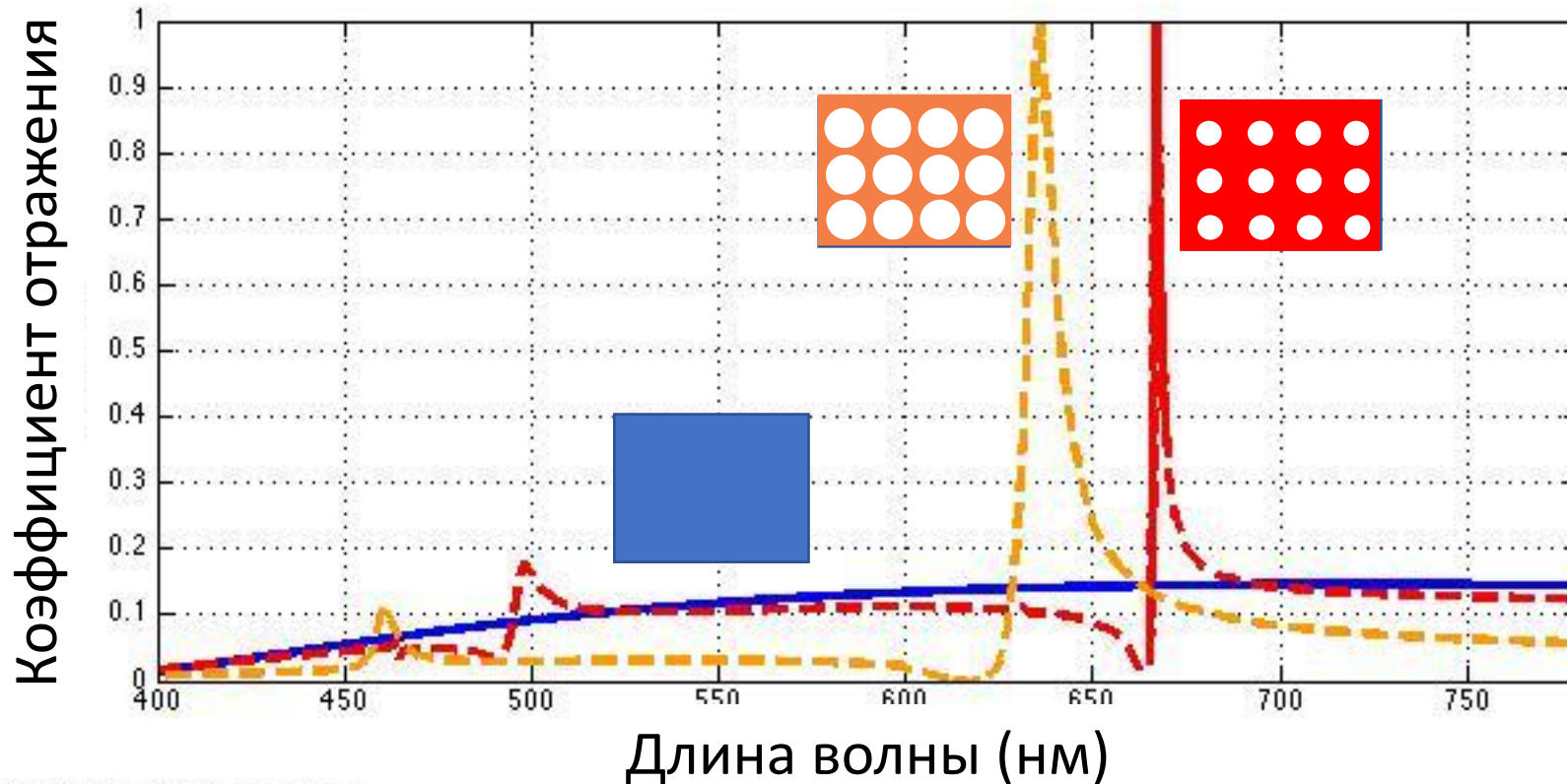




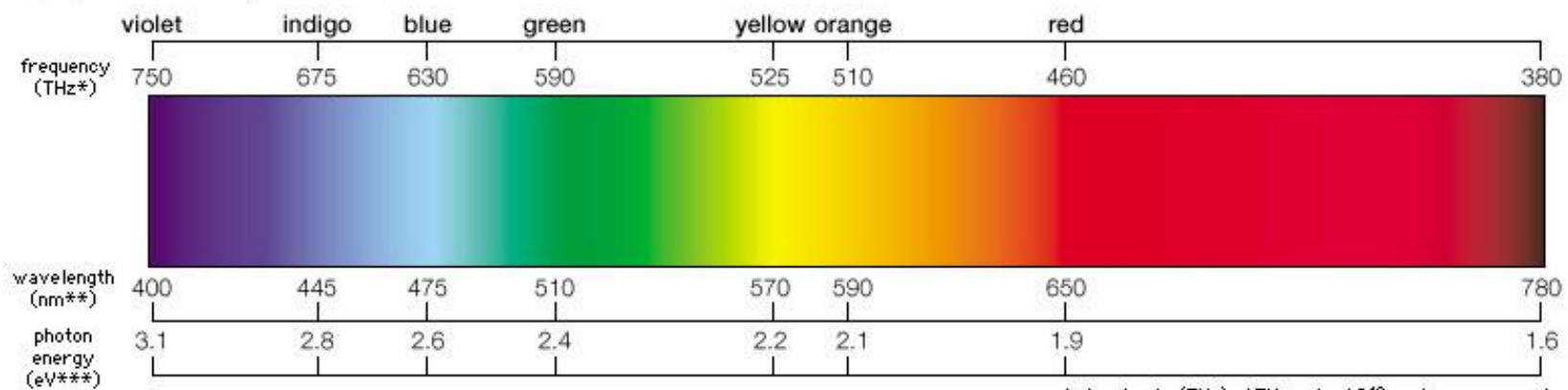
Light, the visible spectrum



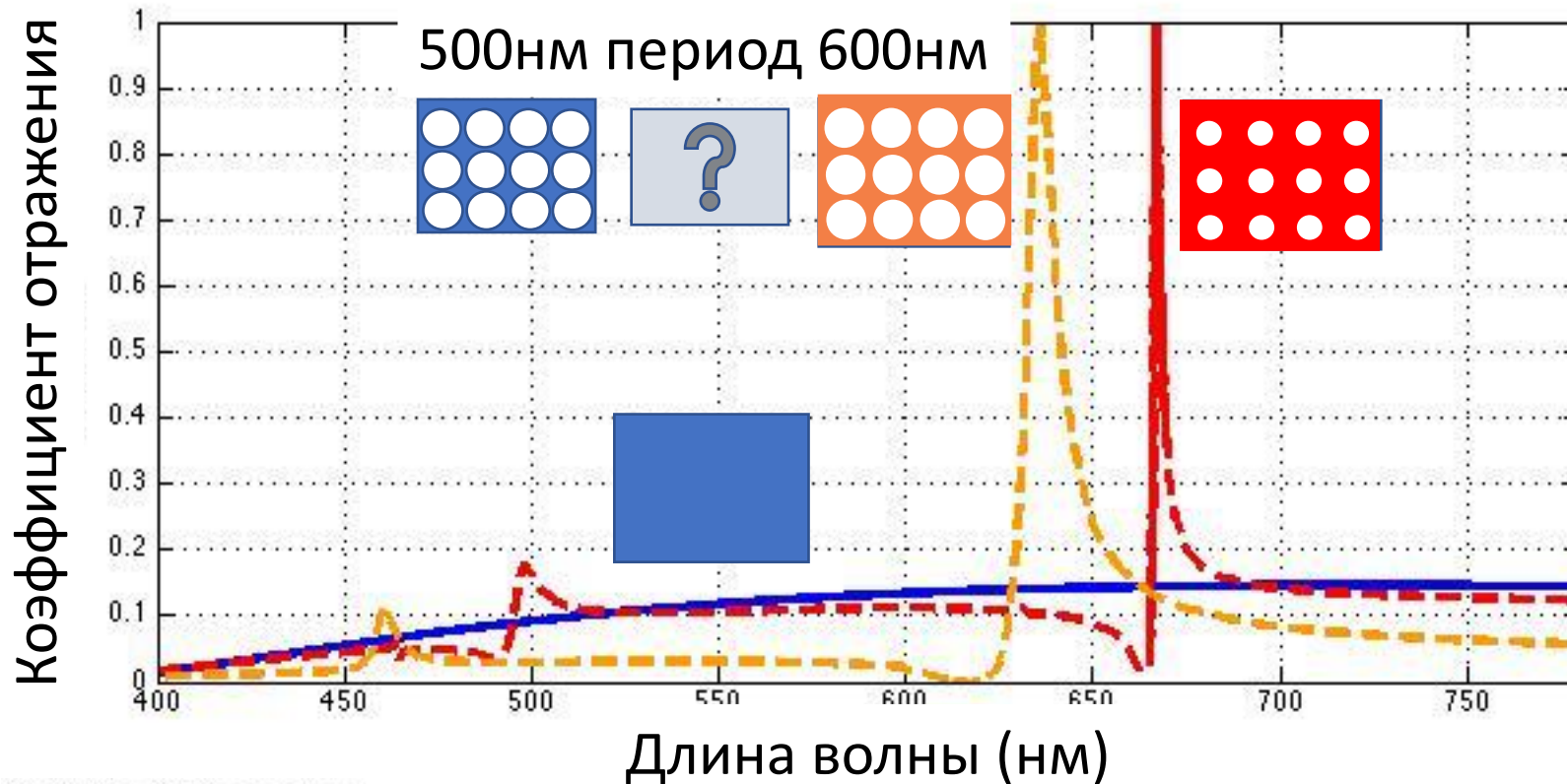
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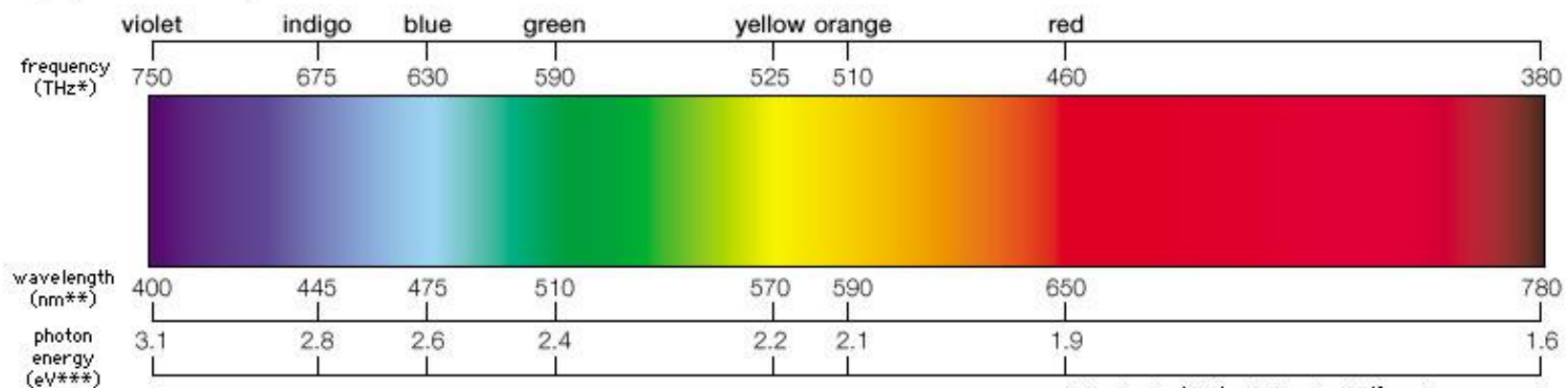
Light, the visible spectrum



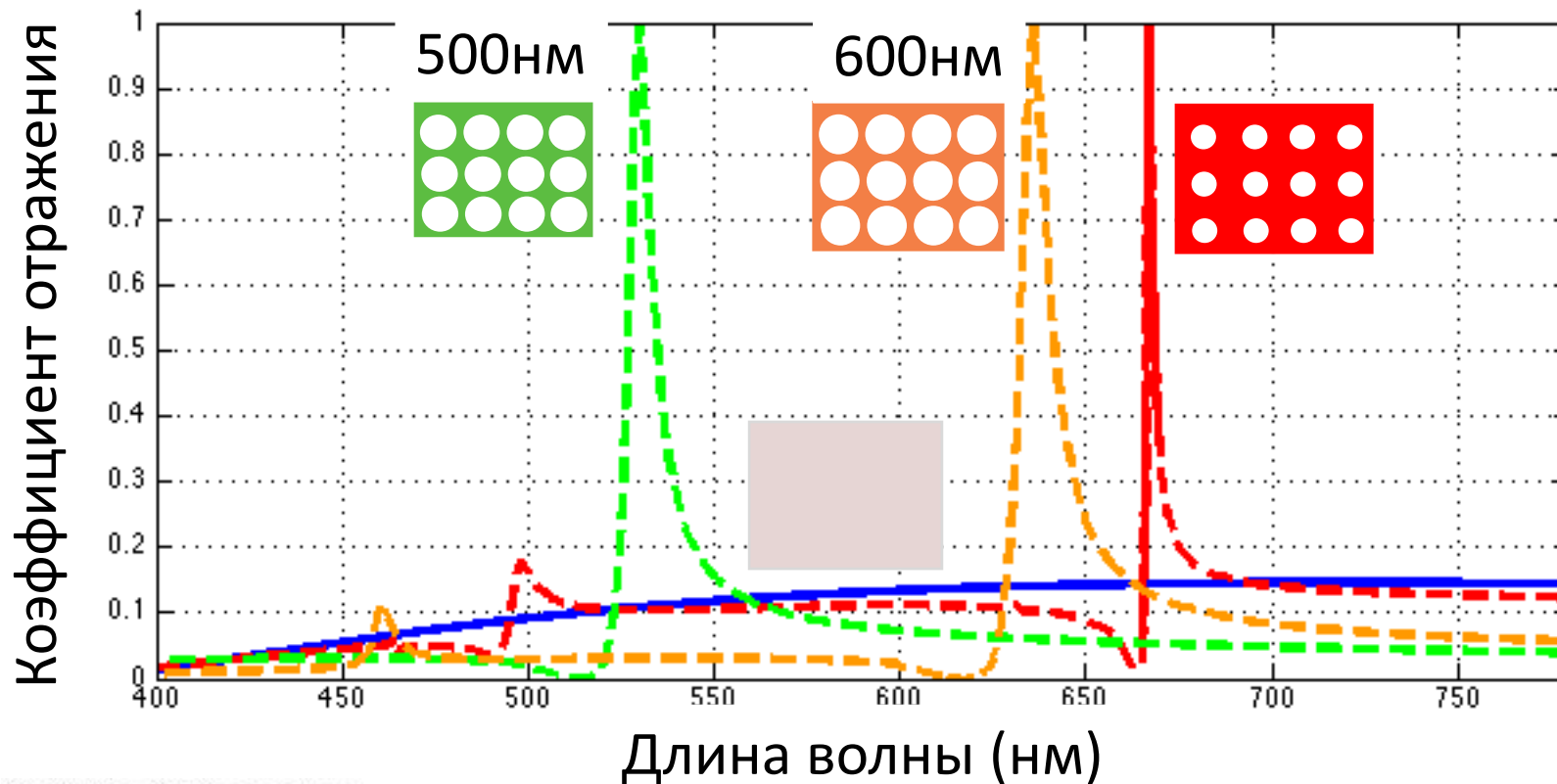
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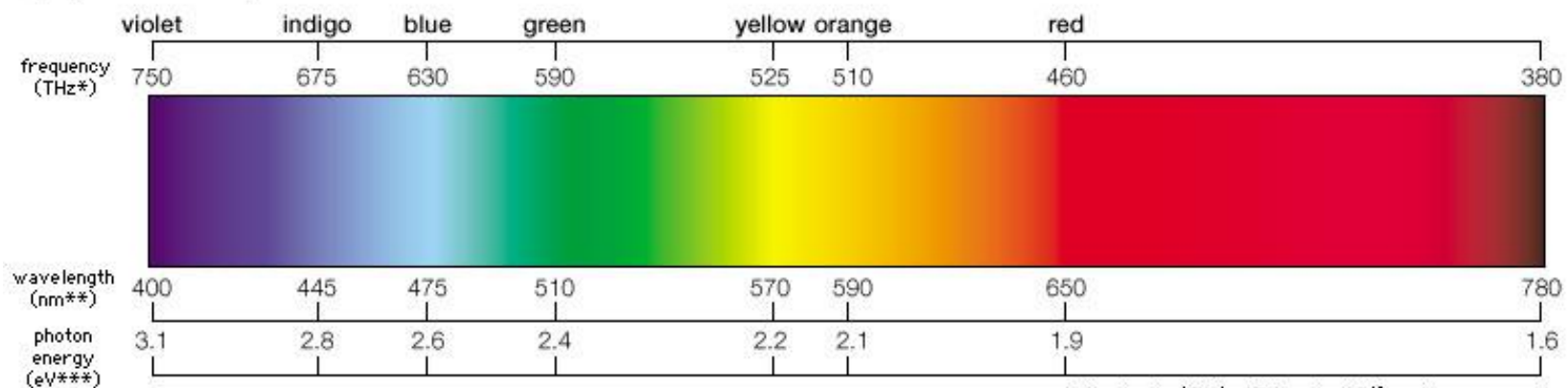
Light, the visible spectrum



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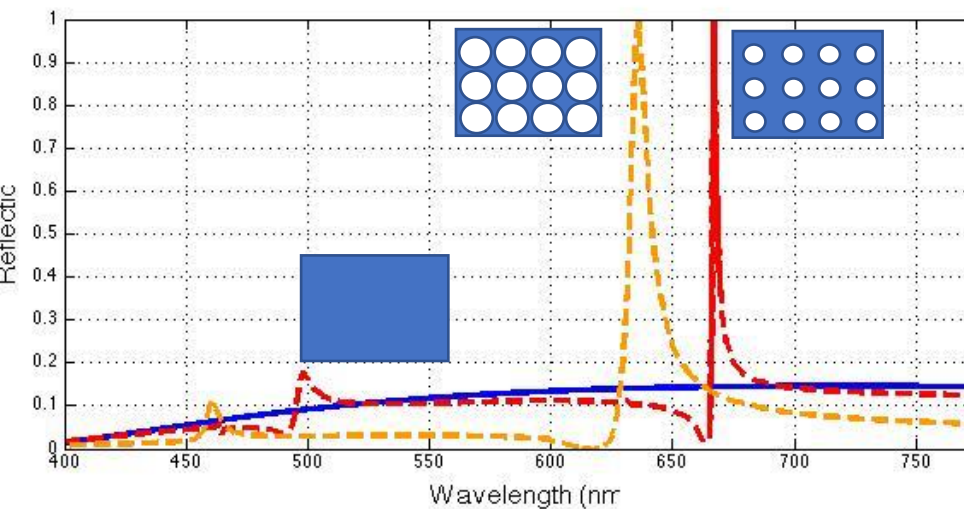


Light, the visible spectrum

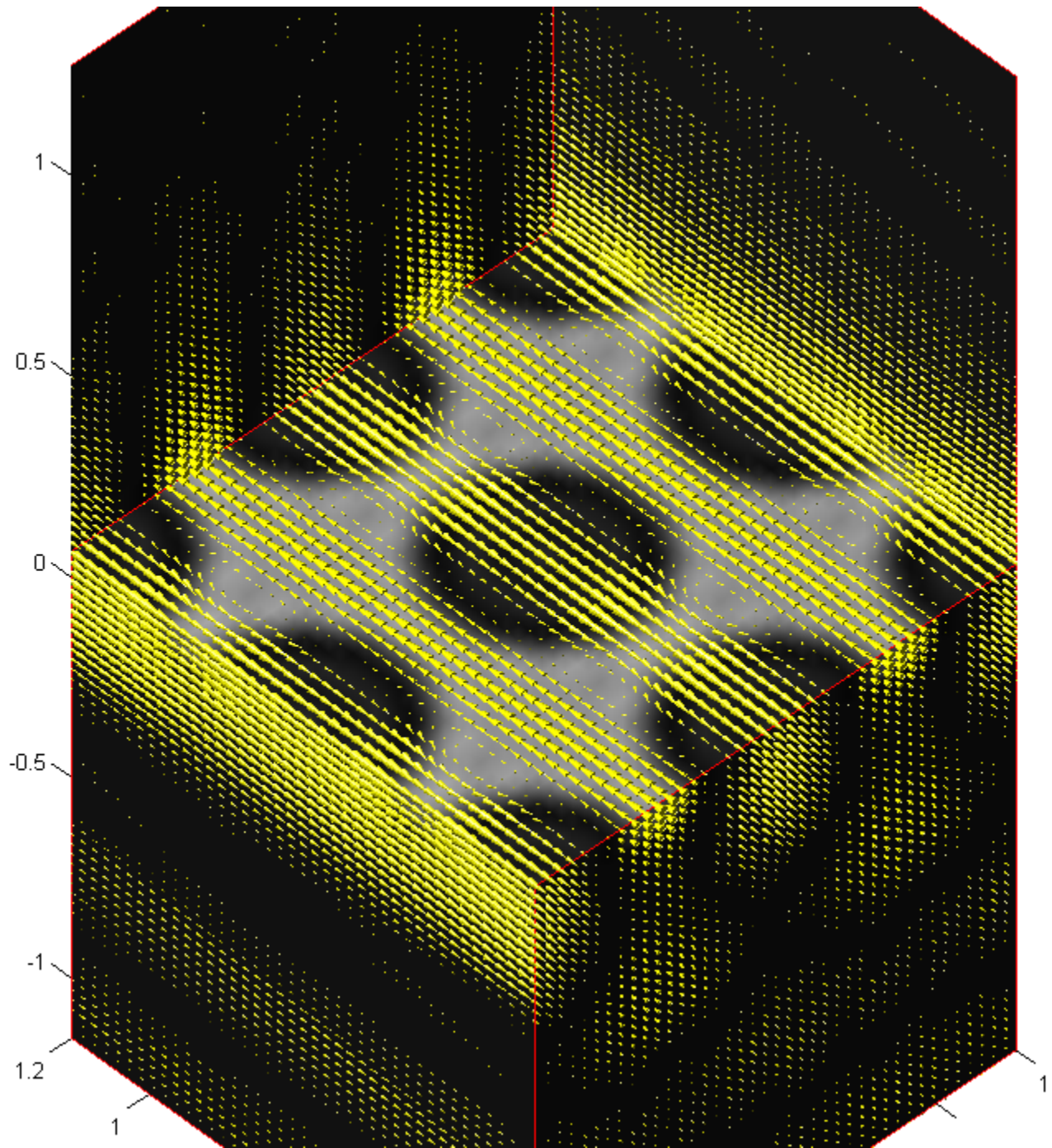


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 \*\* In nanometres (nm); 1 nm =  $1 \times 10^{-9}$  metre.  
 \*\*\* In electron volts (eV).

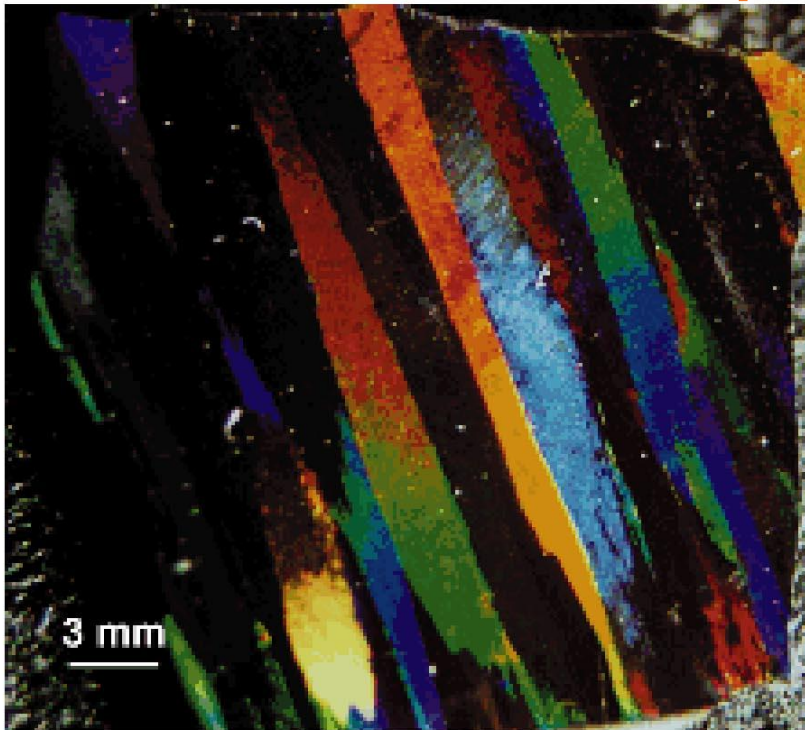




Свет иногда  
усиливается  
вблизи пленки с  
отверстиями

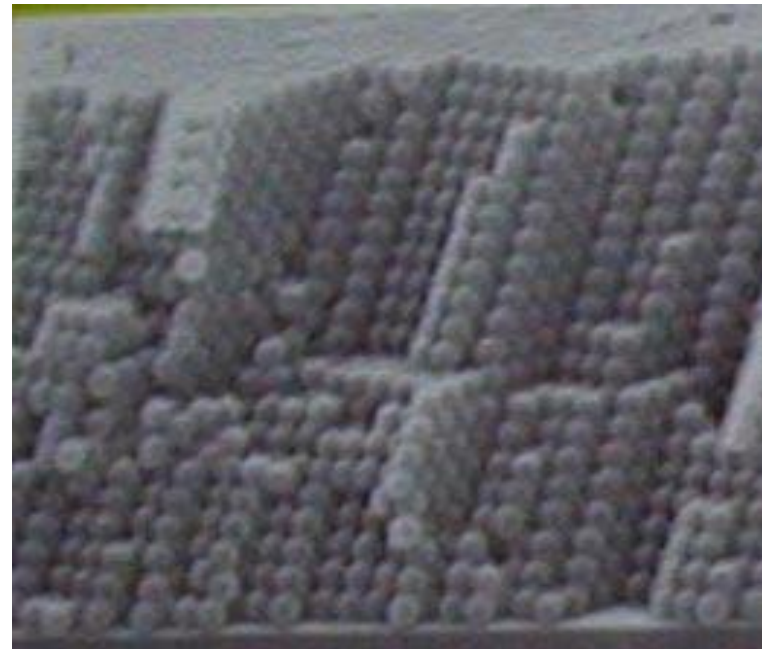
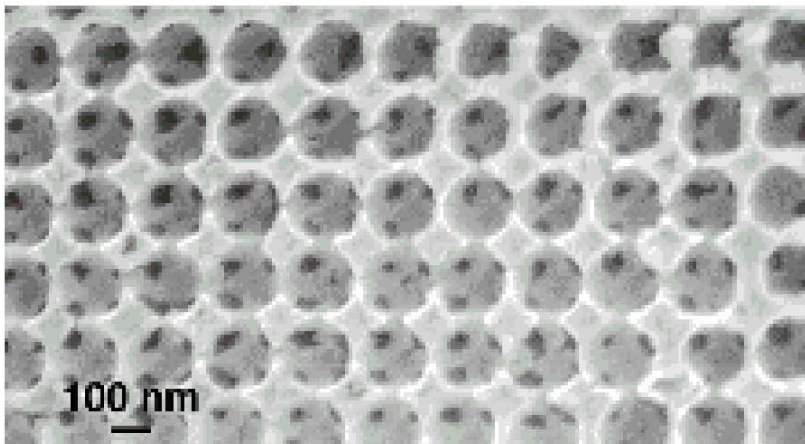


## Фотонные кристаллы в природе



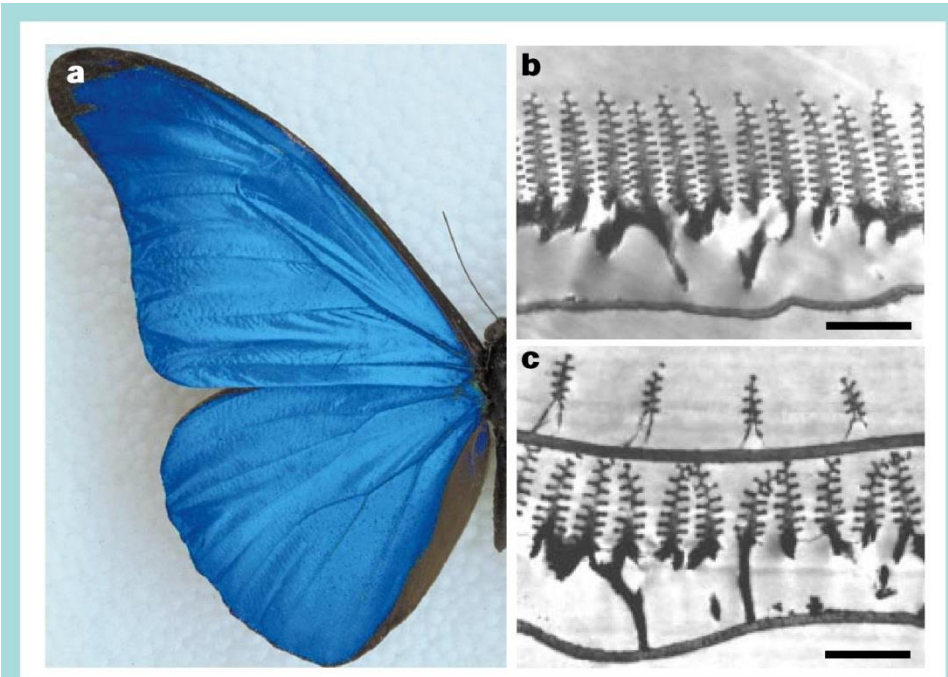
**Опалы:** естественные фотонные кристаллы, образованные слипшимися нано-шариками

На рисунке показаны искусственные опало-подобные структуры из фуллеренов (слева) и кварцевых наночастиц (внизу)

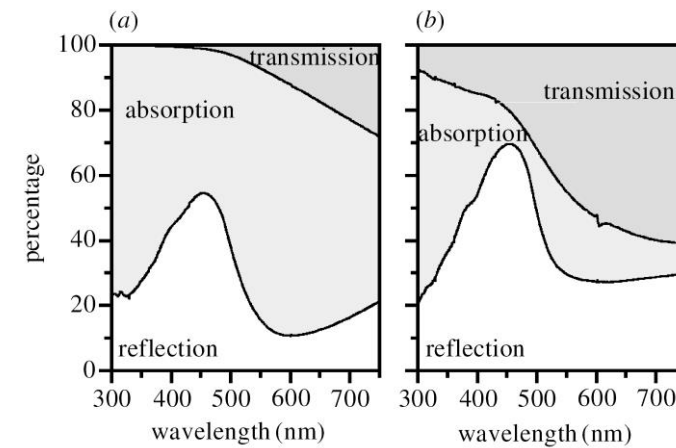




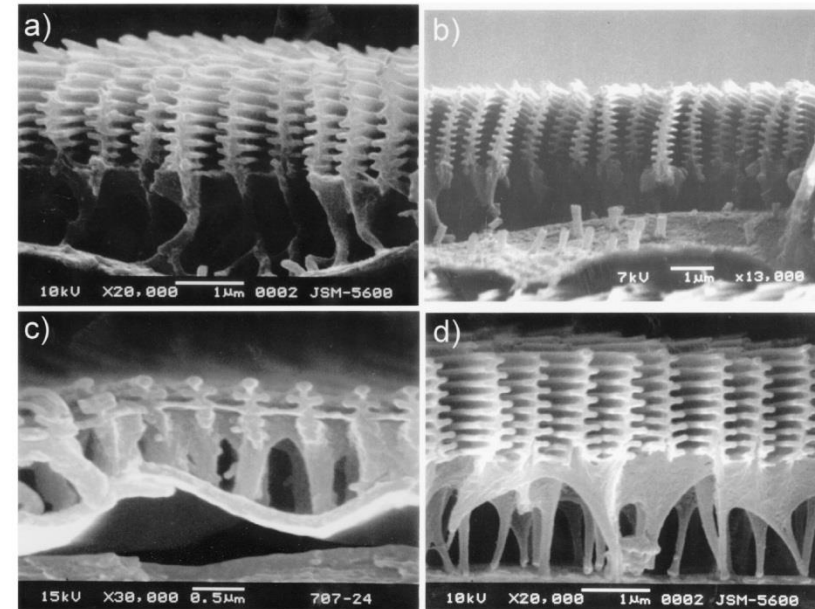
# Окраска бабочки



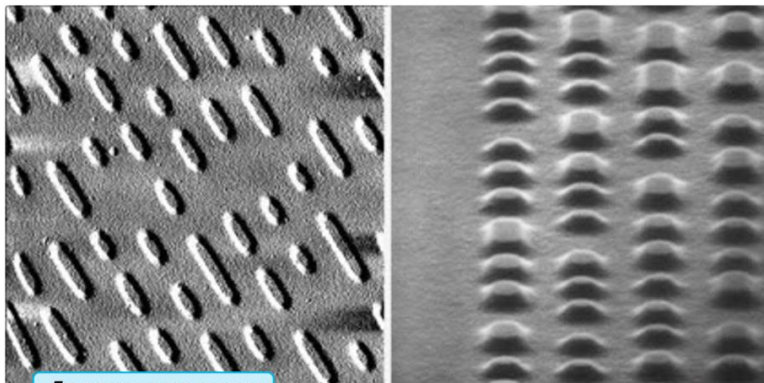
**Figure 3** Iridescence in the butterfly *Morpho rhetenor*. **a**, Real colour image of the blue iridescence from a *M. rhetenor* wing. **b**, Transmission electron micrograph (TEM) images showing wing-scale cross-sections of *M. rhetenor*. **c**, TEM images of a wing-scale cross-section of the related species *M. didius* reveal its discretely configured multilayers. The high occupancy and high layer number of *M. rhetenor* in **b** creates an intense reflectivity that contrasts with the more diffusely coloured appearance of *M. didius*, in which an overlying second layer of scales effects strong diffraction<sup>4</sup>. Bars, **a**, 1 cm; **b**, 1.8  $\mu\text{m}$ ; **c**, 1.3  $\mu\text{m}$ .



**Figure 2.** Percentages of transmission, absorption and reflection for the wings of (a) *Morpho didius* and (b) *Morpho sulkowskyi* measured by a spectrophotometer equipped with an integrated sphere.



**Fig. 3.** Scanning electron microscope images of the cross sections of the iridescent scales of *Morpho* butterflies: (a) a ground scale of *M. didius*, (b) a scale of *M. rhetenor*, (c) a cover scale of *M. adonis* and (d) a scale of *M. sulkowskyi*.

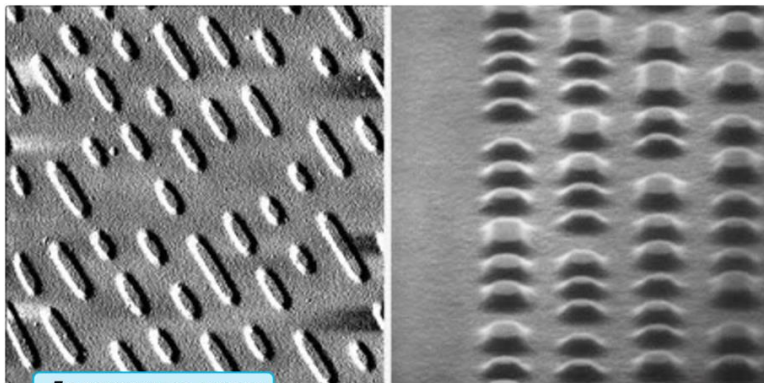


Дорожки на диске

	DVD	CD
размер штрихов (микрон)	0.4	0.83
ширина дорожки	0.74	1.6

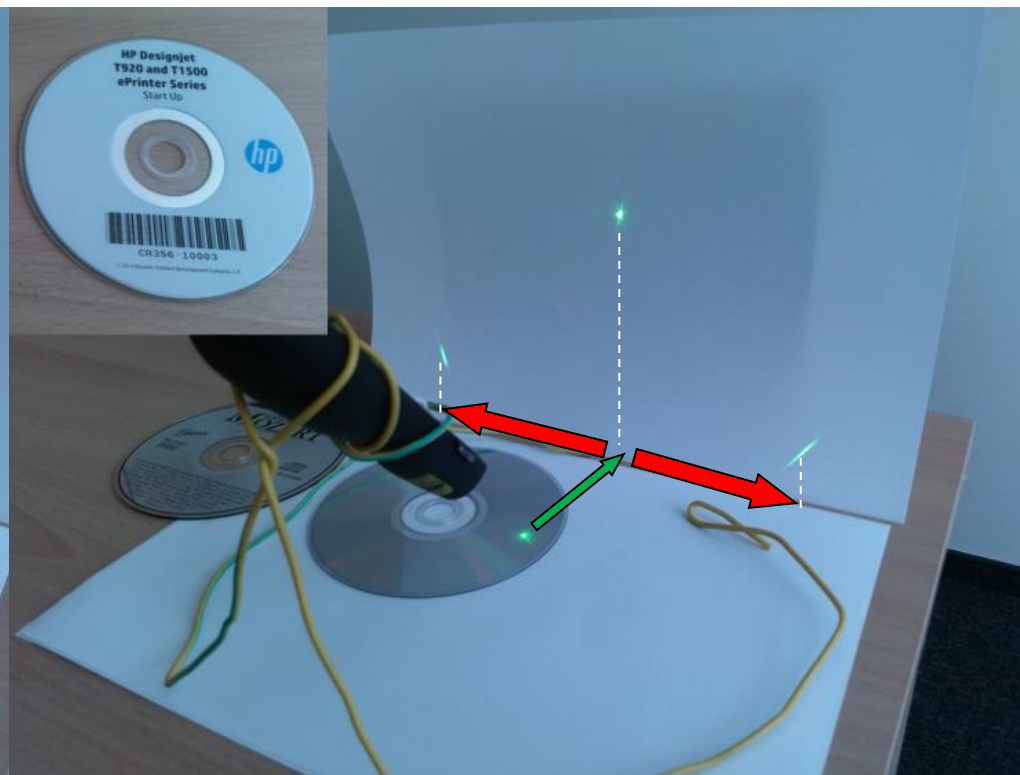
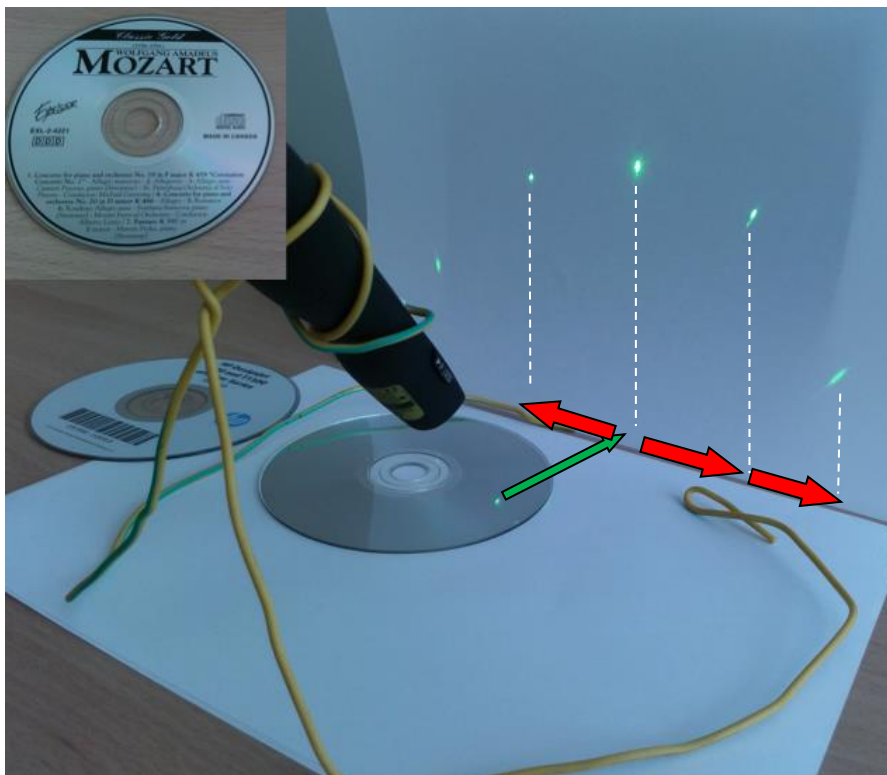


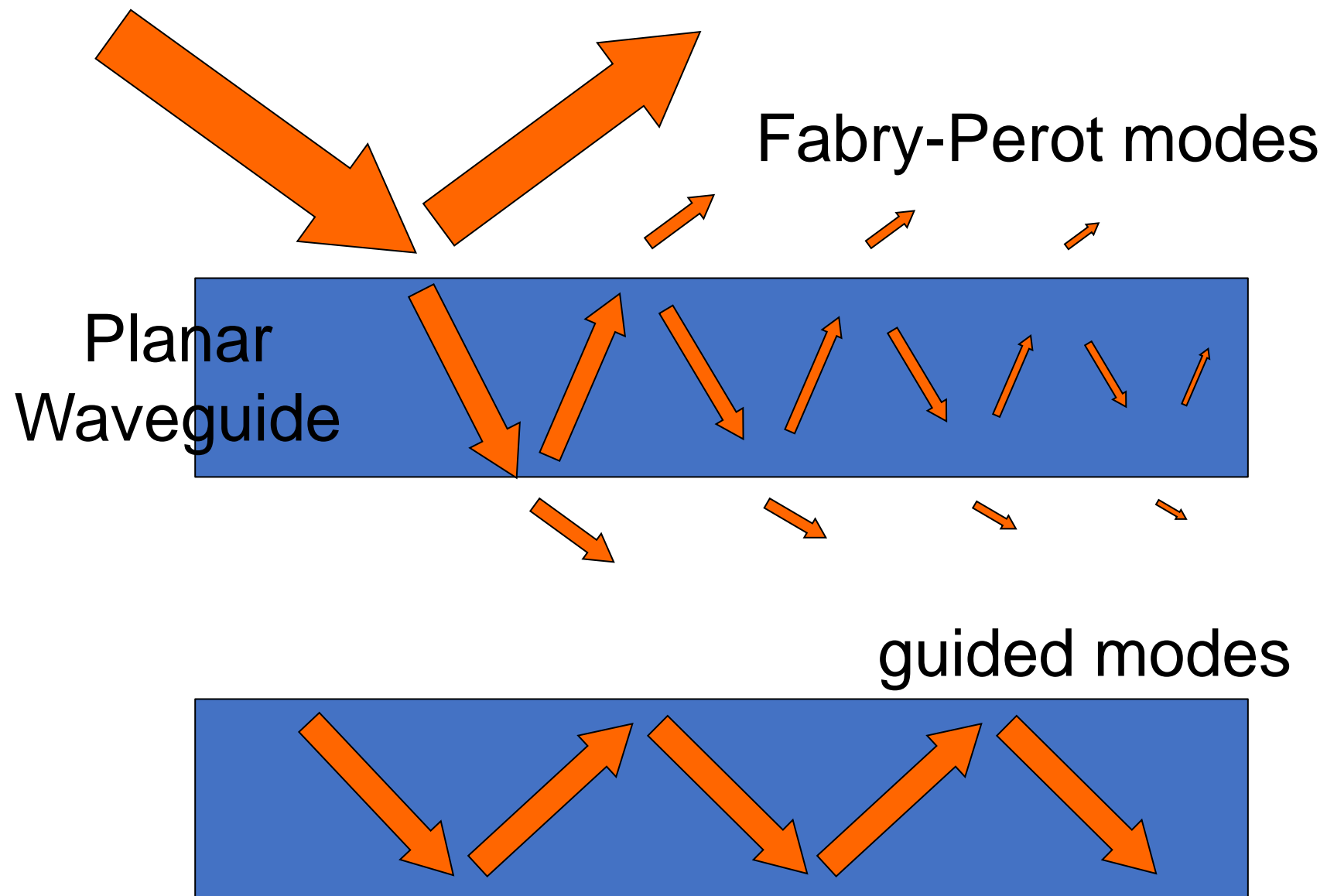




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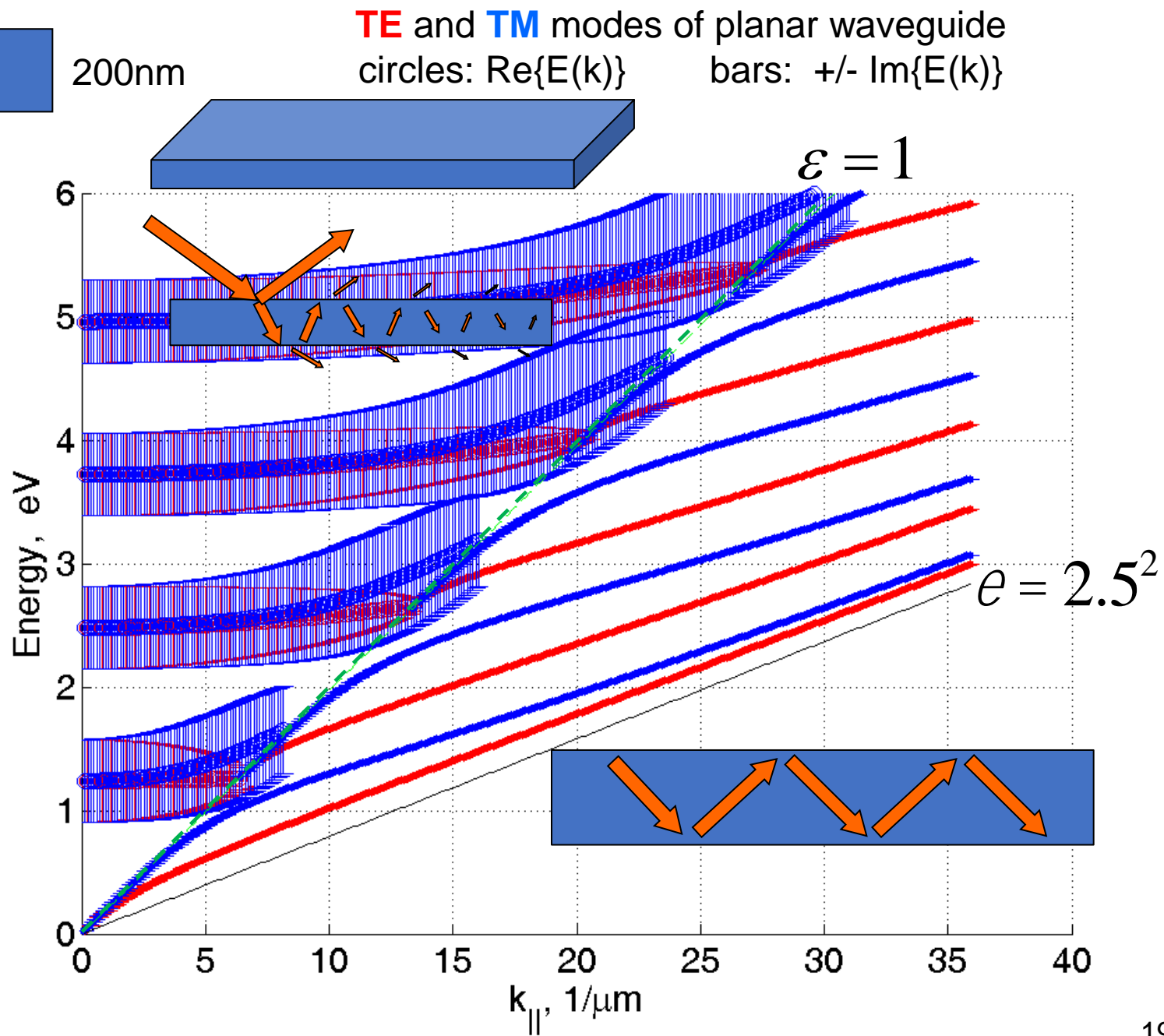


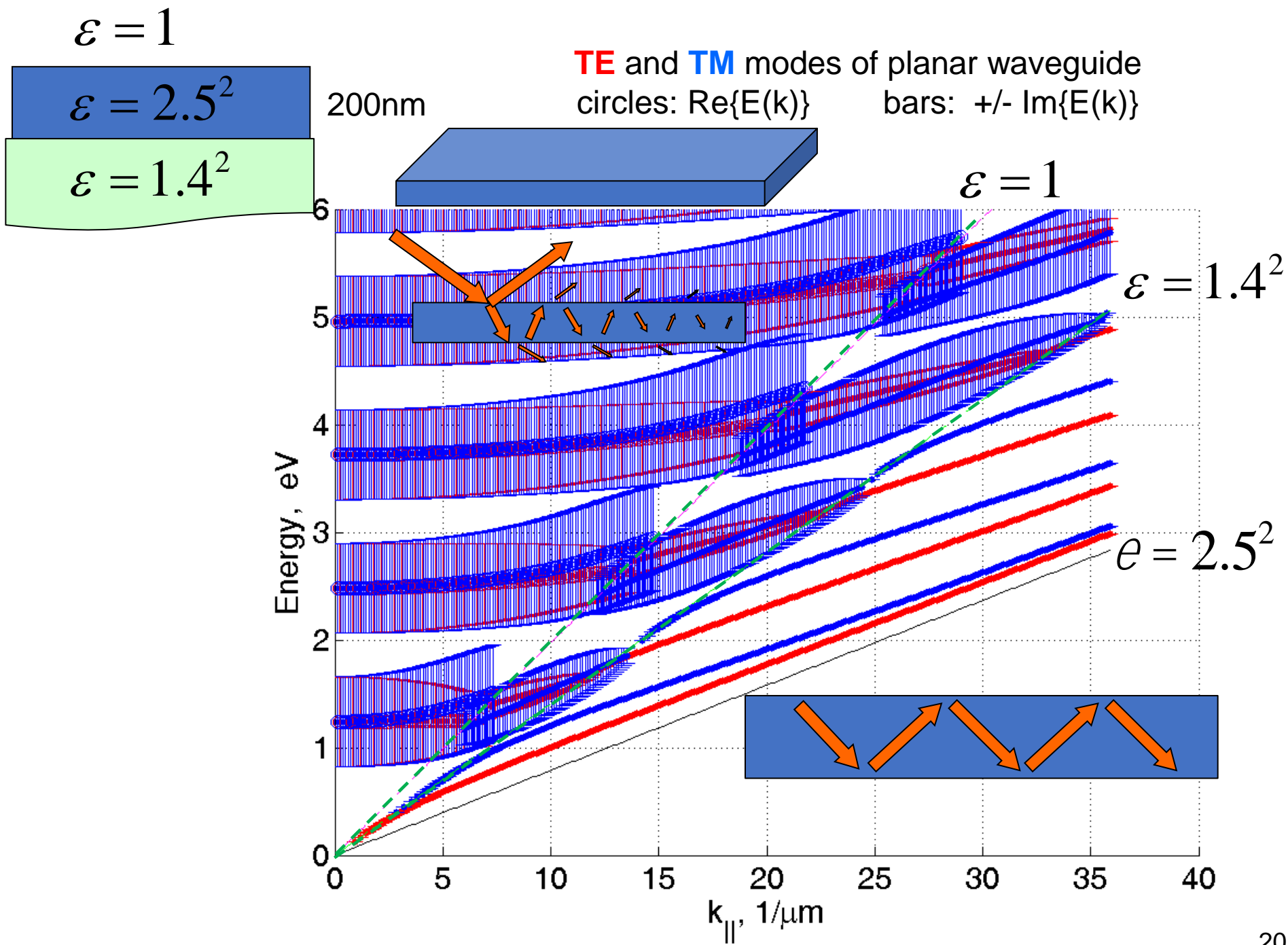


$$\epsilon = 1$$

$$\epsilon = 2.5^2$$

$$\epsilon = 1^2$$



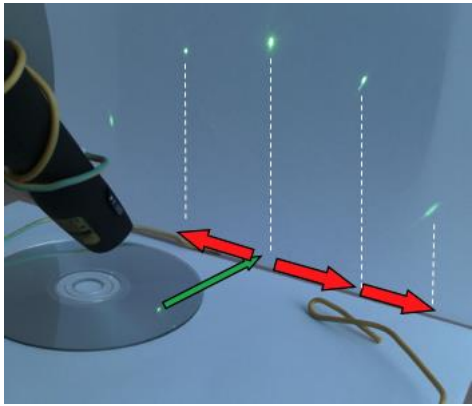
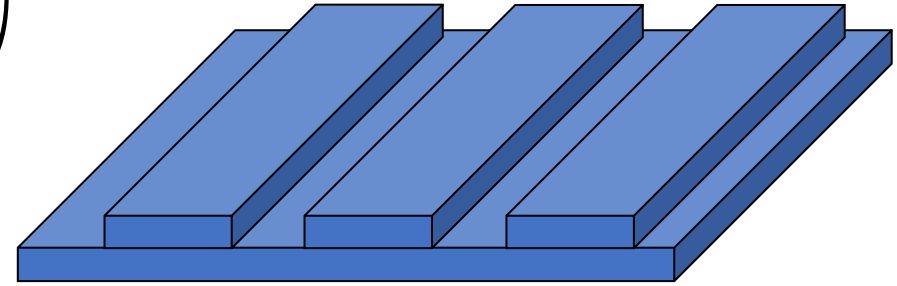




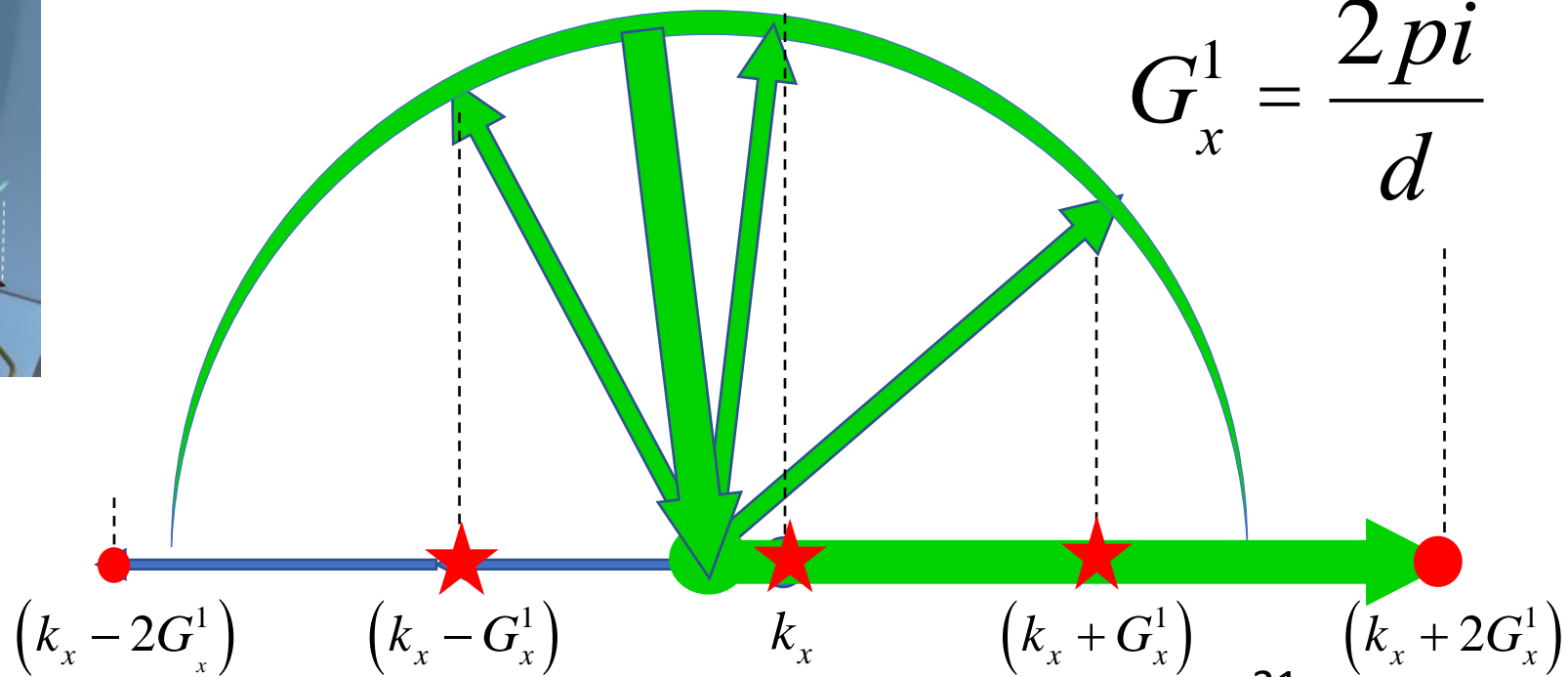
# Quasi-guided modes in *modulated* waveguide

$$k_z \left( G_x^n \right) = \sqrt{\frac{\epsilon_i}{c^2} \omega^2 - \left( k_x + G_x^n \right)^2}$$

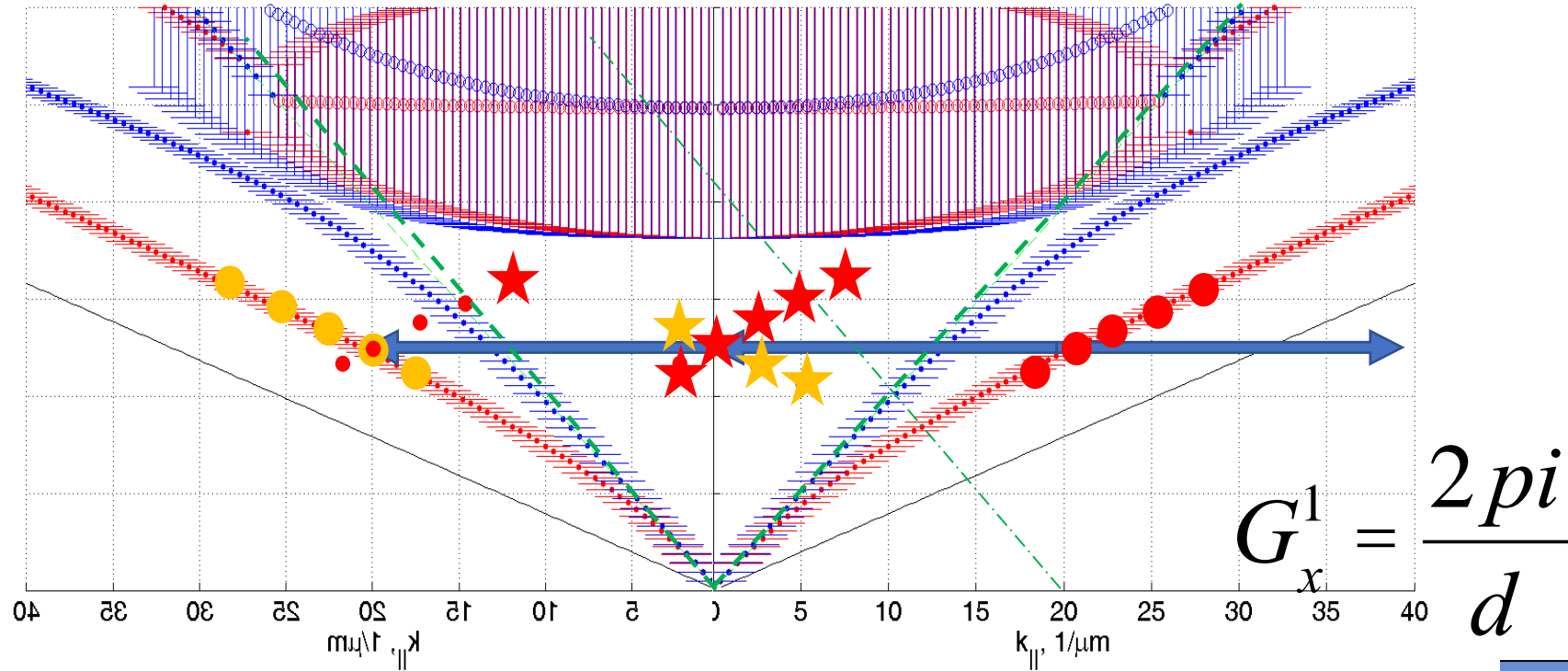
$$E = \hbar \omega$$



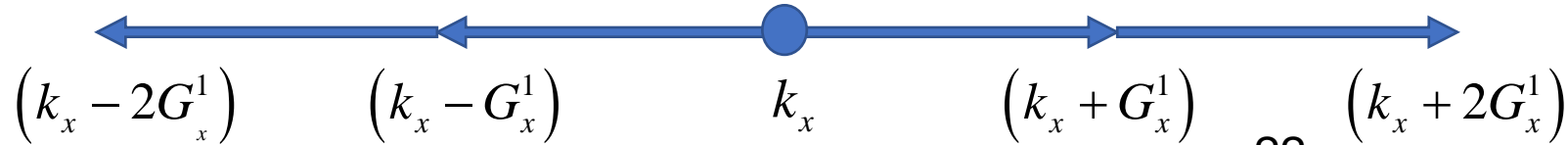
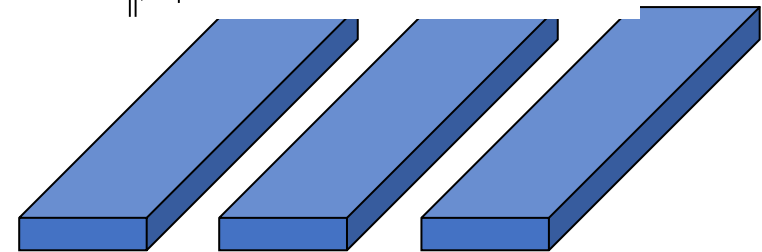
$$G_x^1 = \frac{2\pi}{d}$$



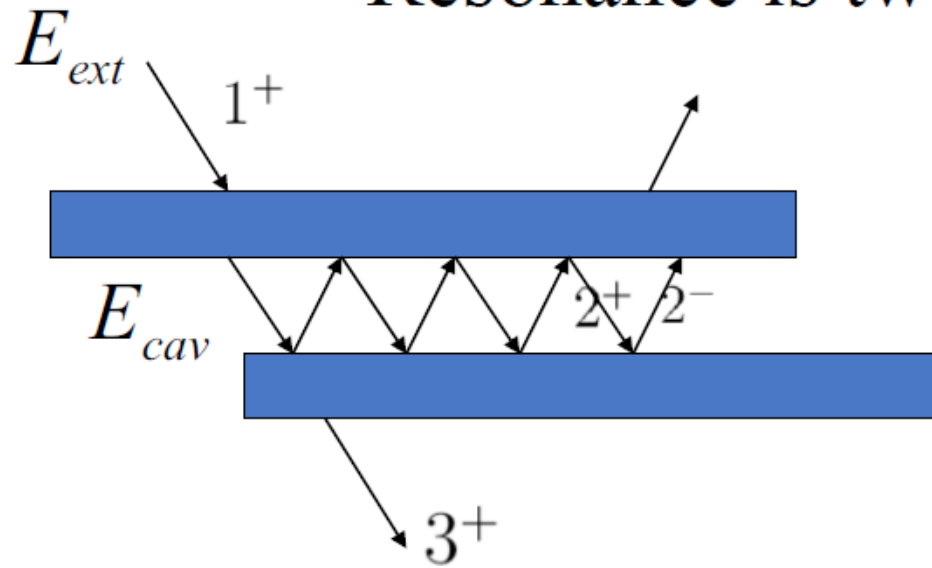
# Quasi-guided modes in *modulated* waveguide



$$k_z(G_x^n) = \sqrt{\frac{\epsilon_i}{c^2} \omega^2 - (k_x + G_x^n)^2}$$



# Resonance is two S-matrix problem



$$\begin{pmatrix} 1^- \\ 2^+ \end{pmatrix} = \begin{pmatrix} 1 \curvearrowright 1 & 1 \searrow 2 \\ 2 \swarrow 1 & 2 \curvearrowleft 2 \end{pmatrix} \begin{pmatrix} 1^+ \\ 2^- \end{pmatrix}$$

$$\begin{pmatrix} 2^- \\ 3^+ \end{pmatrix} = \begin{pmatrix} 2 \curvearrowright 2 & 2 \searrow 3 \\ 3 \swarrow 2 & 3 \curvearrowleft 3 \end{pmatrix} \begin{pmatrix} 2^+ \\ 3^- \end{pmatrix}$$

$$E_{cav} = \frac{1}{\omega - \omega_{cav}} \alpha E_{ext}$$

$$\omega_{cav} = \Omega_0 - i\gamma_0$$

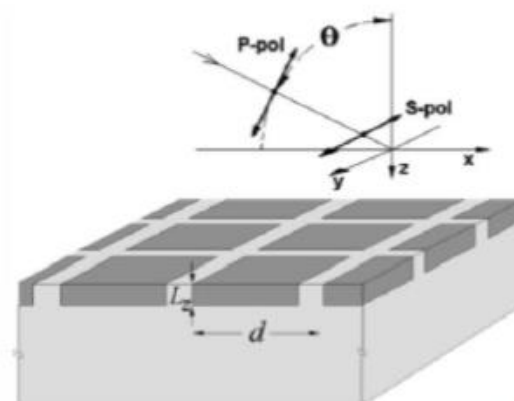
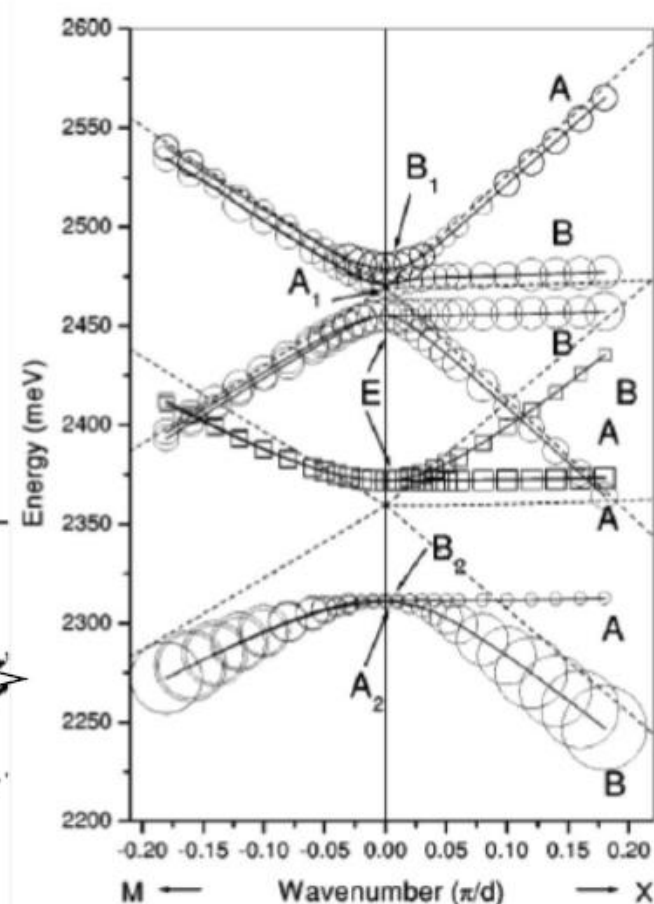
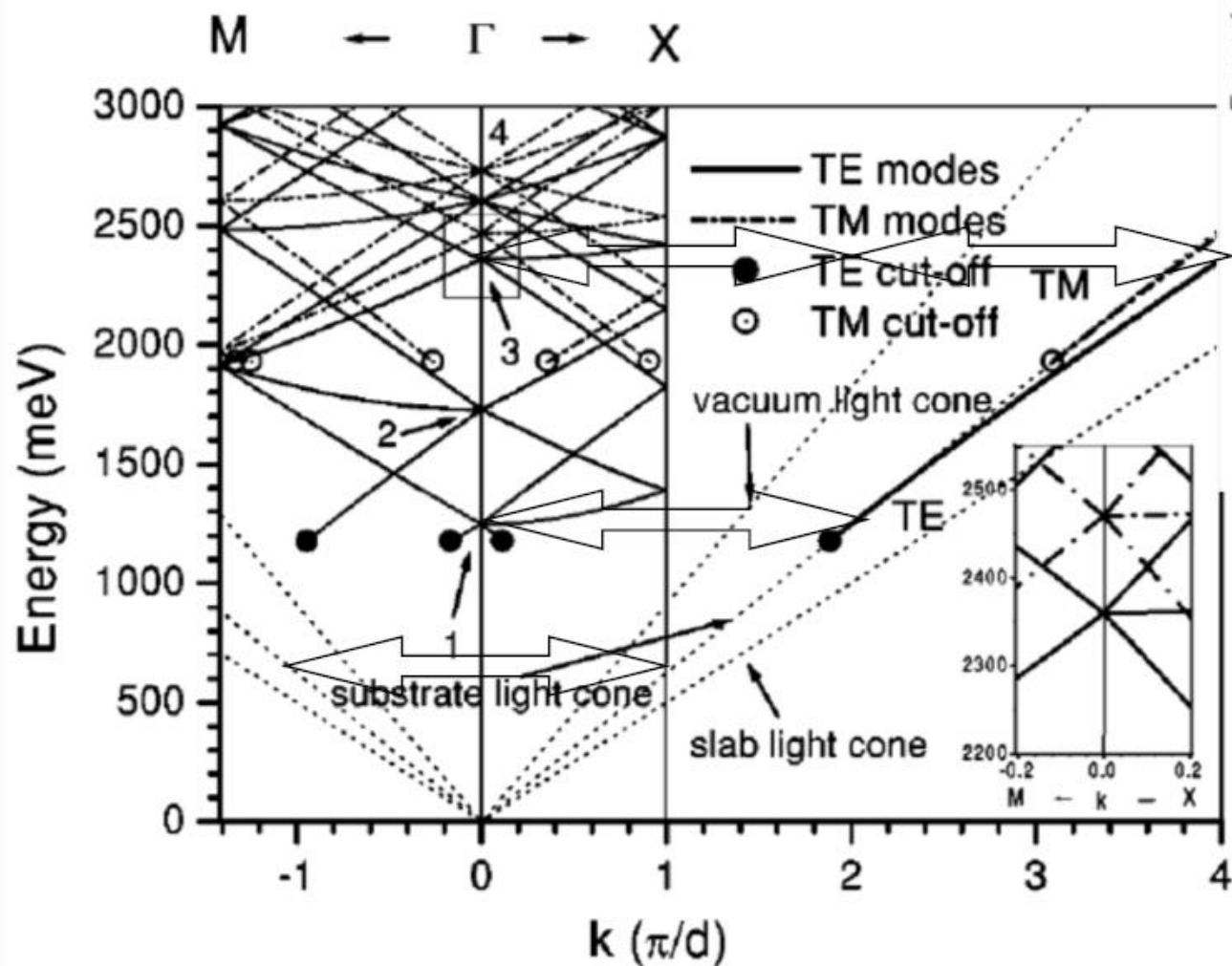
# Quasiguided modes and optical properties of photonic crystal slabs

S. G. Tikhodeev,<sup>1</sup> A. L. Yablonskii,<sup>1</sup> E. A. Muljarov,<sup>1,2</sup> N. A. Gippius,<sup>1,2</sup> and Teruya Ishihara<sup>2</sup>

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<sup>2</sup>Institute of Physical and Chemical Research (RIKEN), Wako 351-0198, Japan

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# Resonant mode coupling of optical resonances in stacked nanostructures

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Sergei G. Tikhodeev<sup>2</sup>, and Harald Giessen<sup>3</sup>

<sup>1</sup> LASMEA, University Blaise Pascal, 24 Avenue des Landais, 63177 Aubière Cedex, France

<sup>2</sup> A. M. Prokhorov General Physics Institute, RAS, Vavilova Street 38, Moscow 119991, Russia

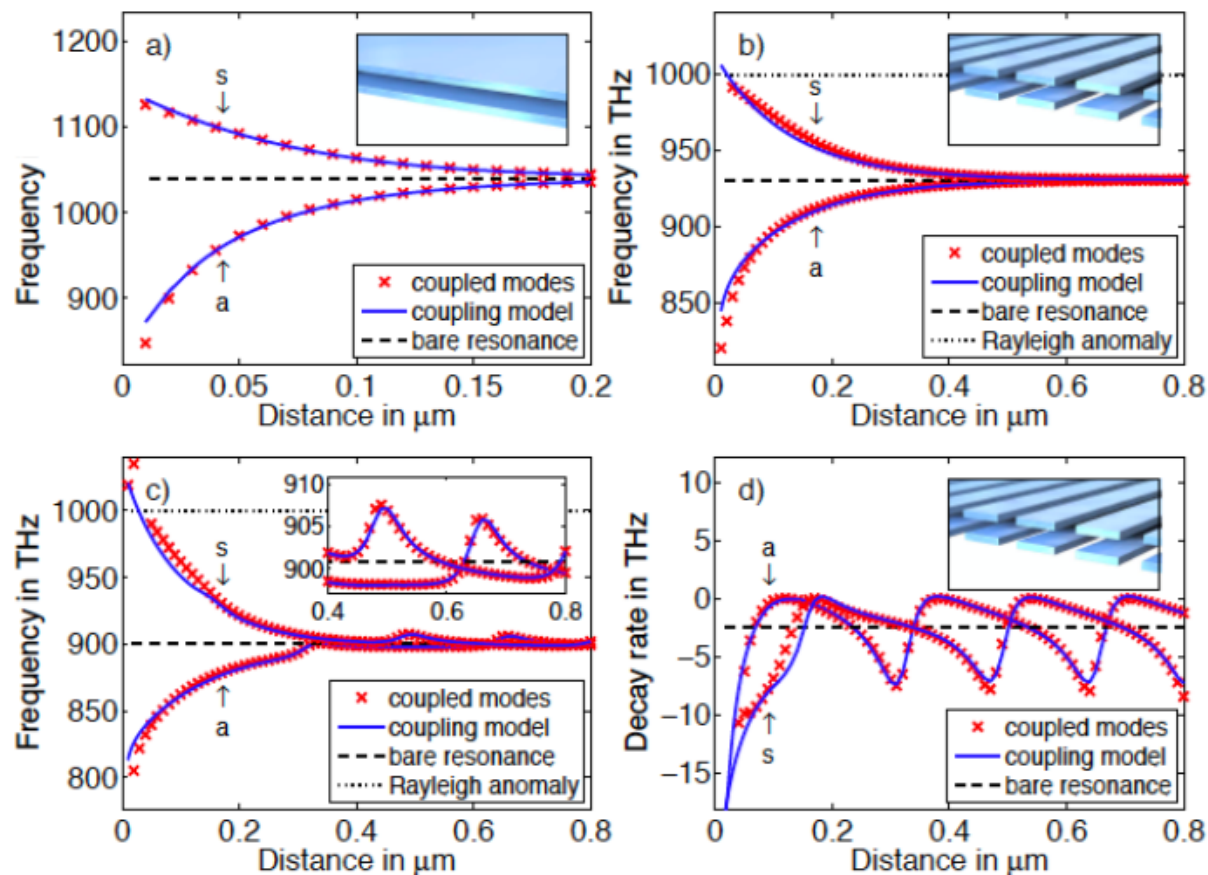
<sup>3</sup> 4<sup>th</sup> Physics Institute and Research Center SCOPE, University of Stuttgart, Pfaffenwaldring 57, 70550 Stuttgart, Germany

\*t.weiss@physik.uni-stuttgart.de

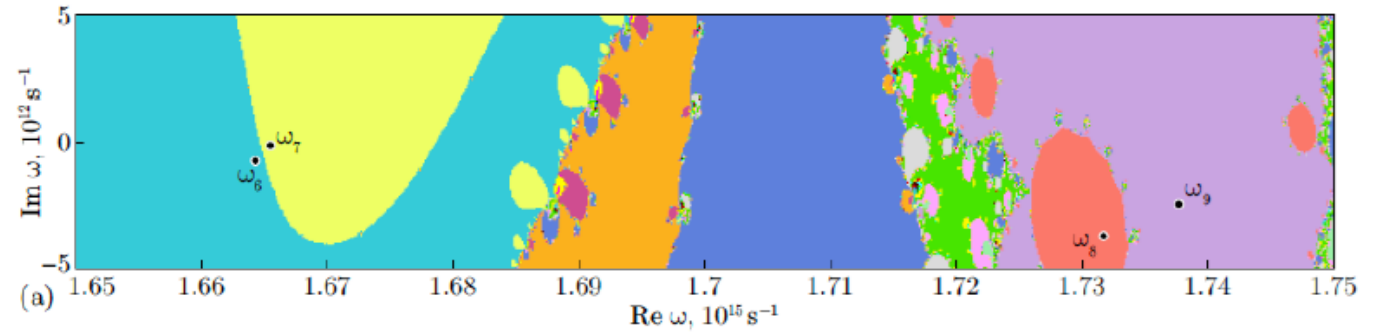
Received 4 Feb 2010; revised 24 Mar 2010; accepted 24 Mar 2010; published 26 Mar 2010  
29 March 2010 / Vol. 18, No. 7 / OPTICS EXPRESS 7569

$$(\omega - \omega_n^a)\alpha_n = \sum_{m=1}^N \langle I_{u,n}^a | \tilde{S}_{ud}^b | \mathbf{O}_{d,m}^a \rangle \alpha_m + \sum_{m=1}^M \langle I_{u,n}^a | \mathbf{O}_{u,m}^b \rangle \beta_m$$

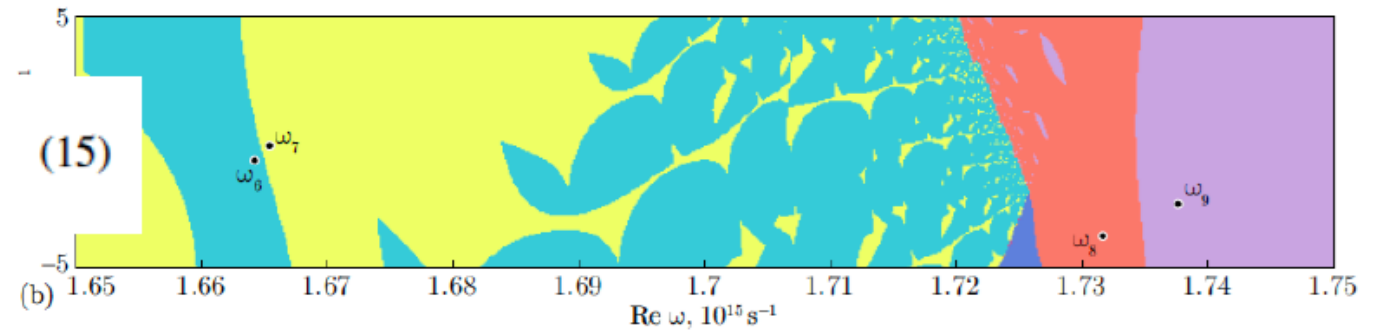
$$(\omega - \omega_n^b)\beta_n = \sum_{m=1}^N \langle I_{d,n}^b | \mathbf{O}_{d,m}^a \rangle \alpha_m + \sum_{m=1}^M \langle I_{d,n}^b | \tilde{S}_{du}^a | \mathbf{O}_{u,m}^b \rangle \beta_m$$



$$1/\det S(\omega) = 0. \quad (10)$$



$$\omega_{n+1} = \omega_n - \min \text{eig} \left( S^{-1}(\omega_n), \frac{dS^{-1}}{d\omega} \Big|_{\omega_n} \right), \quad (15)$$



$$\omega_{n+1} = \omega_n + 2 \min \text{eig} (U_r^\dagger S'(\omega_n) V_r \Sigma_r^{-1}). \quad (19)$$

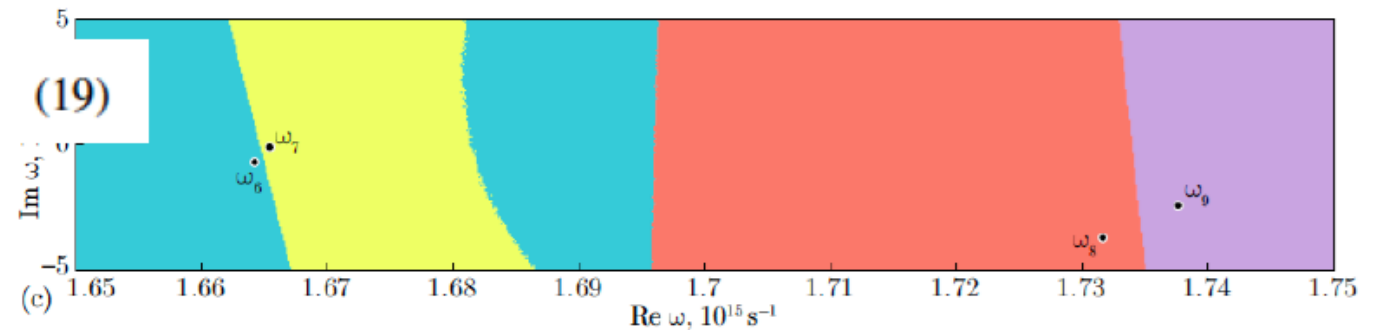
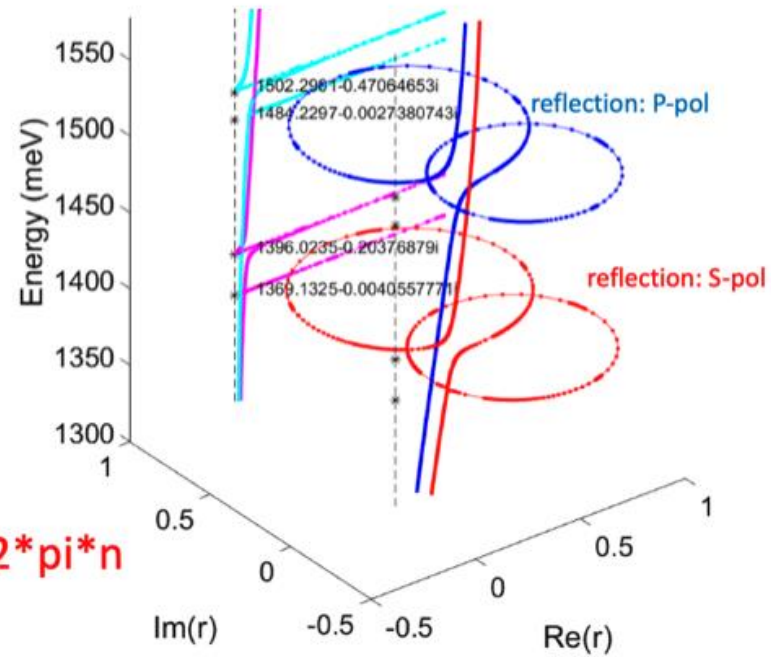
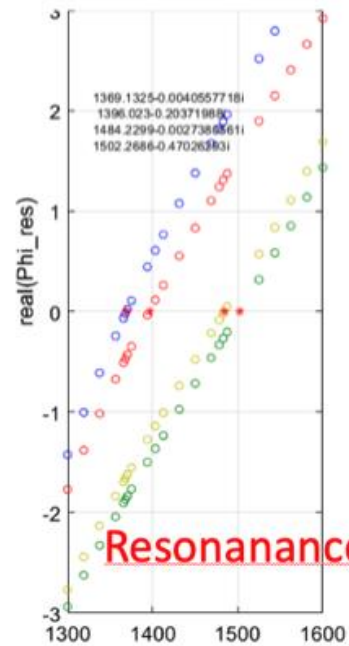
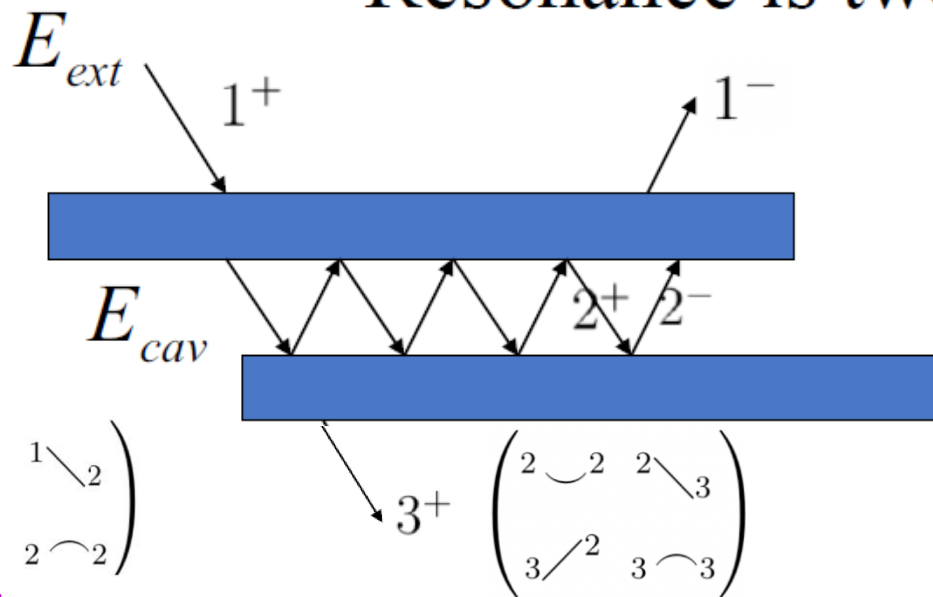


Fig. 2. (Color online) Basins of attraction (different colors denote different attraction poles): (a) Newton for Eq. (10); (b) Eq. (15); (c) Eq. (19).

## 2. Сколько можно считать одно и то же ?!



# Resonance is two S-matrix problem

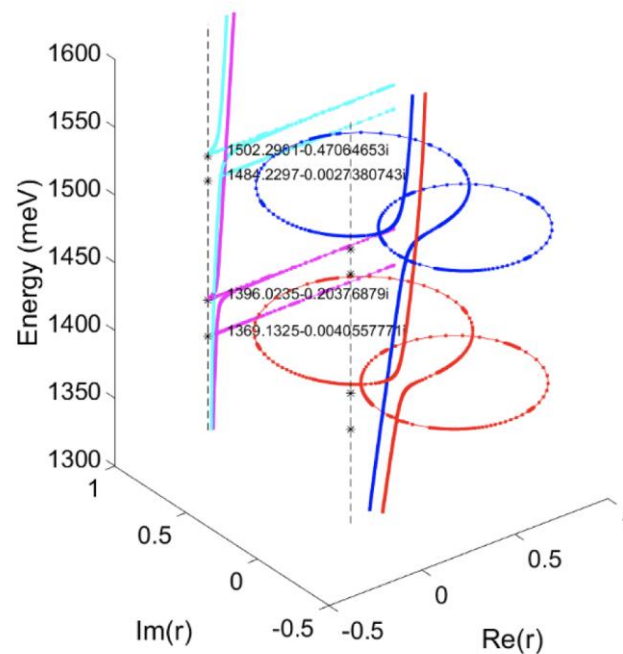
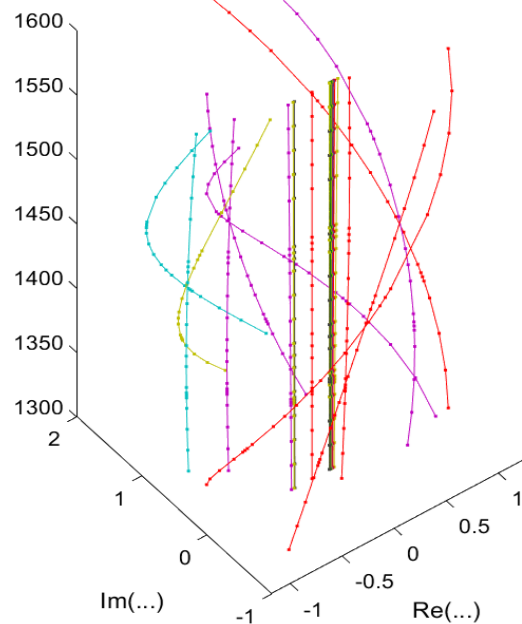
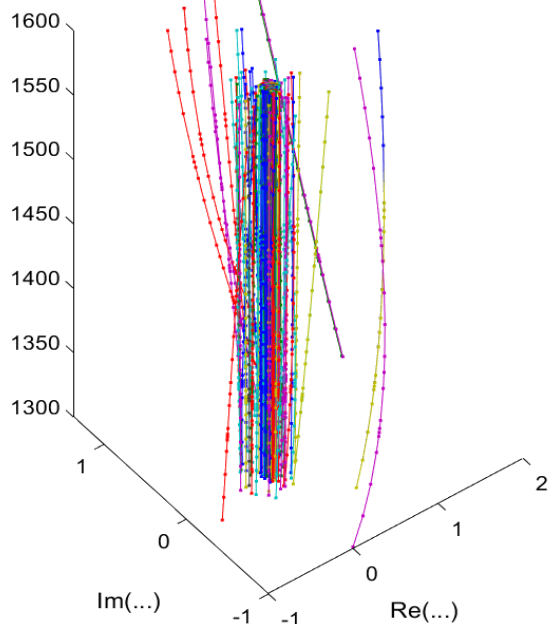


$$\begin{pmatrix} 1 \curvearrowright 1 & 1 \searrow 2 \\ 2 \swarrow 1 & 2 \curvearrowleft 2 \end{pmatrix}$$

$$\begin{pmatrix} 2 \curvearrowright 2 & 2 \searrow 3 \\ 3 \swarrow 2 & 3 \curvearrowleft 3 \end{pmatrix}$$

$$\begin{pmatrix} 1^- \\ 2^+ \end{pmatrix} = \begin{pmatrix} 1 \curvearrowright 1 & 1 \searrow 2 \\ 2 \swarrow 1 & 2 \curvearrowleft 2 \end{pmatrix} \begin{pmatrix} 1^+ \\ 2^- \end{pmatrix}$$

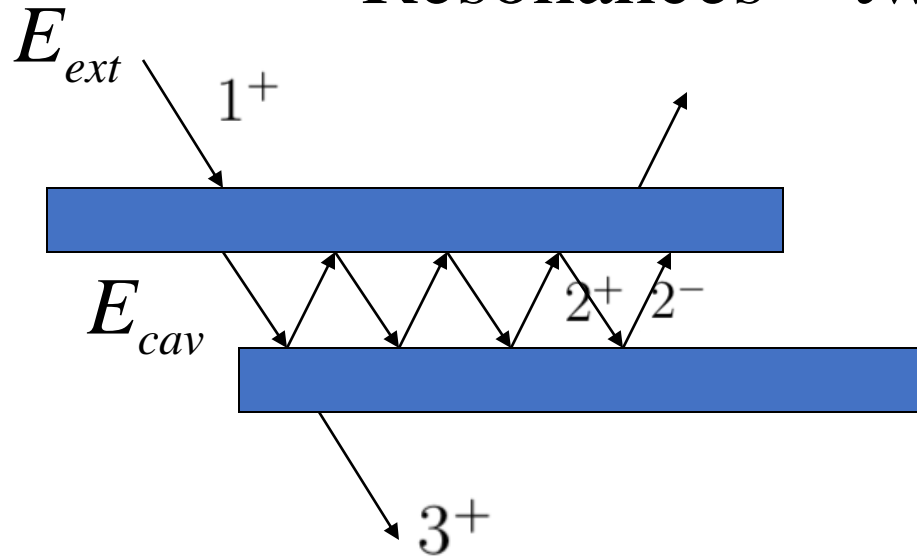
$$\begin{pmatrix} 2^- \\ 3^+ \end{pmatrix} = \begin{pmatrix} 2 \curvearrowright 2 & 2 \searrow 3 \\ 3 \swarrow 2 & 3 \curvearrowleft 3 \end{pmatrix} \begin{pmatrix} 2^+ \\ 3^- \end{pmatrix}$$



$$\begin{pmatrix} 1 \curvearrowright 1 & 1 \searrow 3 \\ 3 \swarrow 1 & 3 \curvearrowleft 3 \end{pmatrix}$$



# Resonances – two S-matrix problem



$$\begin{pmatrix} 1^- \\ 2^+ \end{pmatrix} = \begin{pmatrix} 1 \smile 1 & 1 \searrow 2 \\ 2 \swarrow 1 & 2 \frown 2 \end{pmatrix} \begin{pmatrix} 1^+ \\ 2^- \end{pmatrix}$$

$$\begin{pmatrix} 2^- \\ 3^+ \end{pmatrix} = \begin{pmatrix} 2 \smile 2 & 2 \searrow 3 \\ 3 \swarrow 2 & 3 \frown 3 \end{pmatrix} \begin{pmatrix} 2^+ \\ 3^- \end{pmatrix}$$

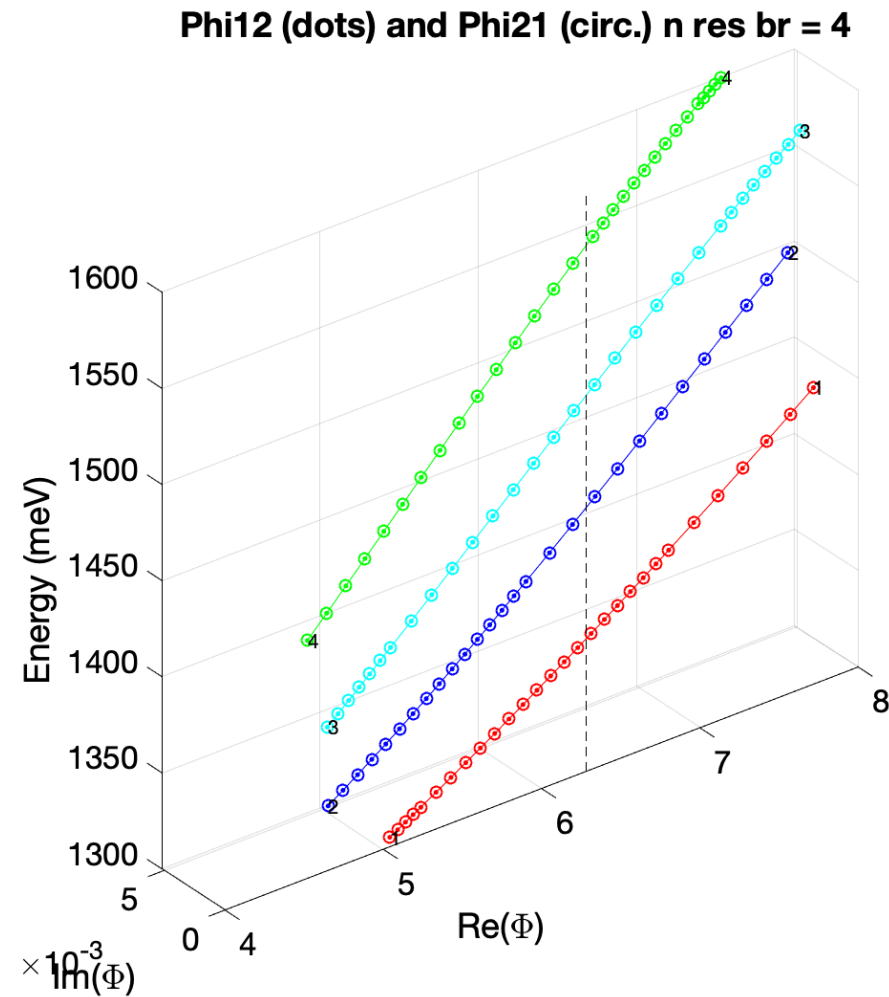
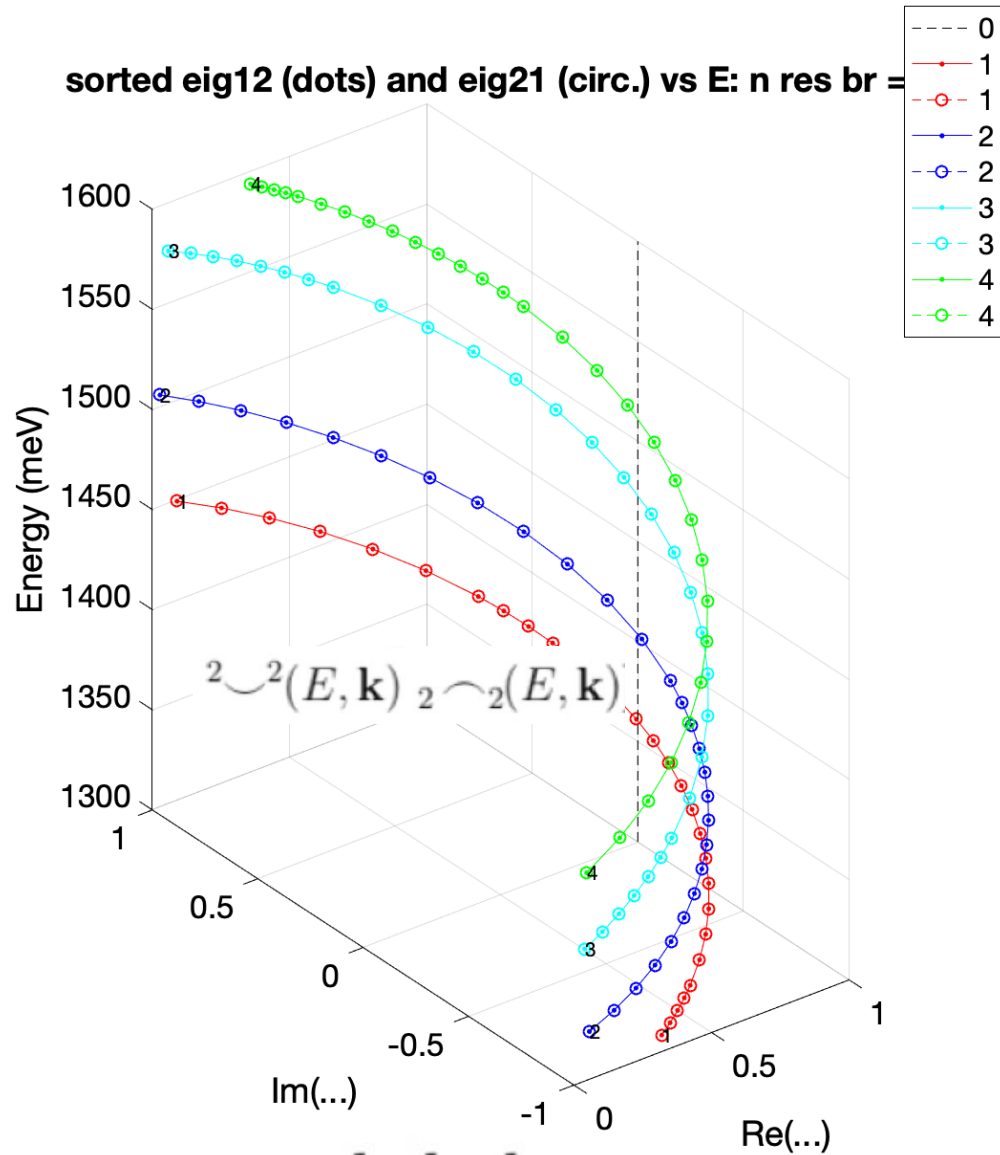
$$\begin{pmatrix} 1 \smile 1 & 1 \searrow 3 \\ 3 \swarrow 1 & 3 \frown 3 \end{pmatrix} = \begin{pmatrix} 1 \smile 1 + 1 \searrow 2 \smile 2 \swarrow 1 + \dots & 1 \searrow 2 \searrow 3 + 1 \searrow 2 \smile 2 \frown 2 \searrow 3 + \dots \\ 3 \swarrow 2 \swarrow 1 + 3 \swarrow 2 \frown 2 \smile 2 \swarrow 1 + \dots & 3 \frown 3 + 3 \swarrow 2 \frown 2 \searrow 3 + \dots \end{pmatrix}$$

$$[1 - 2 \smile 2(E, \mathbf{k}) 2 \frown 2(E, \mathbf{k})]^{-1}$$

$$[1 - 2 \frown 2(E, \mathbf{k}) 2 \smile 2(E, \mathbf{k})]^{-1}$$

$$S_1^{ab} t^b s_2^{ba} t^a X = X e^{i\Phi}, \quad \Phi = \Phi(k_0, k_{||})$$

Roundtrip phase  
of resonant mode



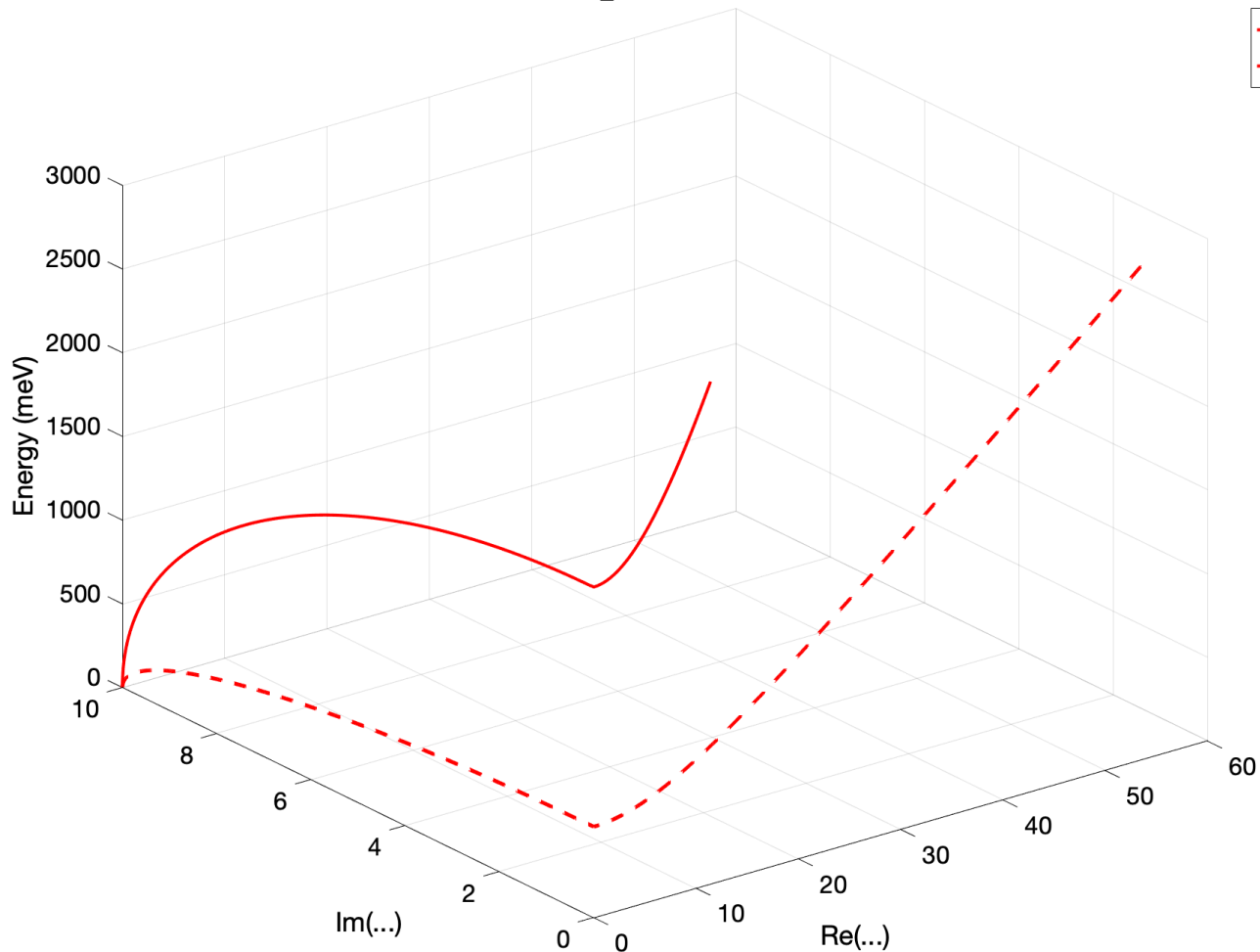
$$S_1^{ab} t^b s_2^{ba} t^a$$

= 1.

Resonance

$$\Phi(k_0, k_{||}) = 2 * \pi * n$$

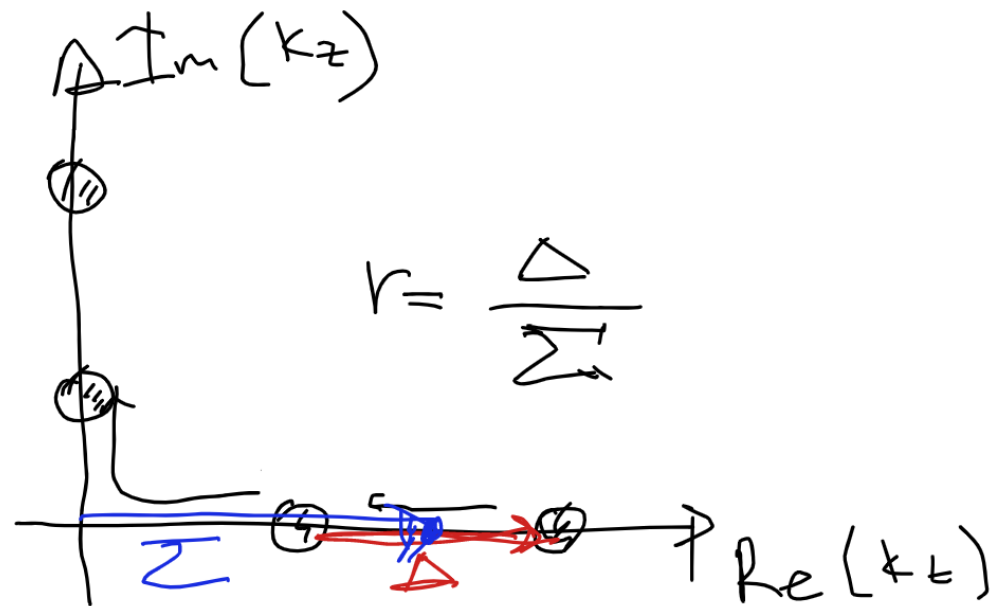
**K<sub>z</sub> vs Energy**



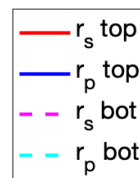
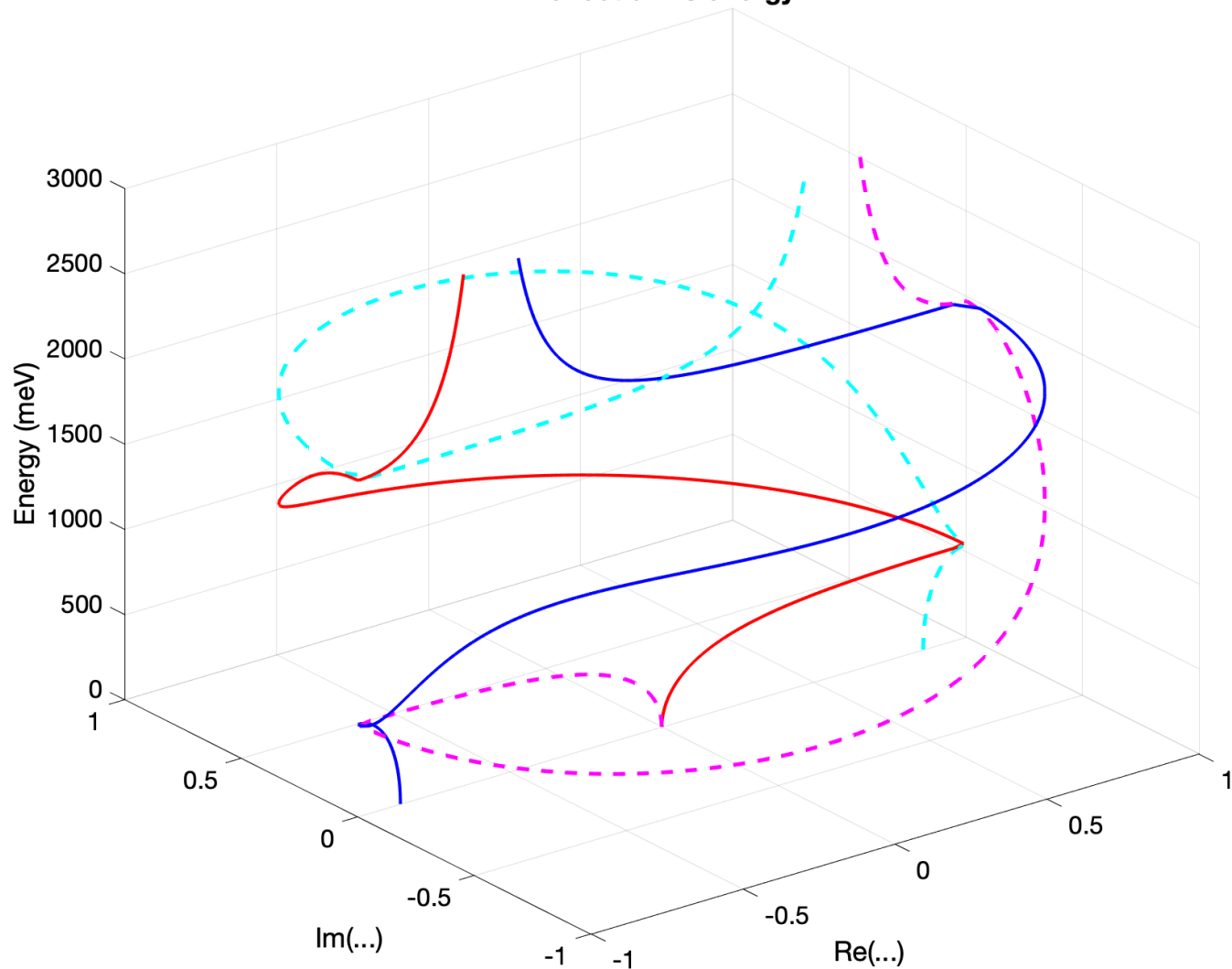
— K<sub>z</sub> top  
 - - - K<sub>z</sub> bot

$$k_0 = \frac{\omega}{c}, \quad k_z = \sqrt{k_0^2 \epsilon - k_{\parallel}^2}, \quad k_y = 0$$

$$r_s^w = \frac{k_z^w - k_z^i}{k_z^w + k_z^i} \quad r_p^w = \frac{k_z^i \epsilon^w - k_z^w \epsilon^i}{k_z^i \epsilon^w + k_z^w \epsilon^i}$$

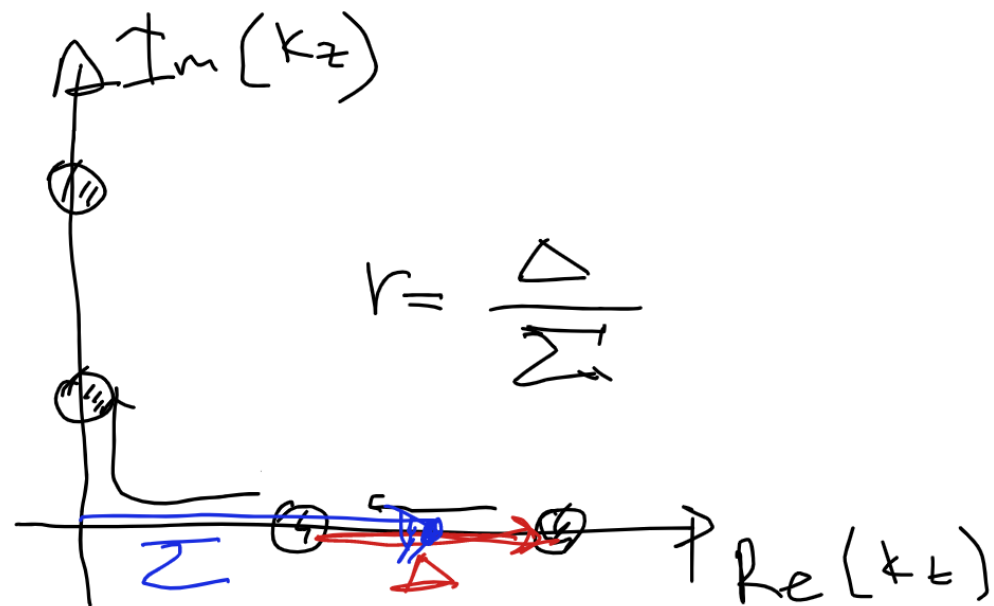


Reflection vs energy



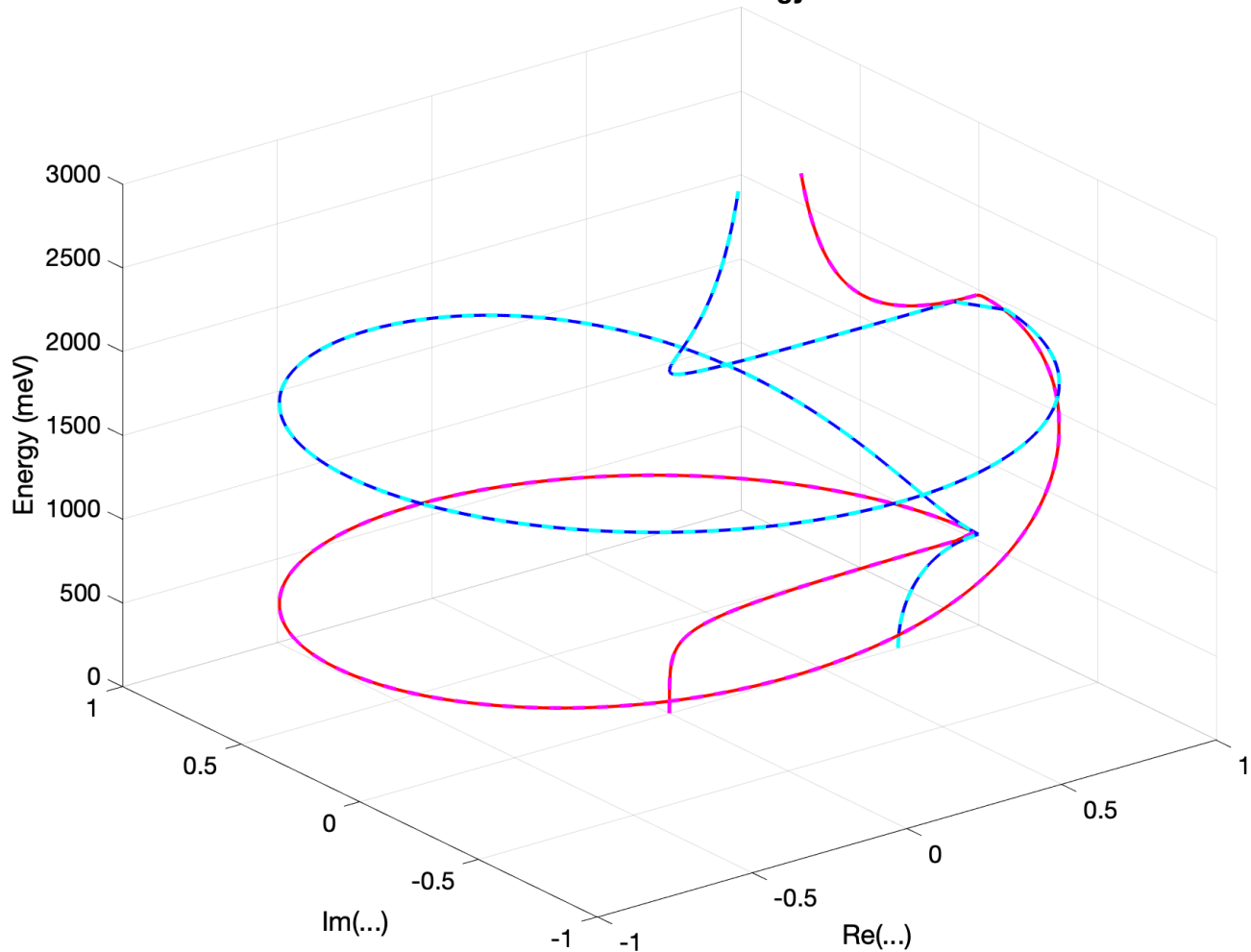
$$k_0 = \frac{\omega}{c}, \quad k_z = \sqrt{k_0^2 \epsilon - k_{\parallel}^2}, \quad k_y = 0$$

$$r_s^w = \frac{k_z^w - k_z^i}{k_z^w + k_z^i} \quad r_p^w = \frac{k_z^i \epsilon^w - k_z^w \epsilon^i}{k_z^i \epsilon^w + k_z^w \epsilon^i}$$



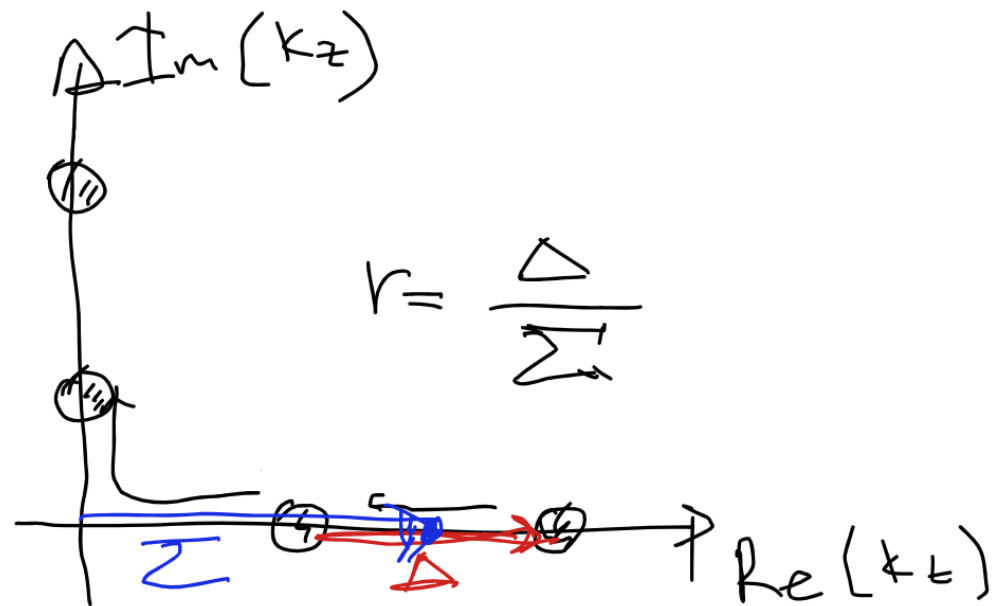


Reflection<sup>2</sup> vs energy



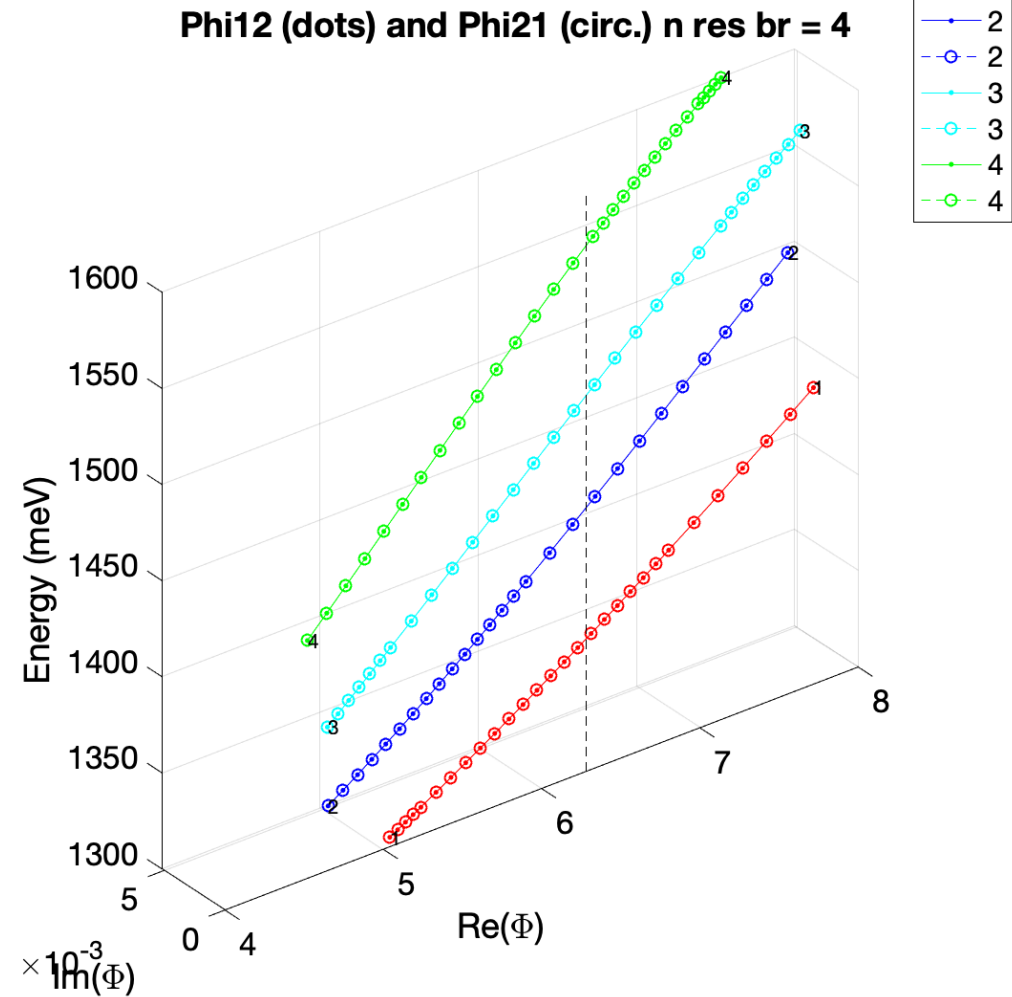
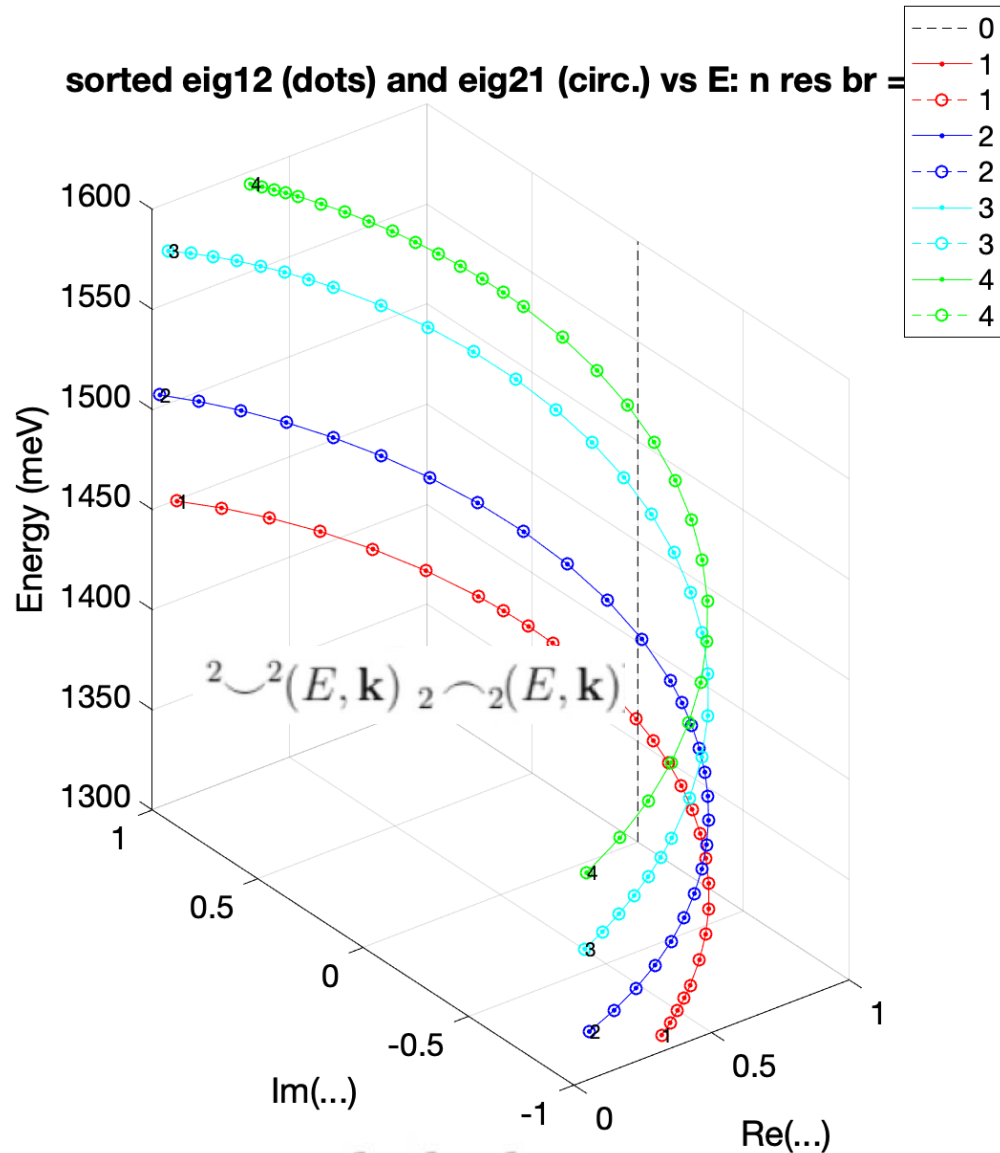
$$k_0 = \frac{\omega}{c}, \quad k_z = \sqrt{k_0^2 \epsilon - k_{||}^2}, \quad k_y = 0$$

$$r_s^w = \frac{k_z^w - k_z^i}{k_z^w + k_z^i} \quad r_p^w = \frac{k_z^i \epsilon^w - k_z^w \epsilon^i}{k_z^i \epsilon^w + k_z^w \epsilon^i}$$



$$S_1^{ab} t^b s_2^{ba} t^a X = X e^{i\Phi}, \quad \Phi = \Phi(k_0, k_{||})$$

Roundtrip phase  
of resonant mode



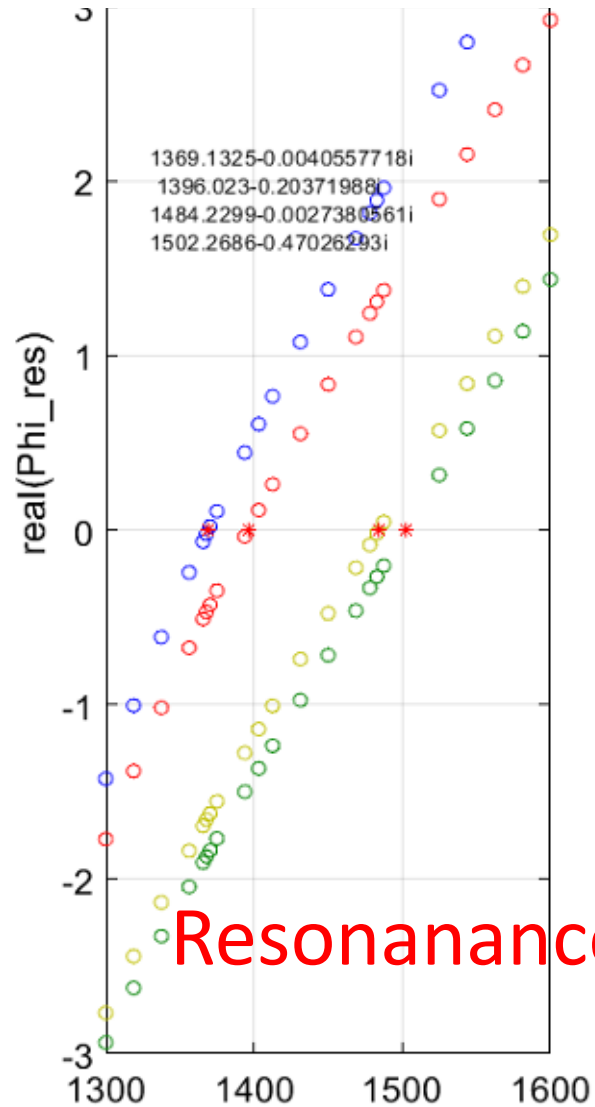
$$S_1^{ab} t^b s_2^{ba} t^a$$

= 1.

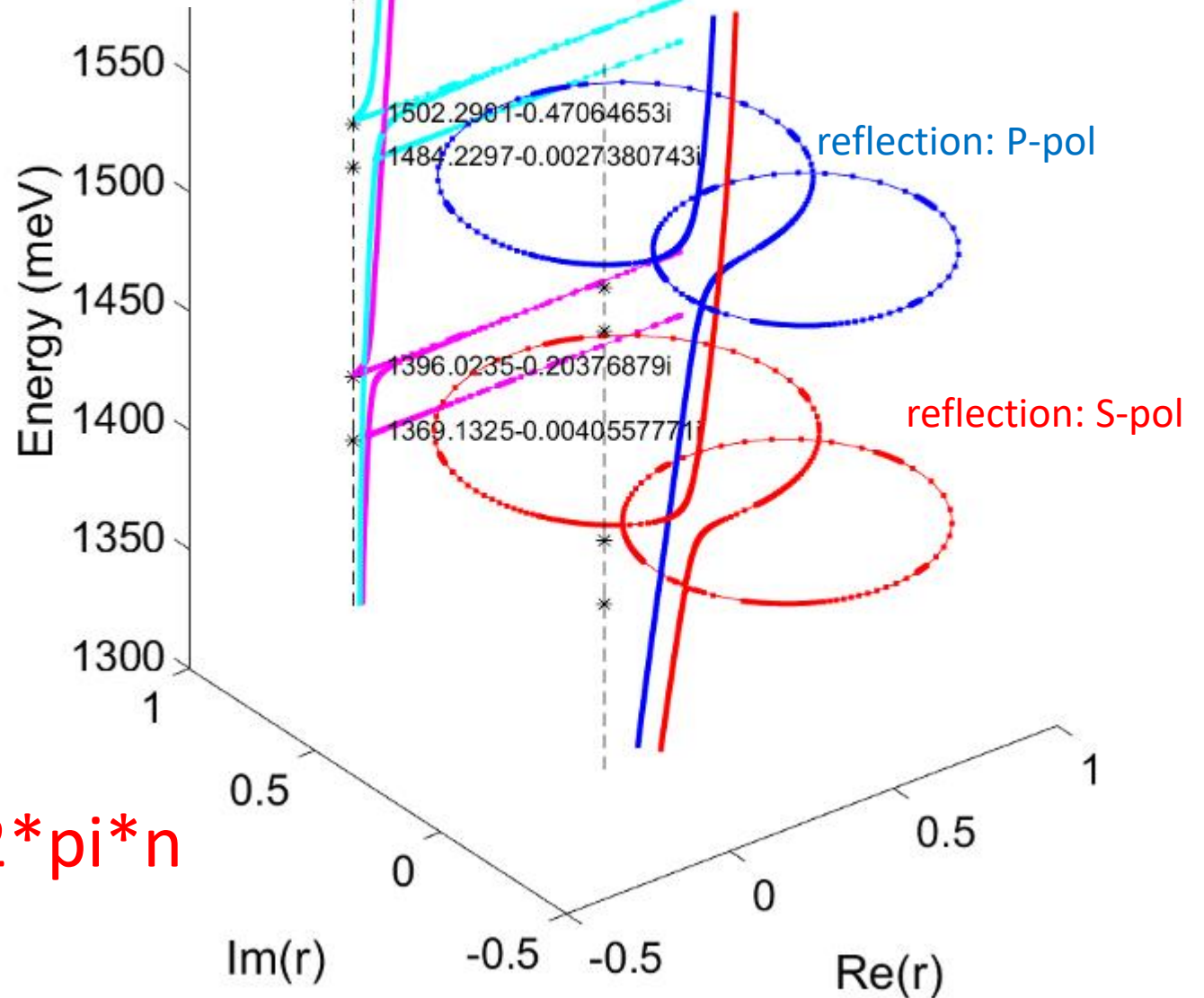
Resonance

$$\Phi(k_0, k_{||}) = 2 * \pi * n$$

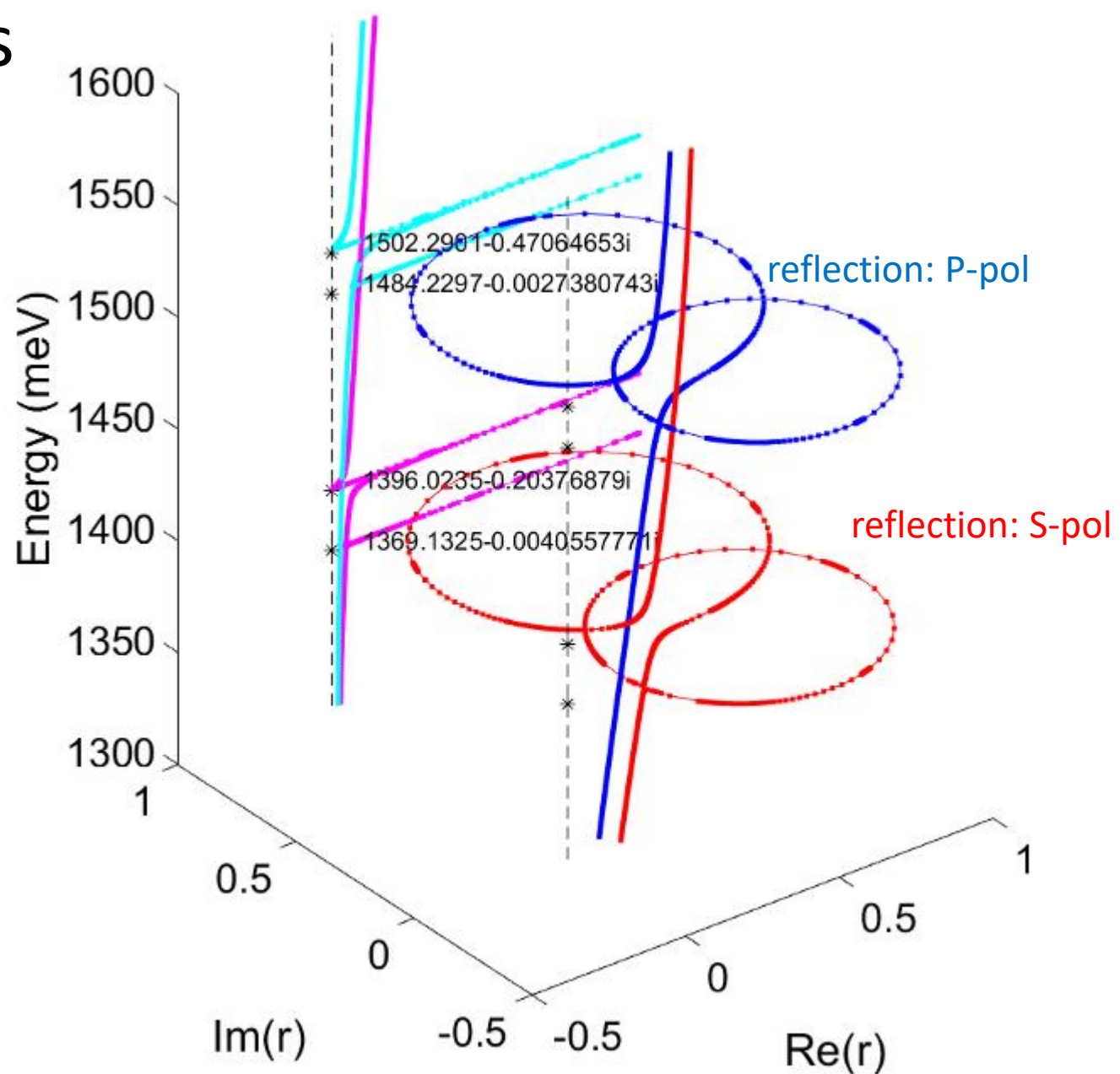
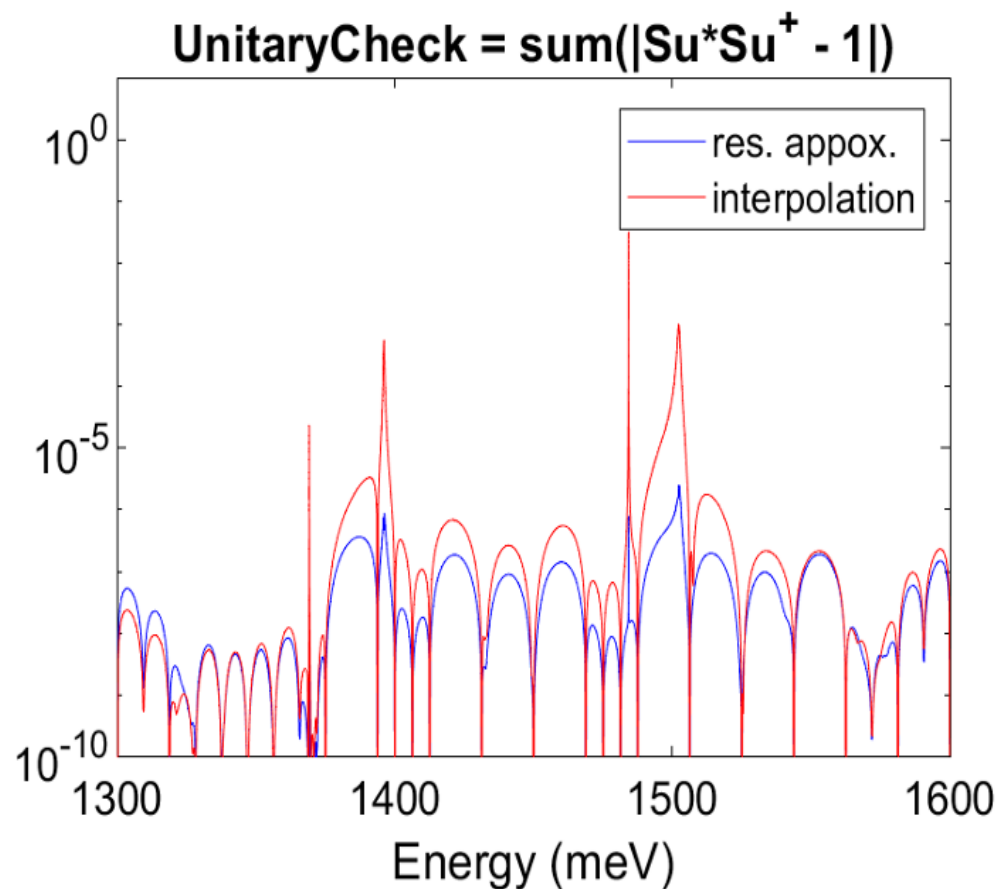
# Roundtrip phase of resonant mode is smooth function of energy



Resonance : Phase =  $2 \cdot \pi \cdot n$

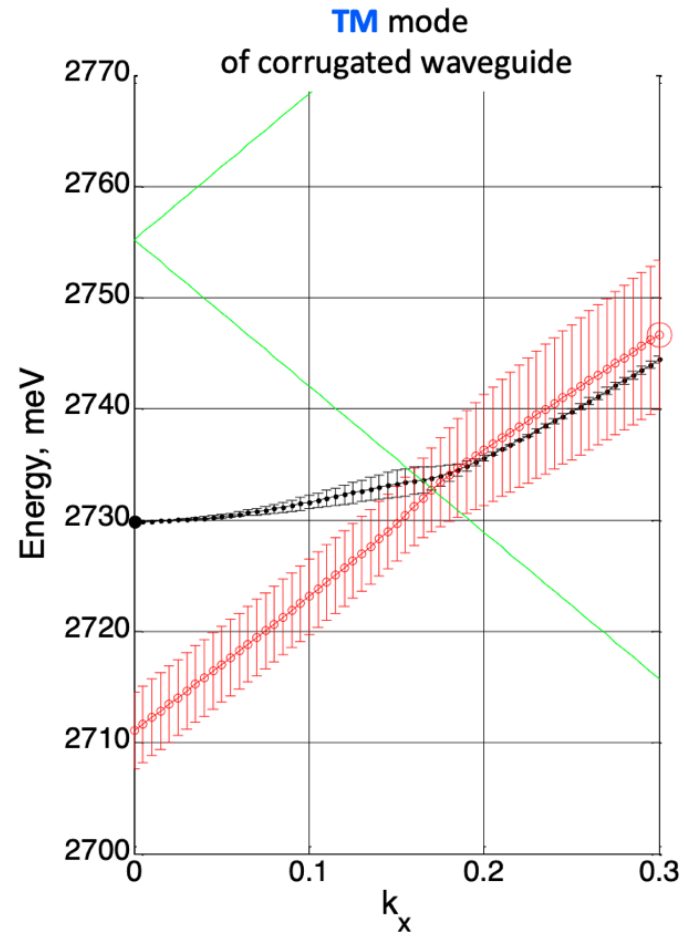
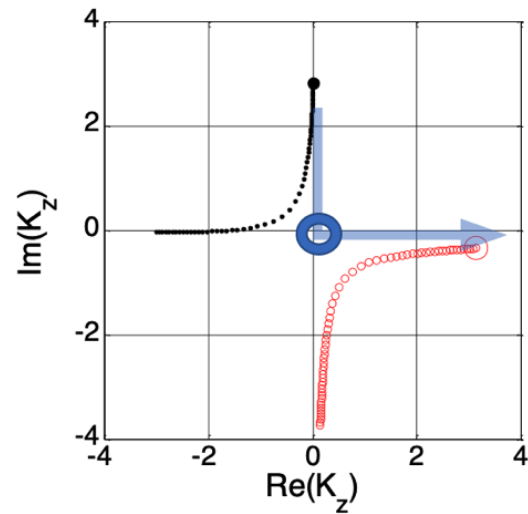
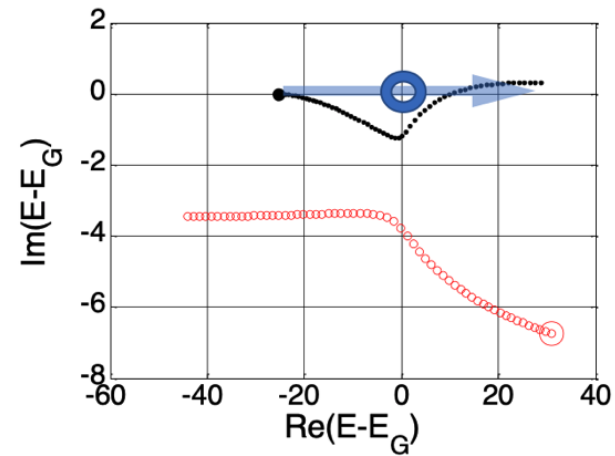


# Reflection near resonances is well described





# 3. Как раскалываются резонансы ?



## ON THRESHOLD PHENOMENA IN CLASSICAL ELECTRODYNAMICS

B. M. BOLOTOVSKIĬ and A. N. LEBEDEV

P. N. Lebedev Physical Institute, Academy of Sciences, U.S.S.R.

Submitted April 21, 1967

Zh. Eksp. Teor. Fiz. 53, 1349-1352 (October, 1967)

It is shown that for a wide class of problems in electrodynamics, the behavior of the amplitudes and phases at the threshold for the production of new proper waves can be determined from the conservation laws. The method proposed, which is analogous to the quantum theory of many-channel nuclear reactions, is employed for an explanation of the Wood anomalies.

$$\psi_n^\pm = \exp \left[ iy \left( k_y - \frac{2\pi n}{d} \right) \pm iz \left[ \frac{\omega^2}{c^2} - \left( k_y - \frac{2\pi n}{d} \right)^2 \right]^{1/2} \right],$$

$$n = 0, \pm 1, \pm 2 \dots$$

$$\kappa_j = \left[ \frac{\omega^2}{c^2} - \left( k_y - \frac{2\pi j}{d} \right)^2 \right]^{1/2}$$

CONTRIBUTION TO THE THEORY OF THRESHOLD PHENOMENA IN  
DIFFRACTION OF ELECTROMAGNETIC WAVES

B. M. BOLOTOVSKIĬ and K. I. KUGEL'

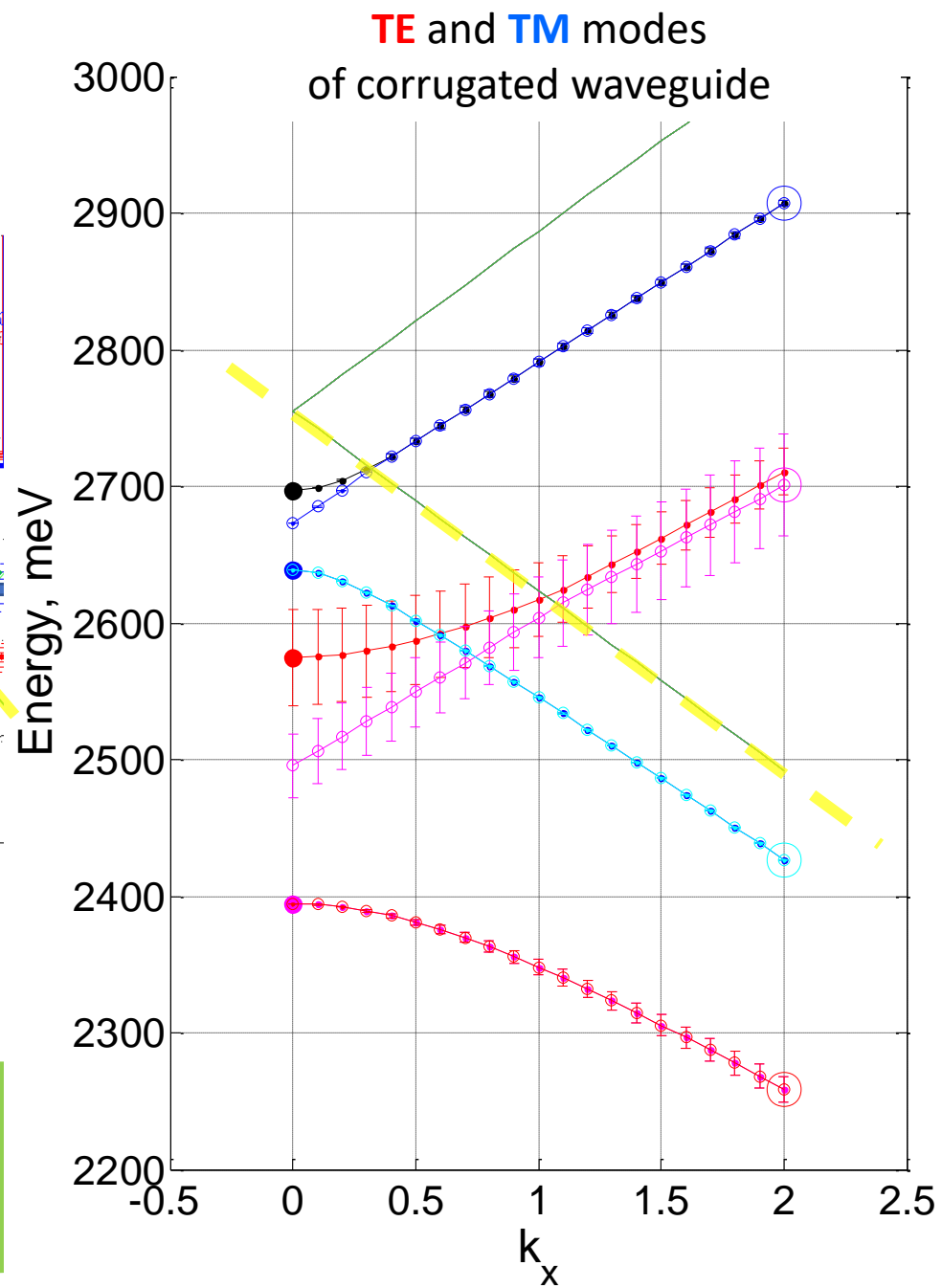
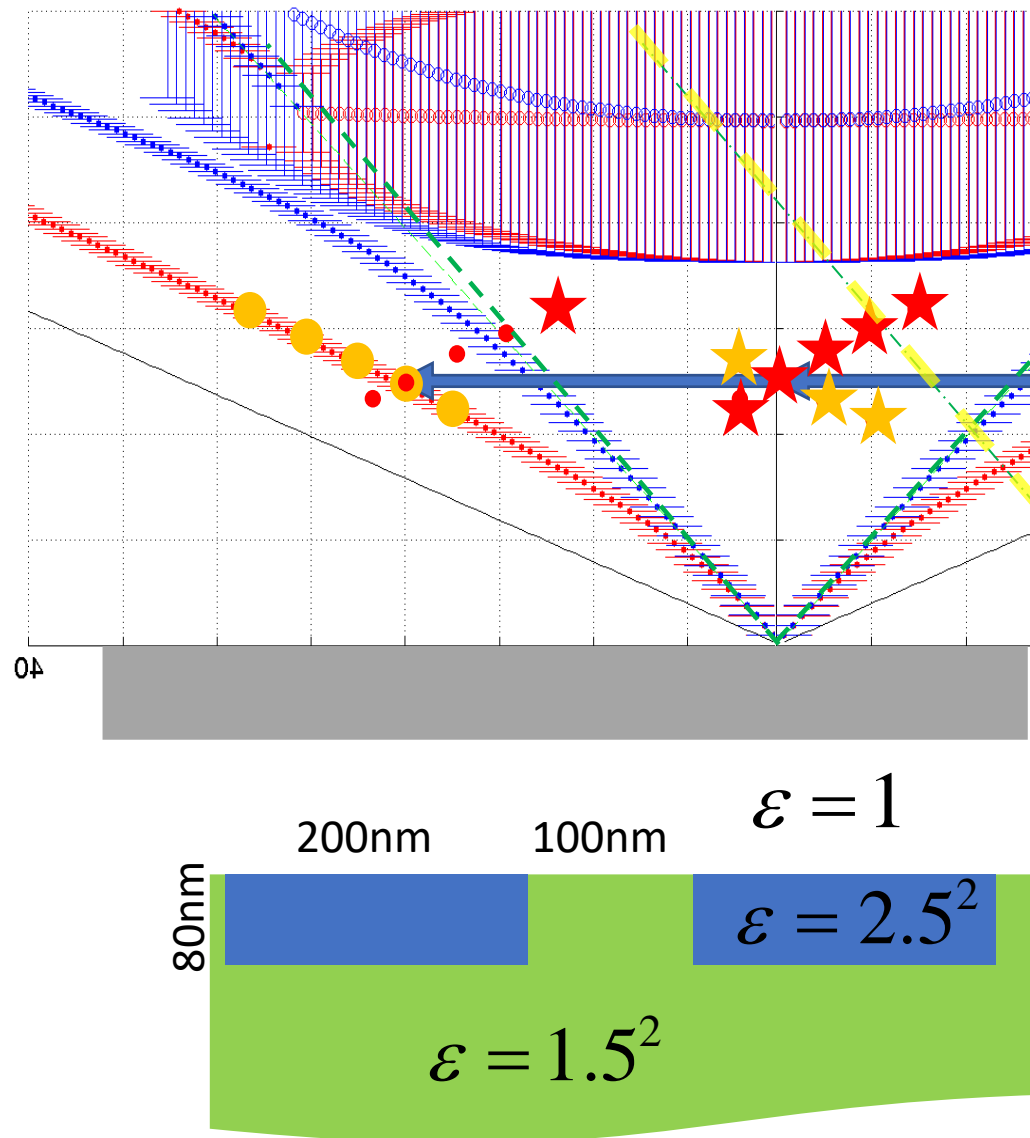
P. N. Lebedev Physics Institute, U.S.S.R. Academy of Sciences

Submitted November 11, 1968

Zh. Eksp. Teor. Fiz. 57, 165-174 (July, 1969)

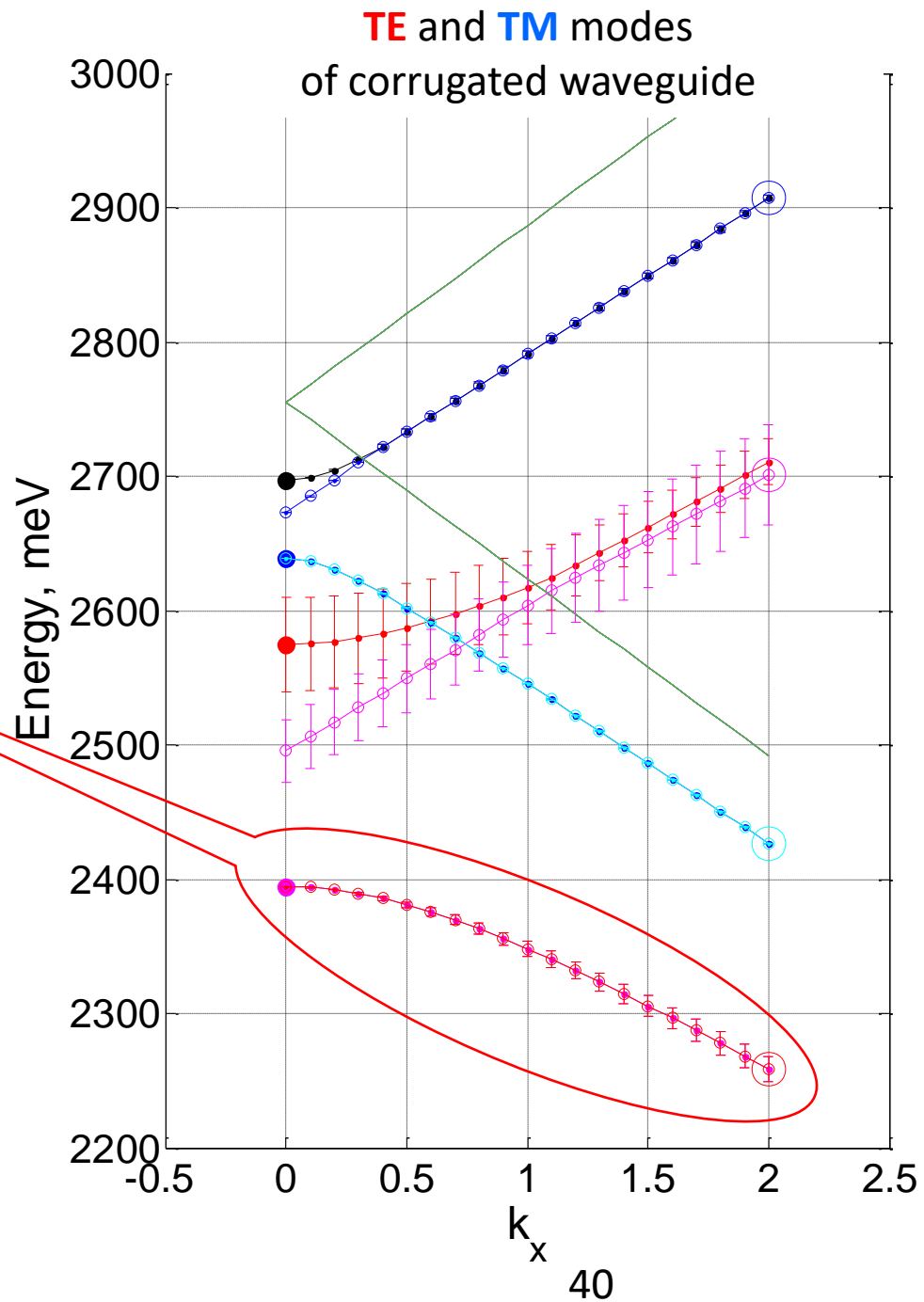
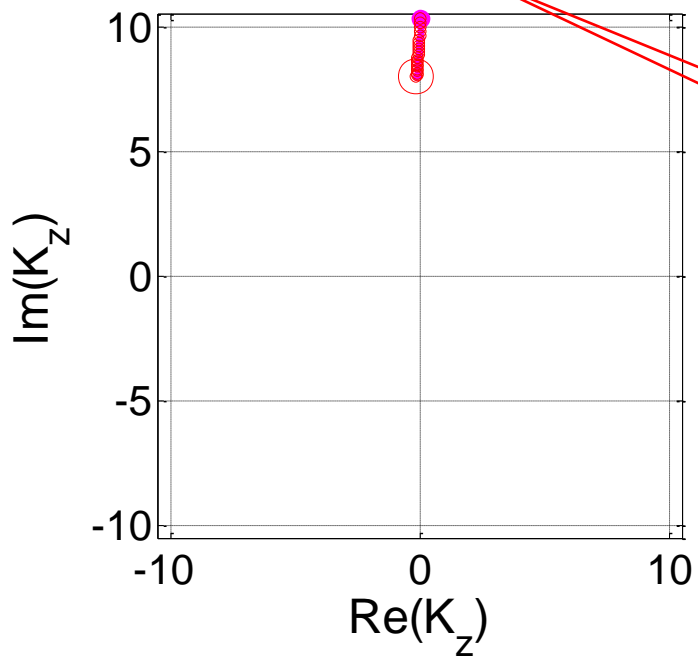
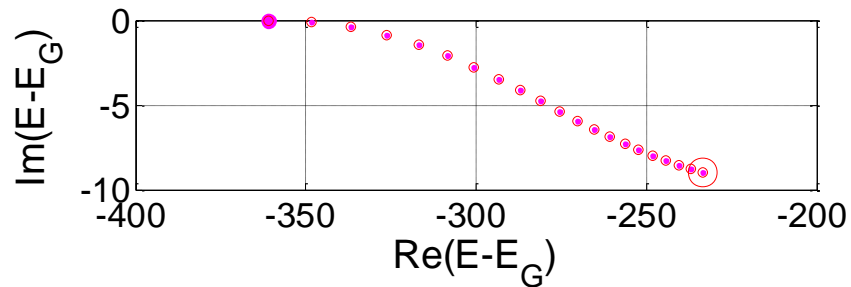
The behavior of the amplitudes and phases of electromagnetic waves at the threshold of appearance of a new electromagnetic wave (a spectrum of a new order) is considered for the case of scattering by a transparent diffraction lattice or by the open end of a cylindrical waveguide.

# Diffraction anomalies



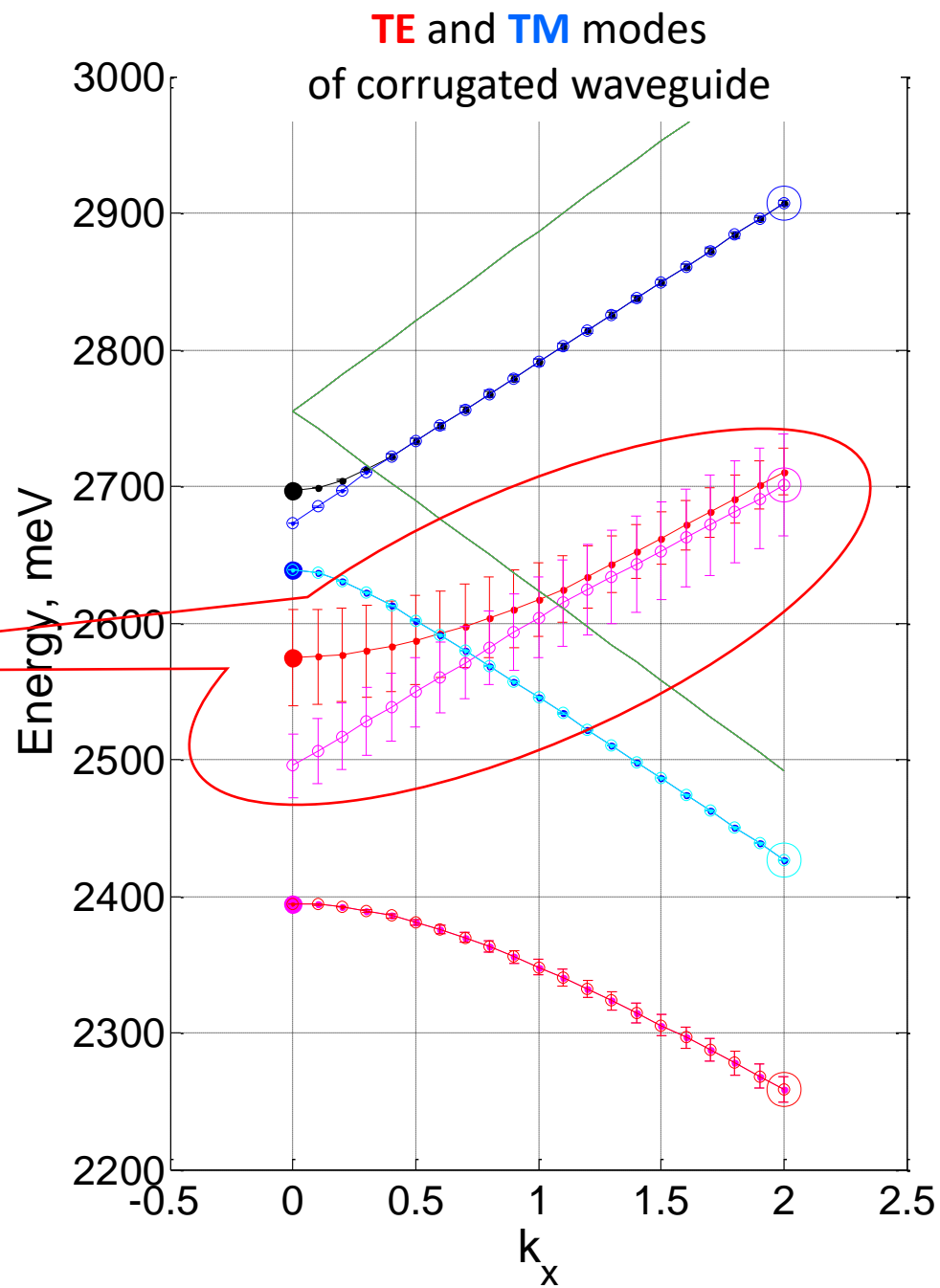
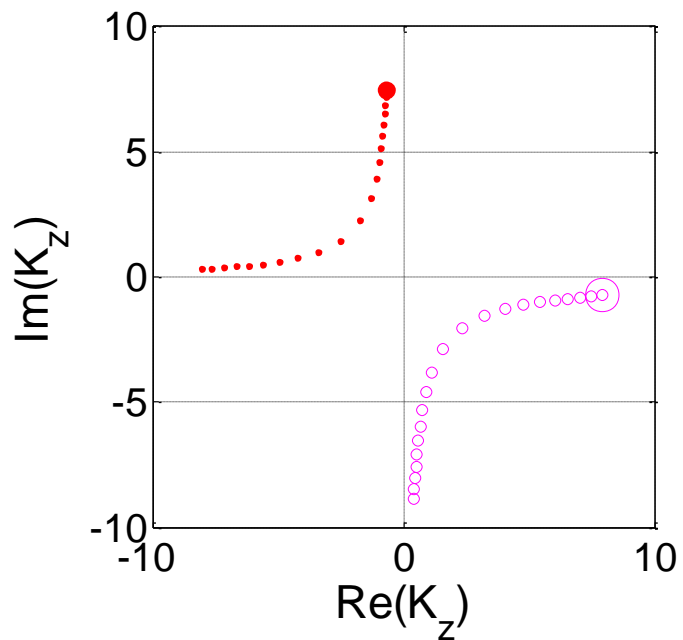
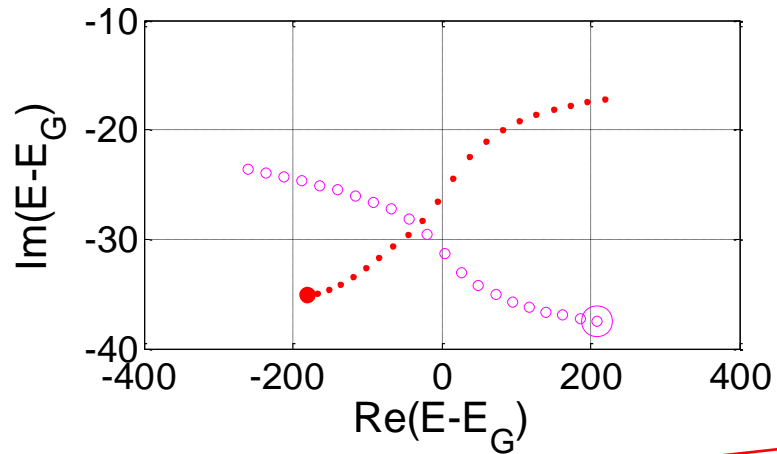
$$k_z(G_x^n) = \sqrt{\frac{\epsilon_i}{c^2} \omega^2 - (k_x + G_x^n)^2}$$

$$k_z(G_x^n) \sim \sqrt{E - E_G}$$

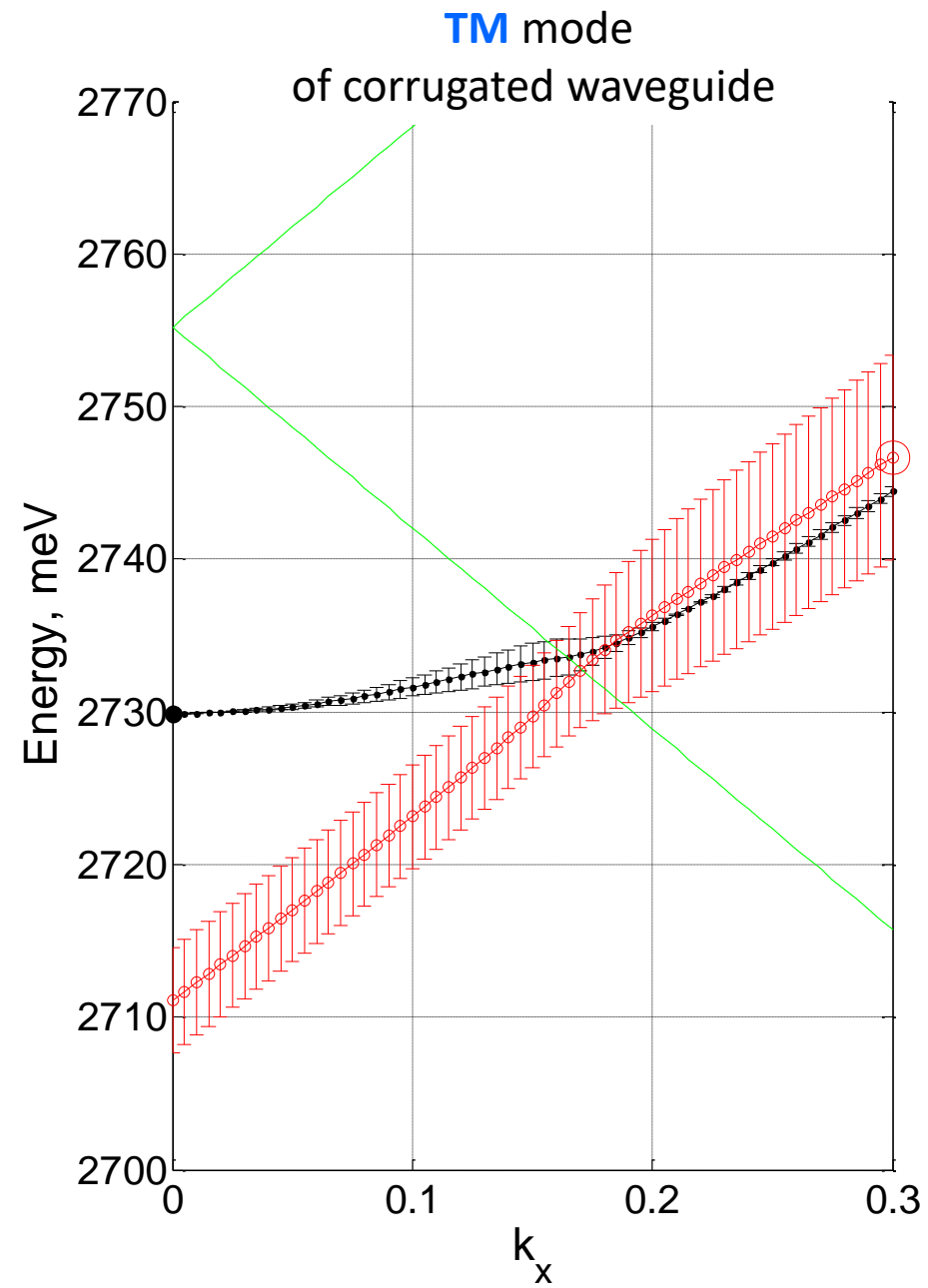
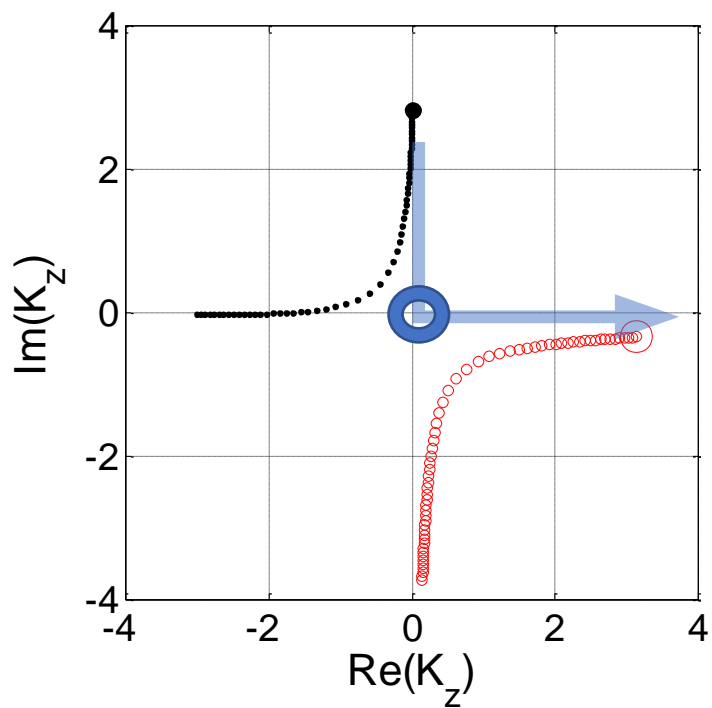
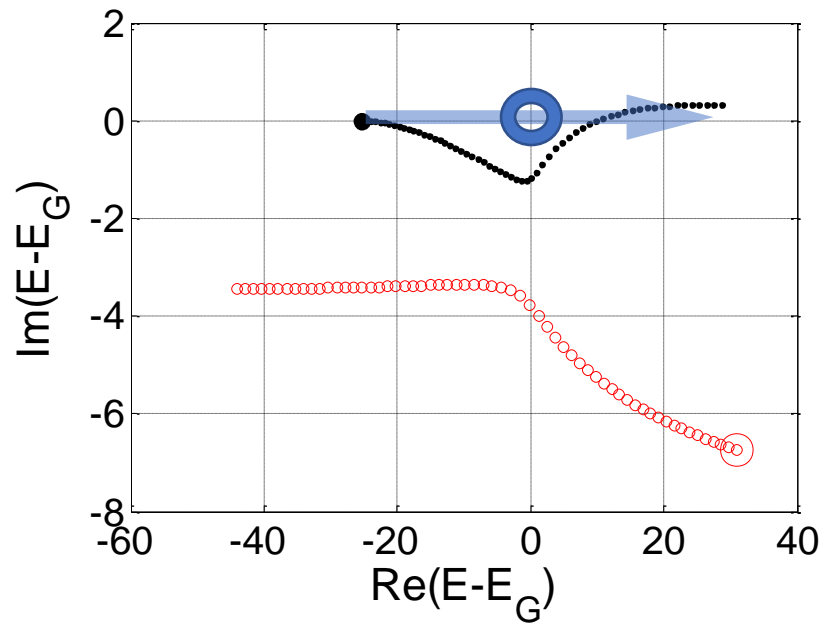


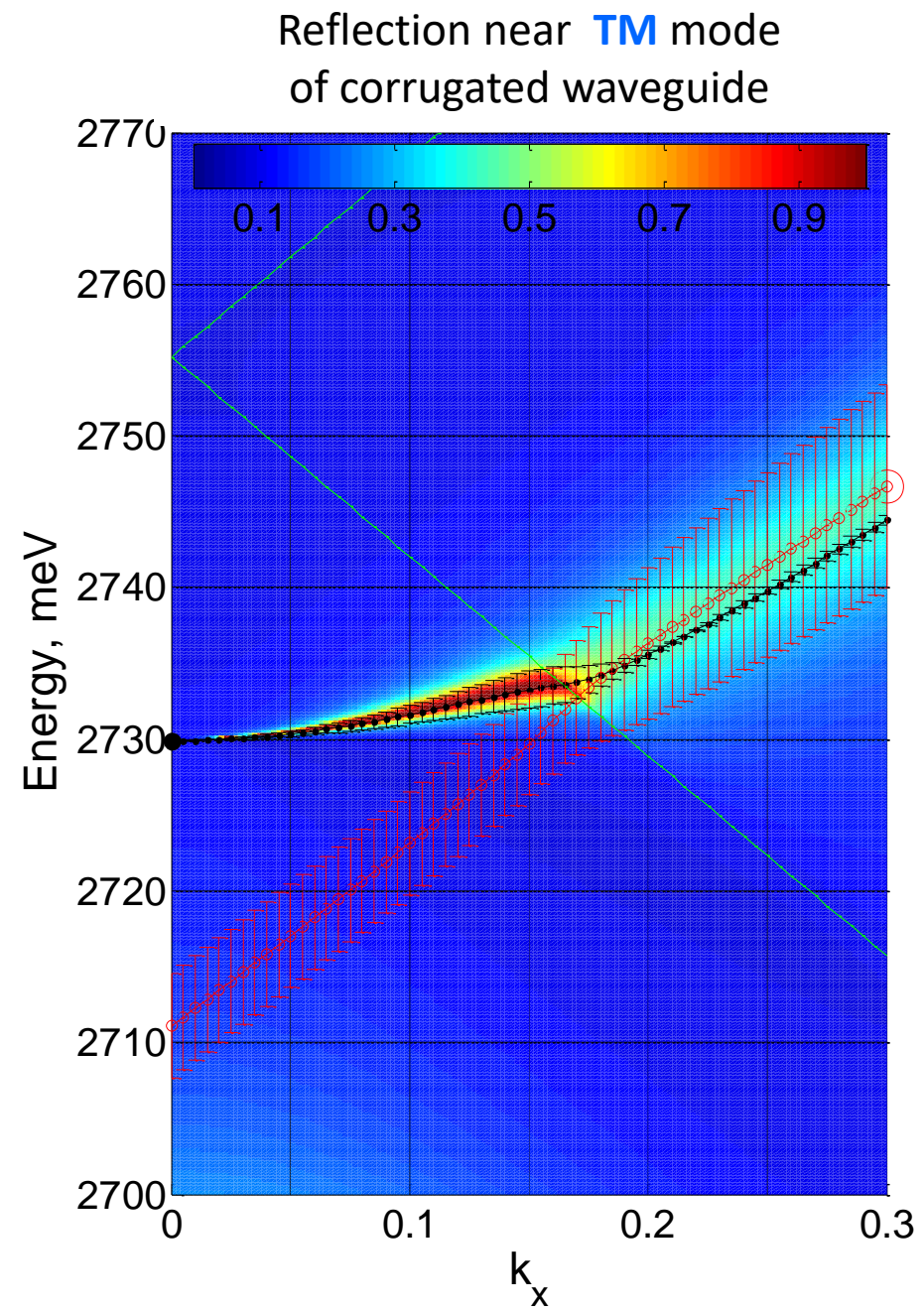
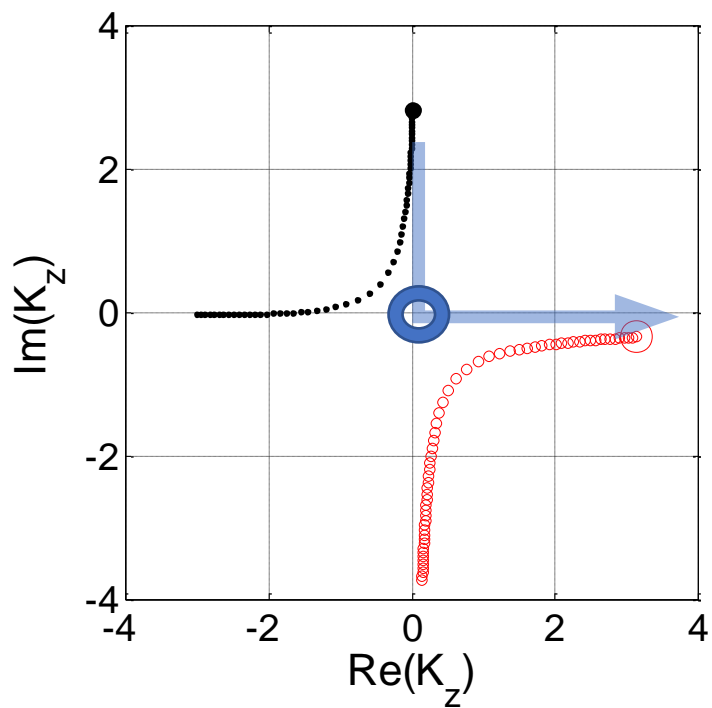
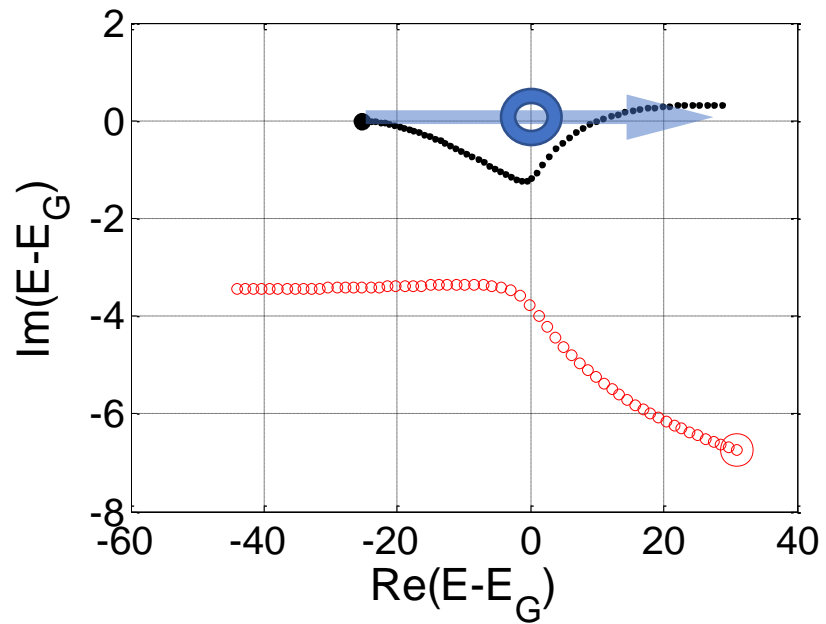
$$k_z(G_x^n) \sim \sqrt{E - E_G}$$

$$E - E_G \sim k_z^2(G)$$









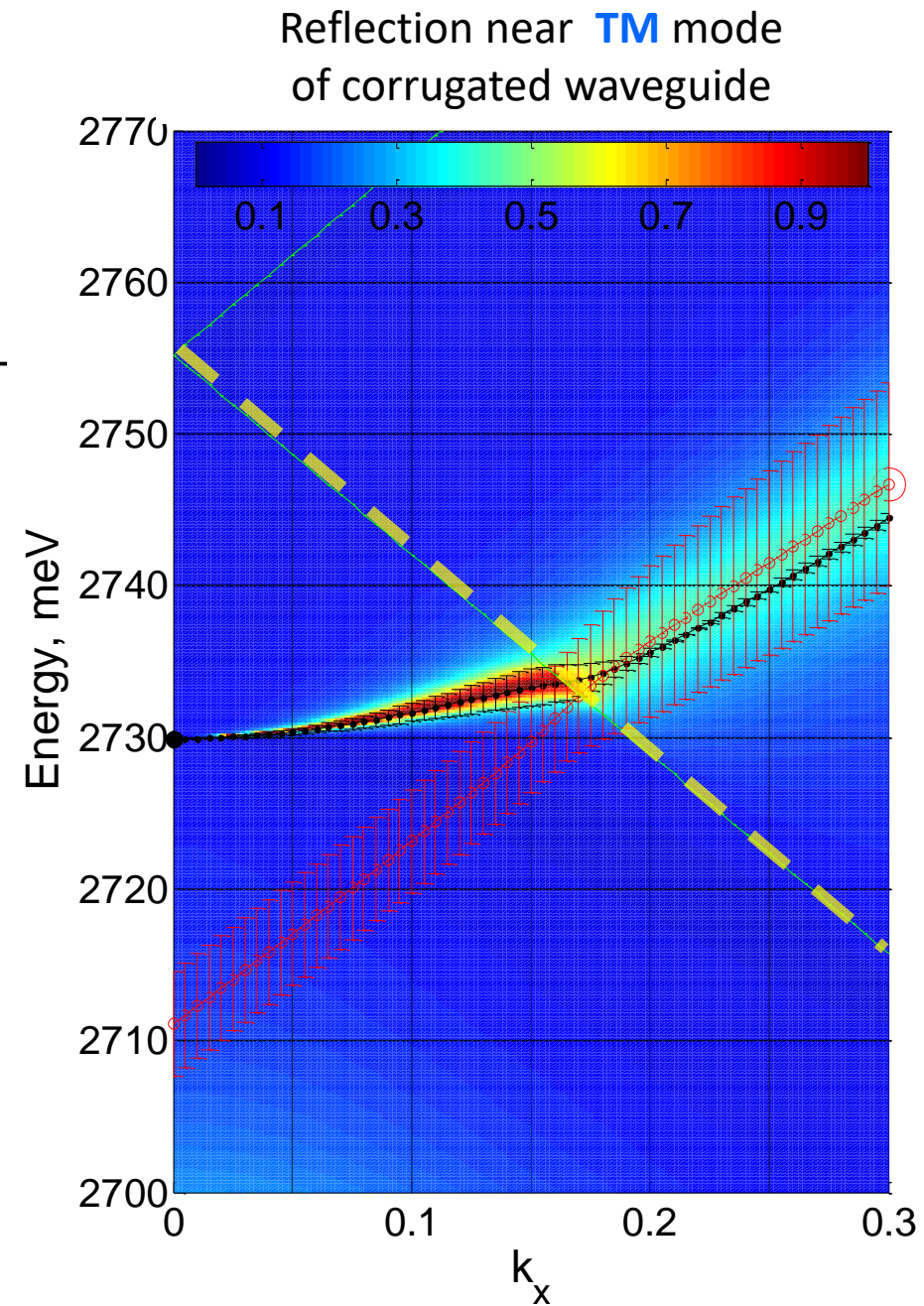
$$E = \hbar\omega$$

$$k_z \left( G_x^n \right) = \sqrt{\frac{\epsilon_i}{c^2} W^2 - \left( k_x + G_x^n \right)^2}$$

$$k_z^2 \propto \left( E - E_G \right)$$

$$E = E_G + Ak_z^2$$

Akimov AB, Gippius NA, Tikhodeev SG.  
JETP Letters 93(8):427 (2011)

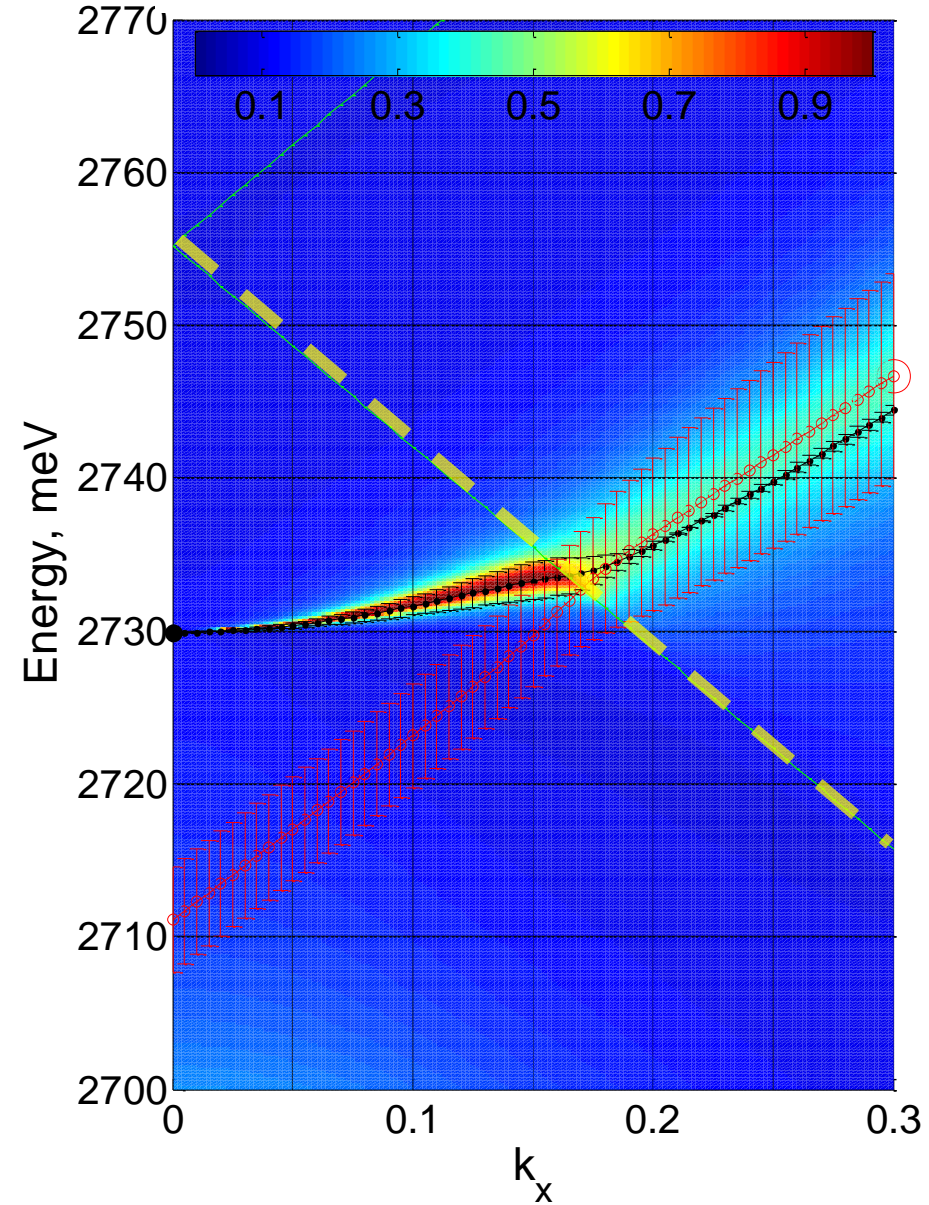




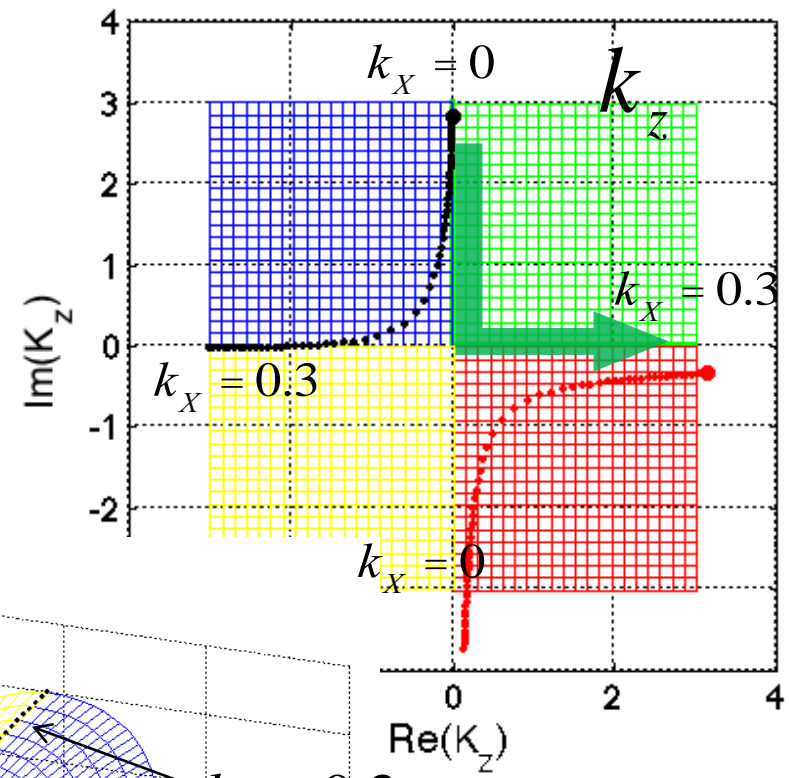
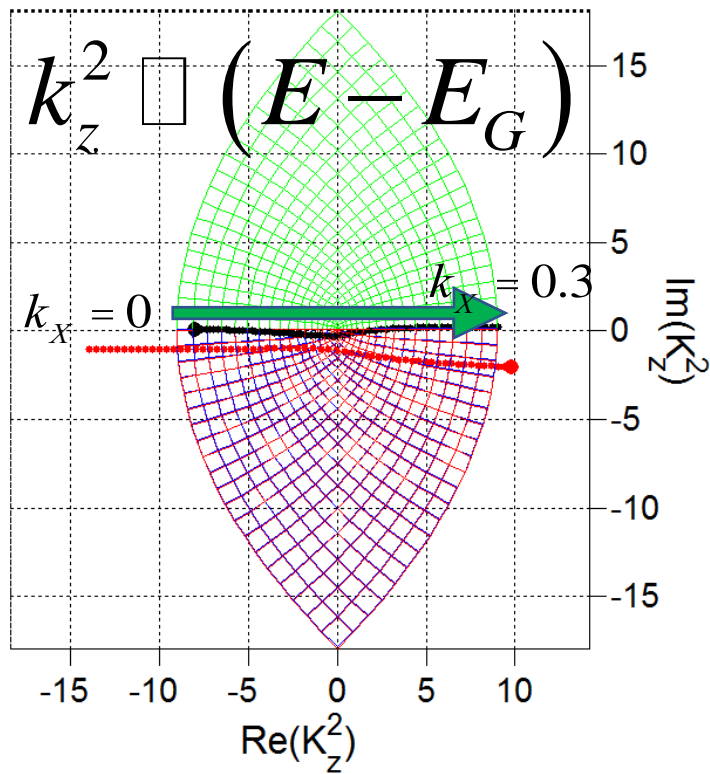
$$\tilde{S} + \sum_r |O_r\rangle \frac{1}{E - E_r} \langle I_r |$$

$$\tilde{S} + \sum_r |o_r\rangle \frac{1}{k_z - k_z^r} \langle i_r |$$

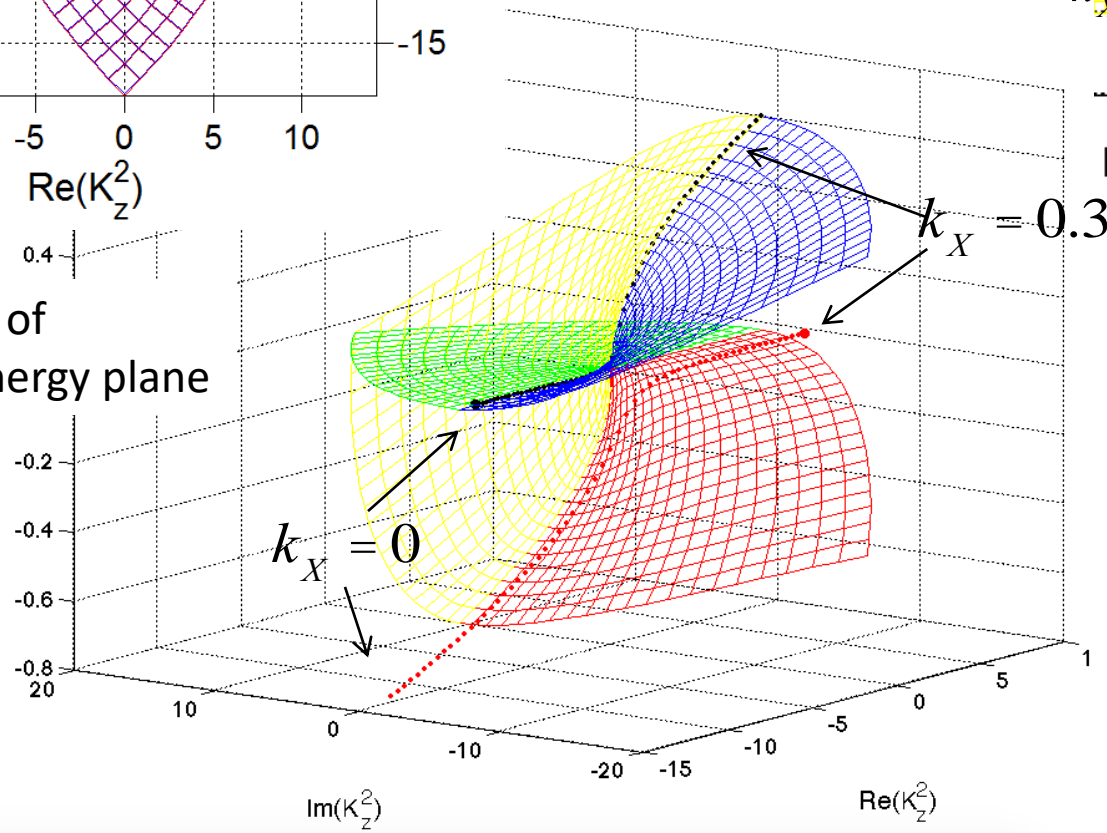
Reflection near **TM** mode  
of corrugated waveguide



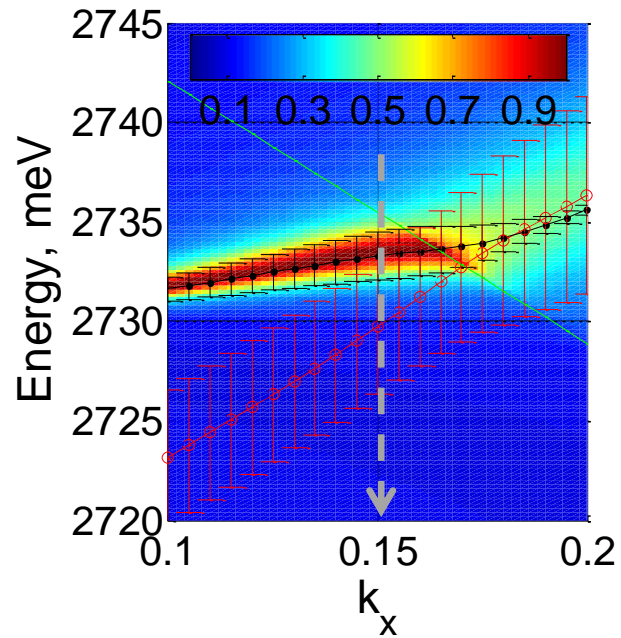
Akimov AB, Gippius NA, Tikhodeev SG.  
JETP Letters 93(8):427 (2011)



Two leaves of complex energy plane

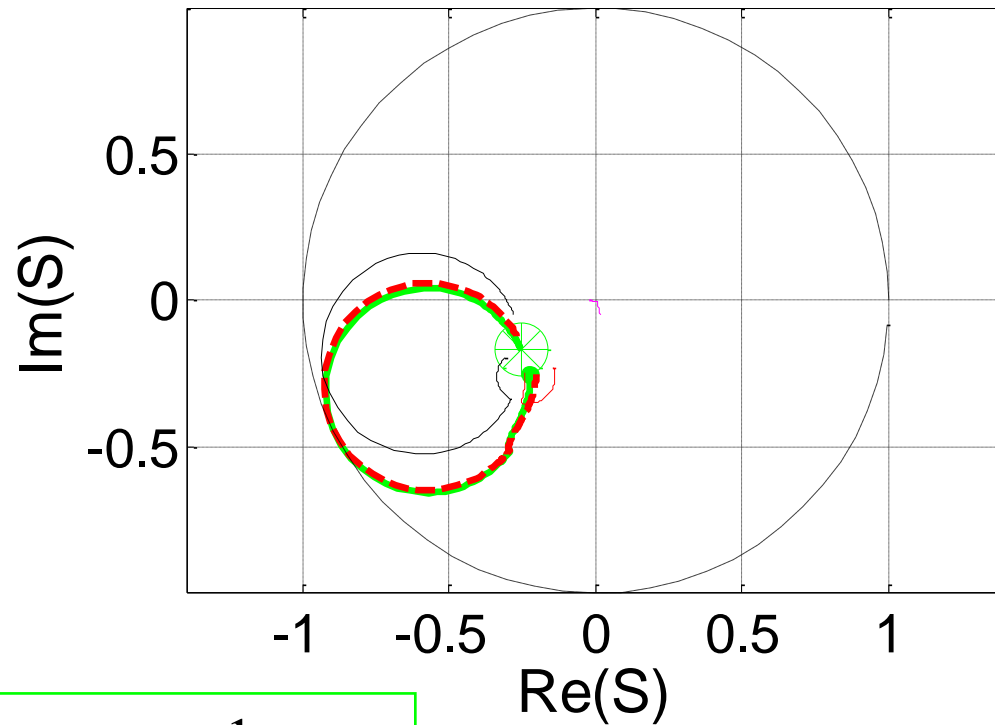
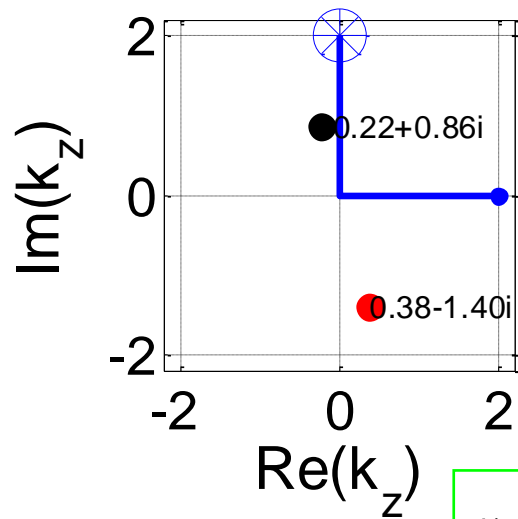




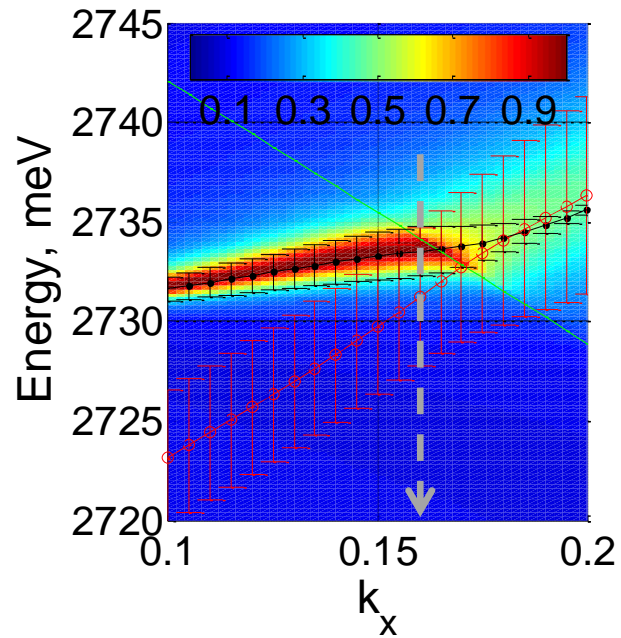


Reflection of corrugated waveguide (green line) and the approximations with two poles (red dashed line) and one pole (thin lines)

$$k_x = 0.15$$

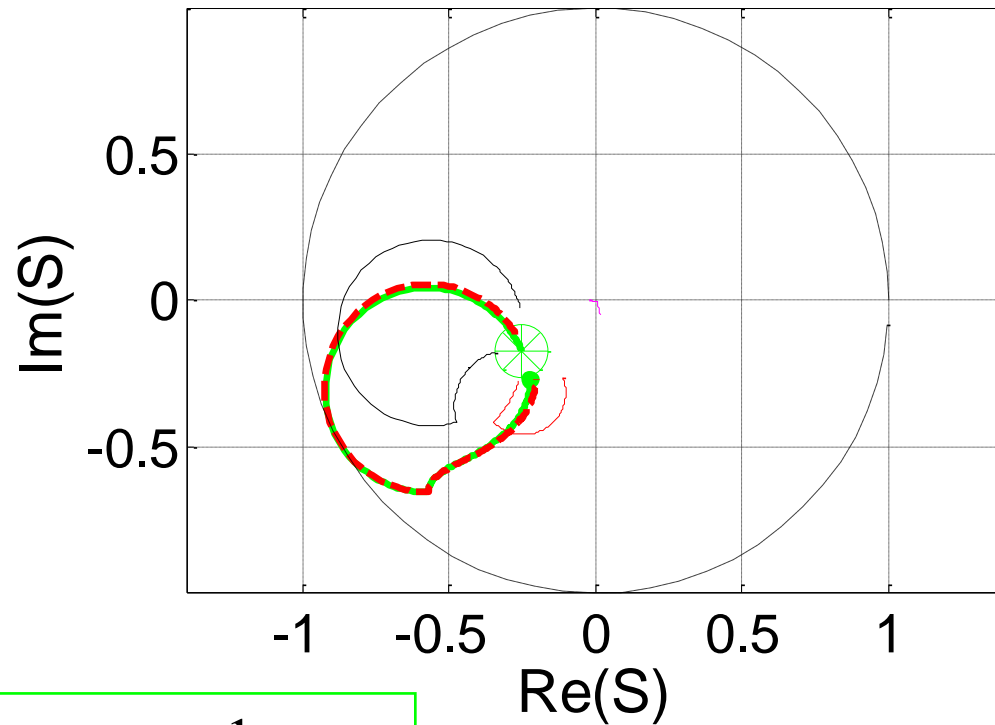
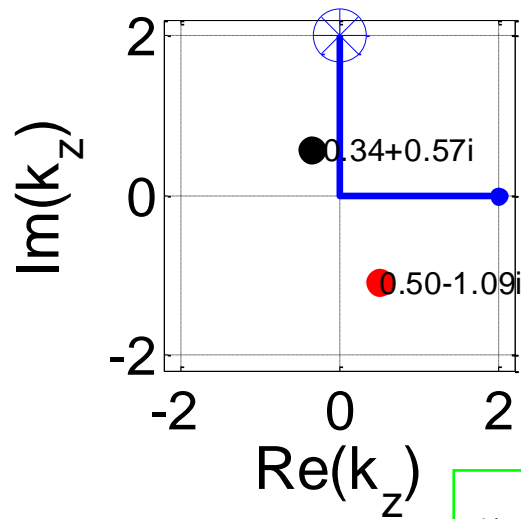


$$\tilde{S} + \sum_r |o_r\rangle \frac{1}{k_z - k_z^r} \langle i_r|$$

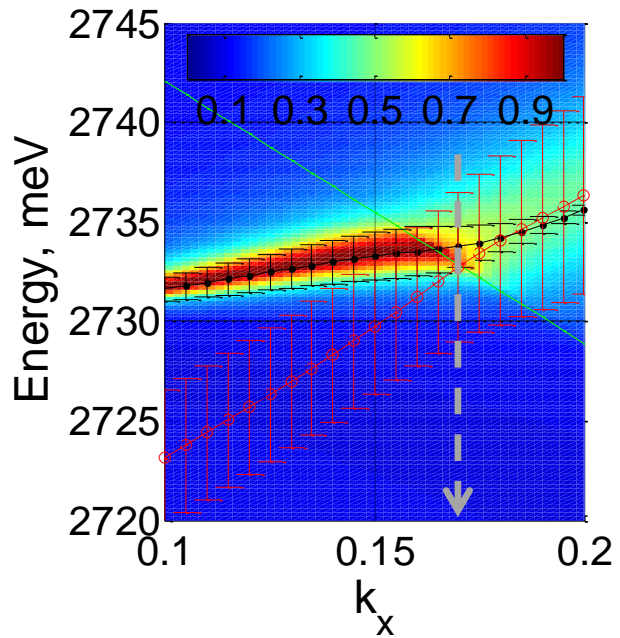


Reflection of corrugated waveguide (green line) and the approximations with two poles (red dashed line) and one pole (thin lines)

$$k_x = 0.16$$

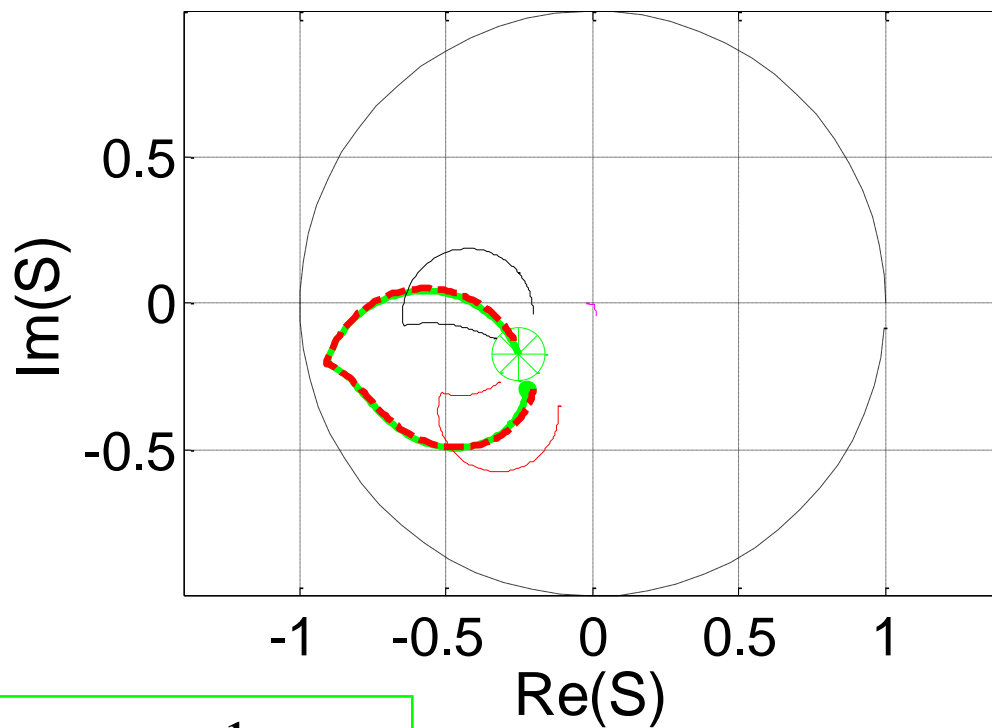
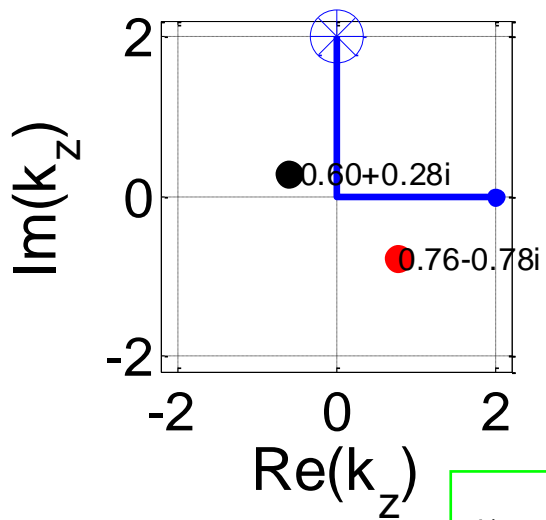


$$\tilde{S} + \sum_r |o_r\rangle \frac{1}{k_z - k_z^r} \langle i_r|$$

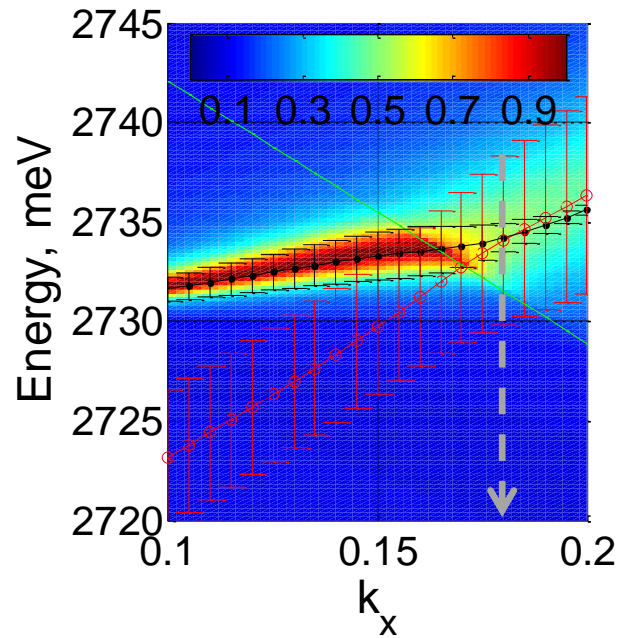


Reflection of corrugated waveguide (green line) and the approximations with two poles (red dashed line) and one pole (thin lines)

$$k_x = 0.17$$

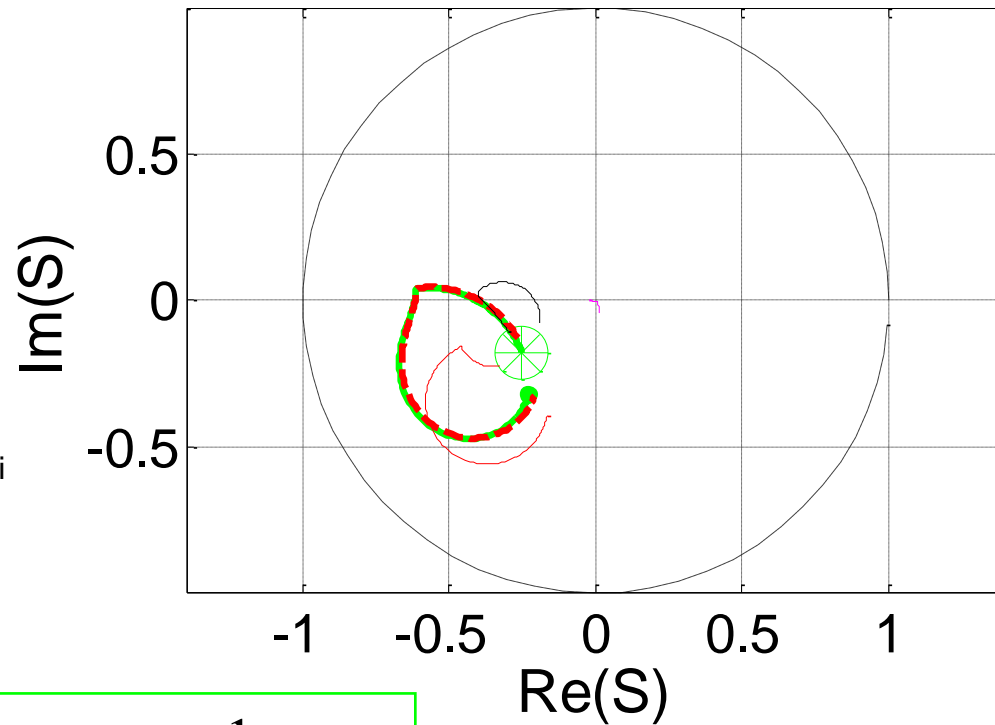
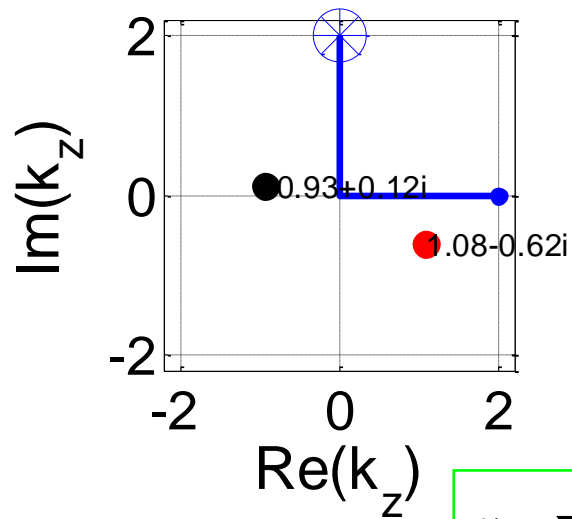


$$\tilde{S} + \sum_r |o_r\rangle \frac{1}{k_z - k_z^r} \langle i_r|$$

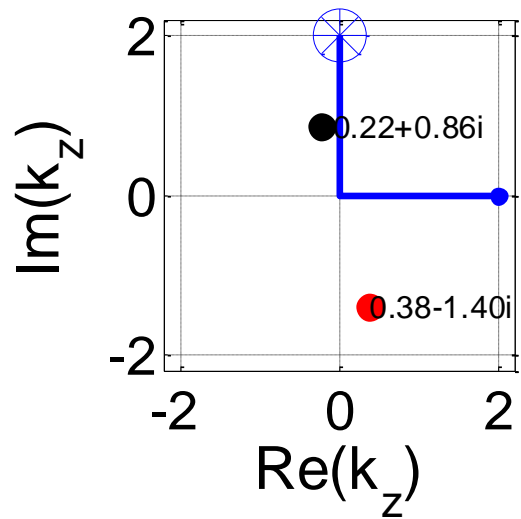
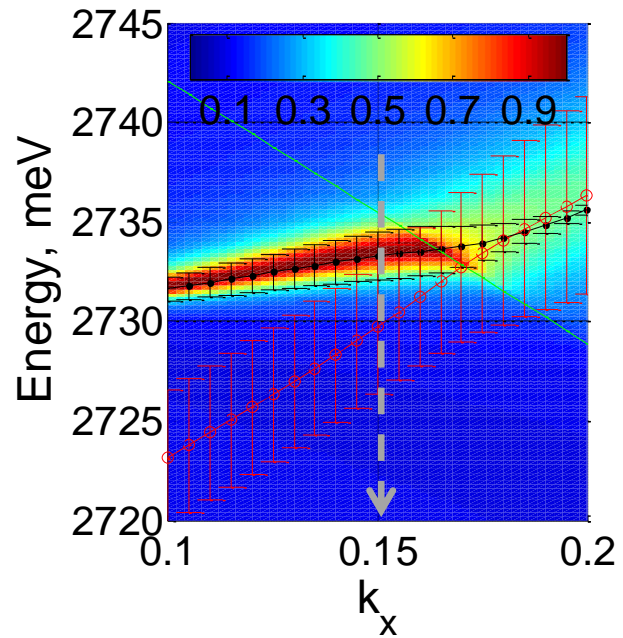


Reflection of corrugated waveguide (green line) and the approximations with two poles (red dashed line) and one pole (thin lines)

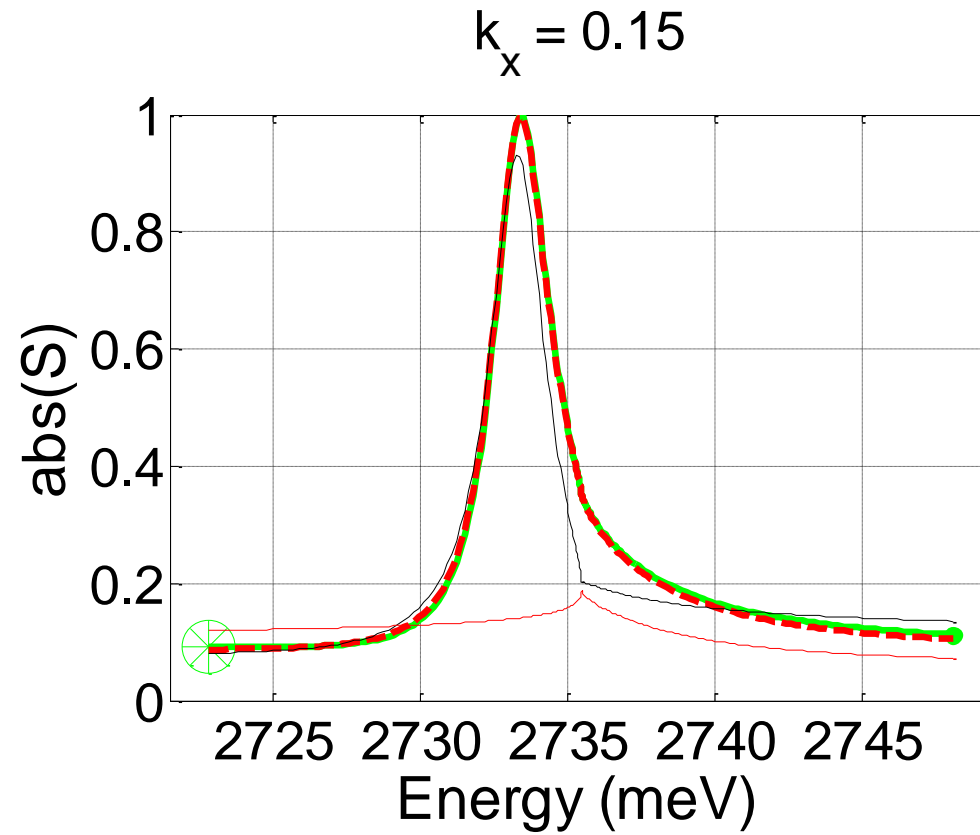
$$k_x = 0.18$$



$$\tilde{S} + \sum_r |o_r\rangle \frac{1}{k_z - k_z^r} \langle i_r|$$

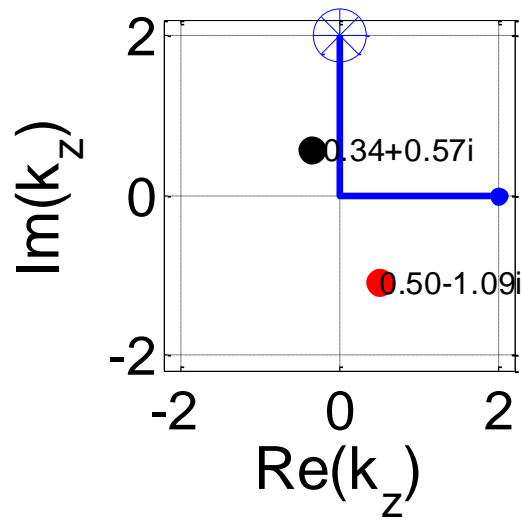
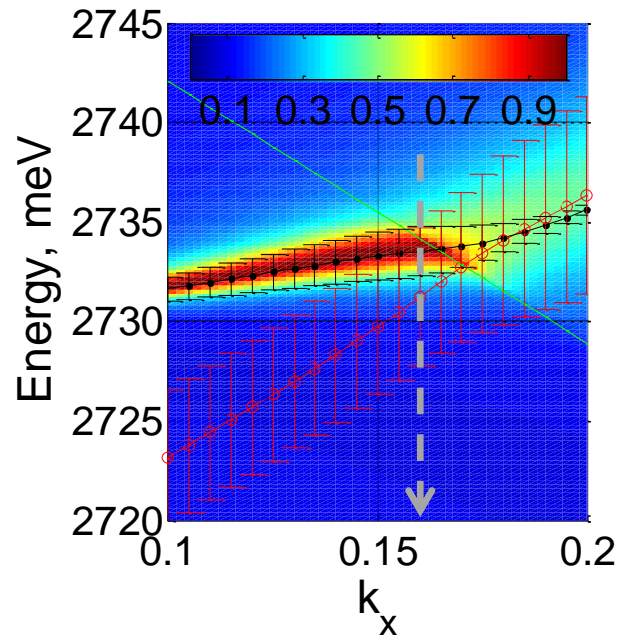


Reflection of corrugated waveguide (green line) and the approximations with two poles (red dashed line) and one pole (thin lines)

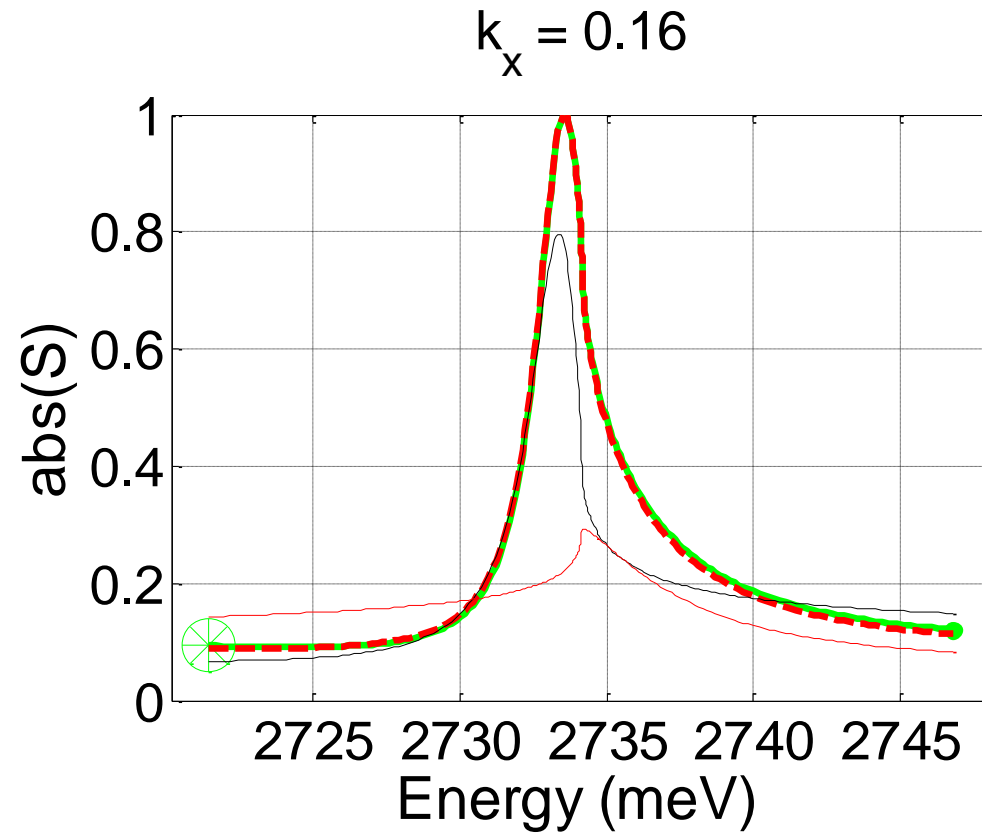


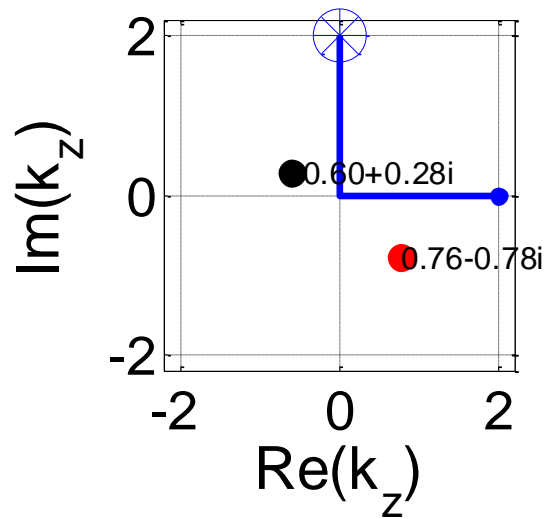
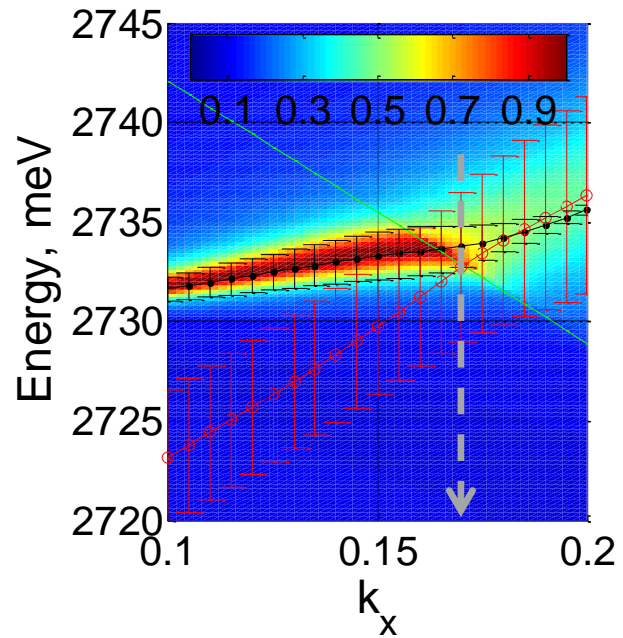
$$E - E_G(k_x) \sim k_z^2(G)$$





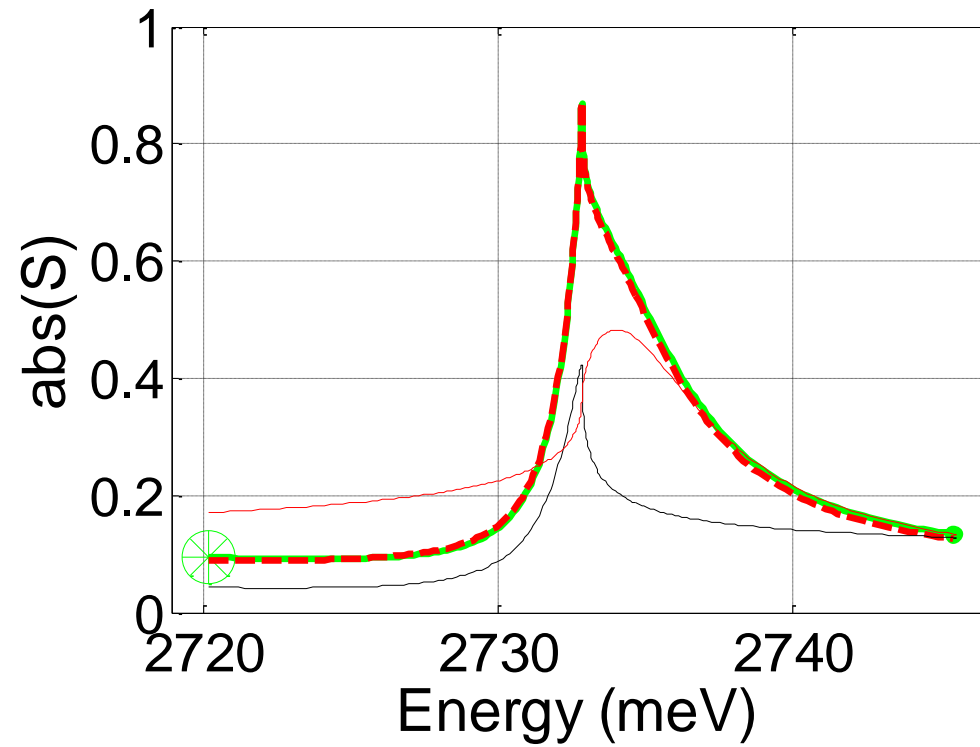
Reflection of corrugated waveguide (green line) and the approximations with two poles (red dashed line) and one pole (thin lines)

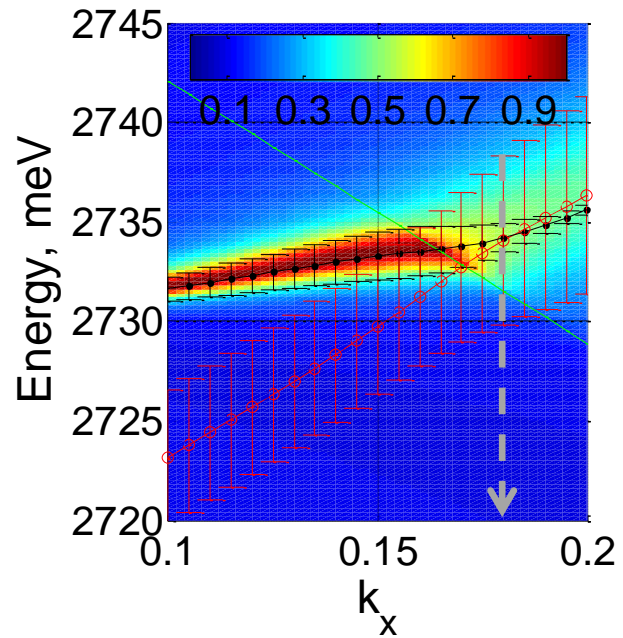




Reflection of corrugated waveguide (green line) and the approximations with two poles (red dashed line) and one pole (thin lines)

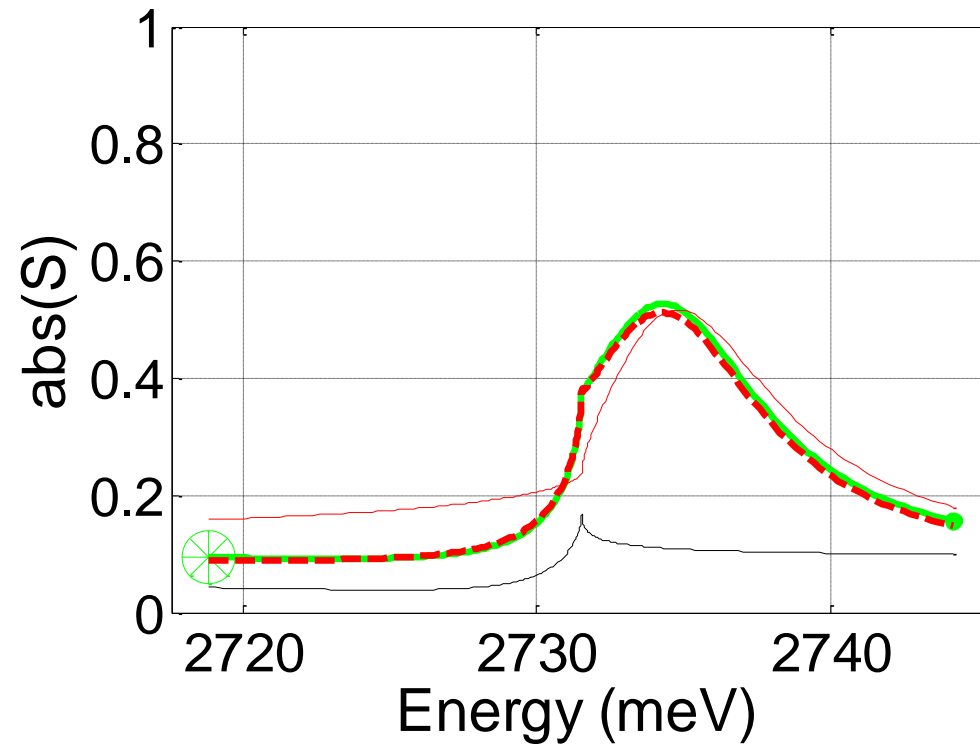
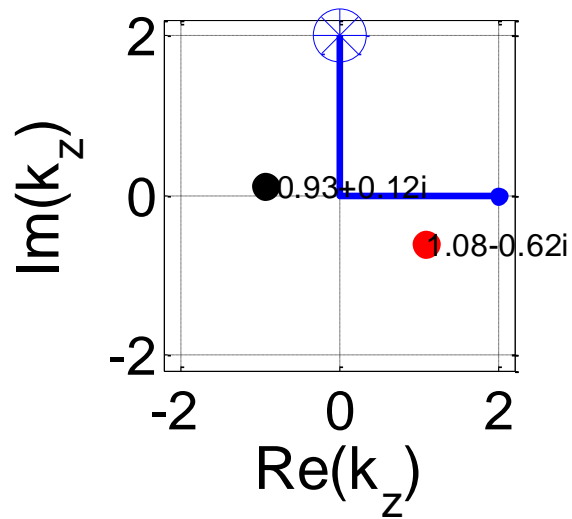
$$k_x = 0.17$$

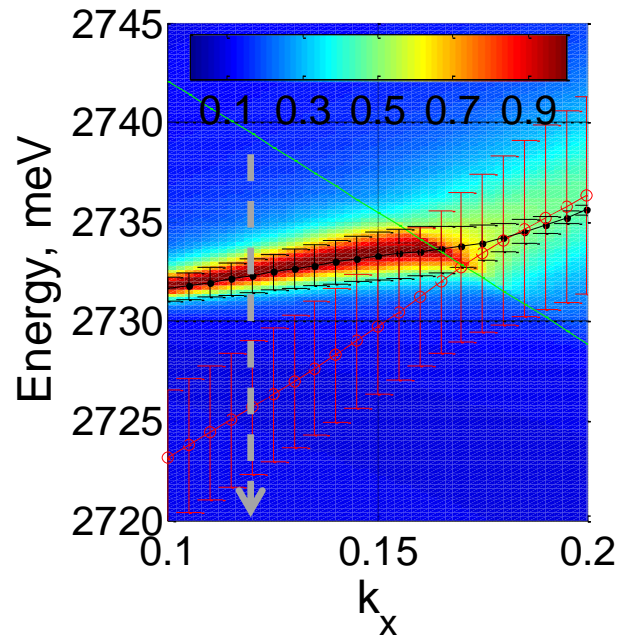




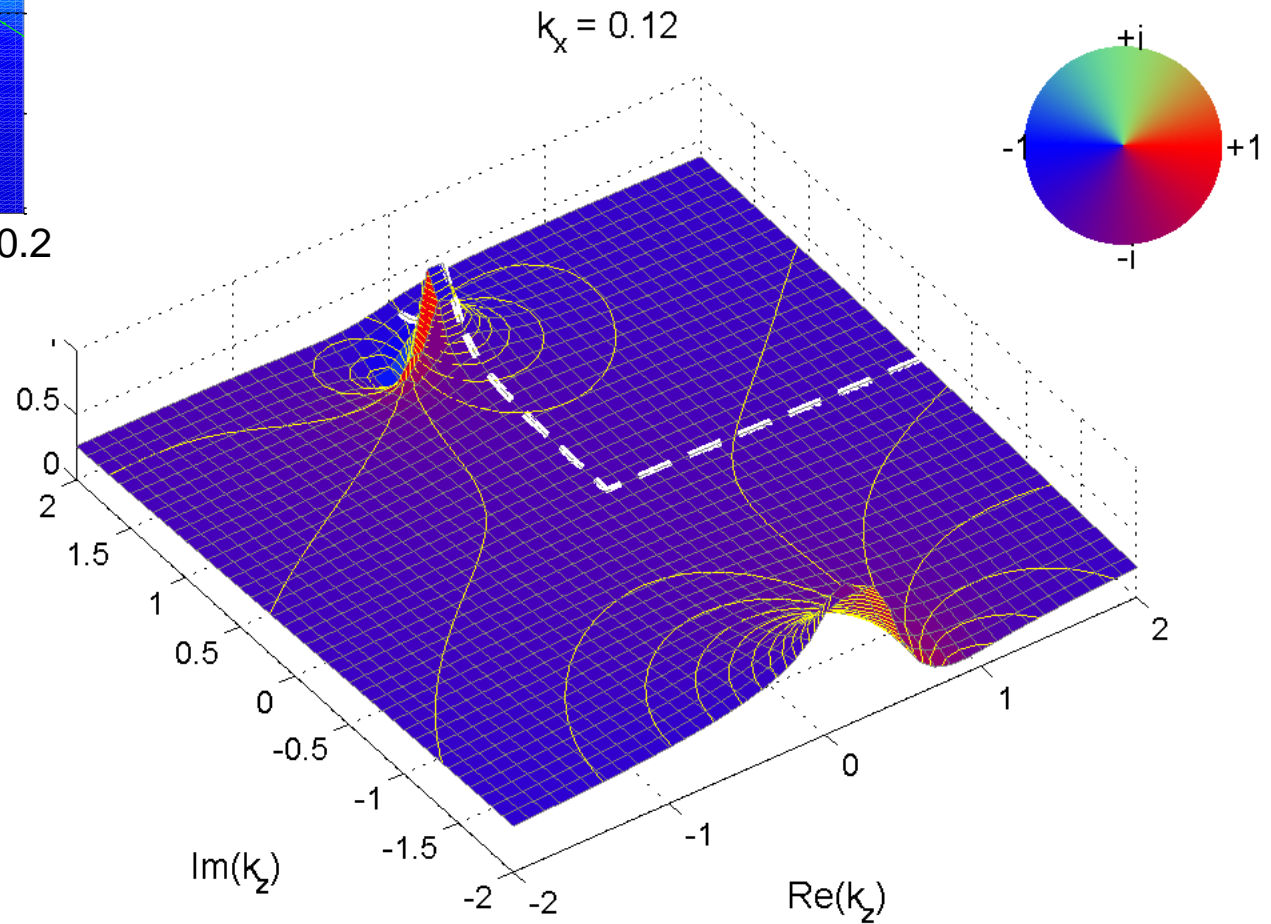
Reflection of corrugated waveguide (green line) and the approximations with two poles (red dashed line) and one pole (thin lines)

$$k_x = 0.18$$

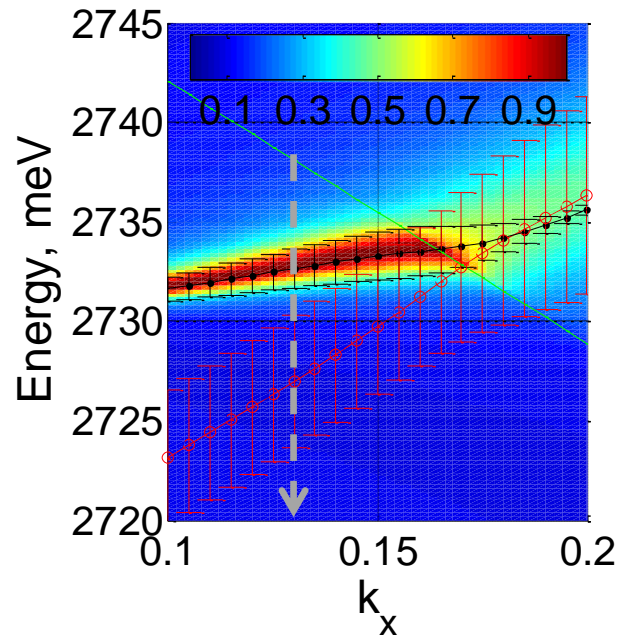




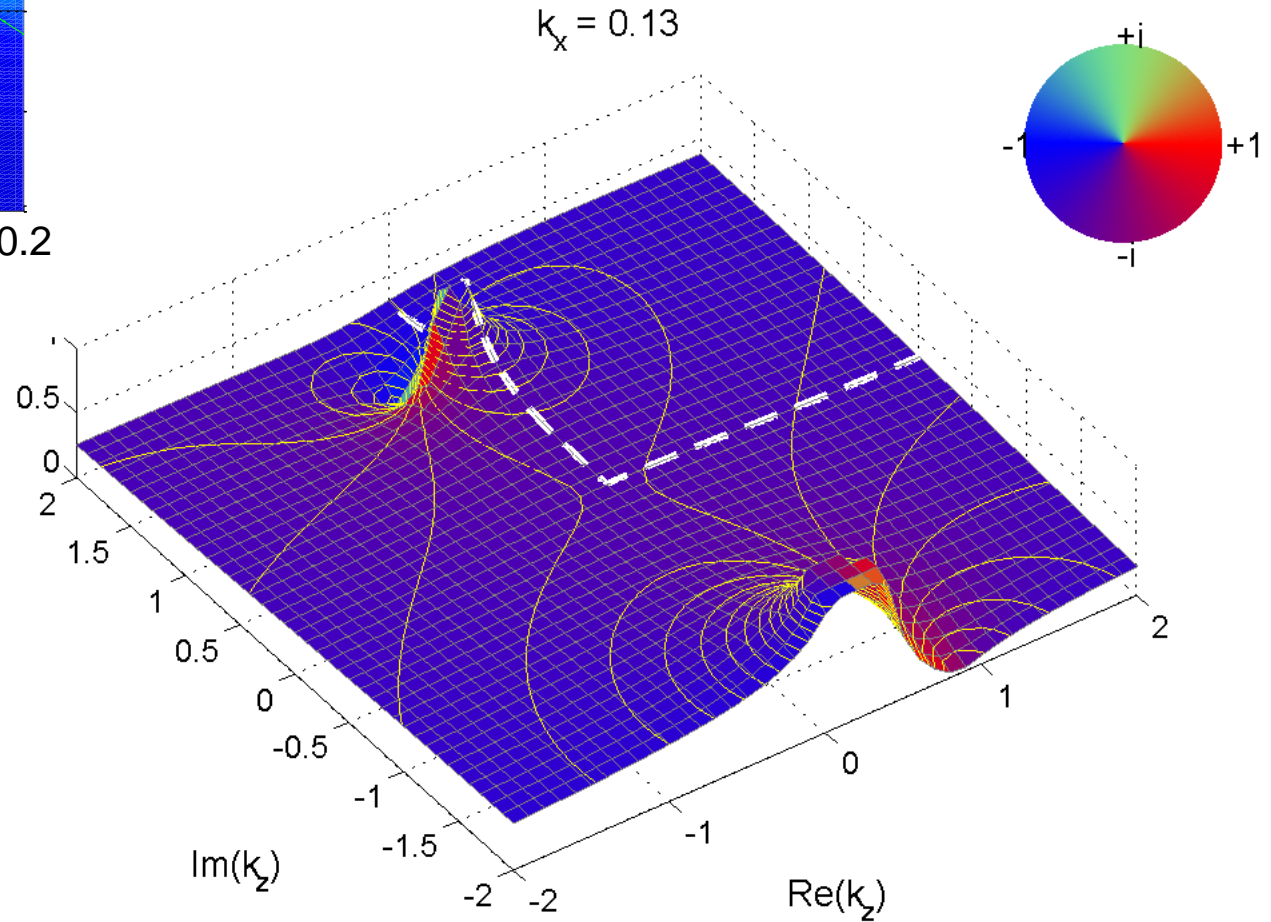
Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction



**Resonantly enhanced Wood-Rayleigh anomaly**

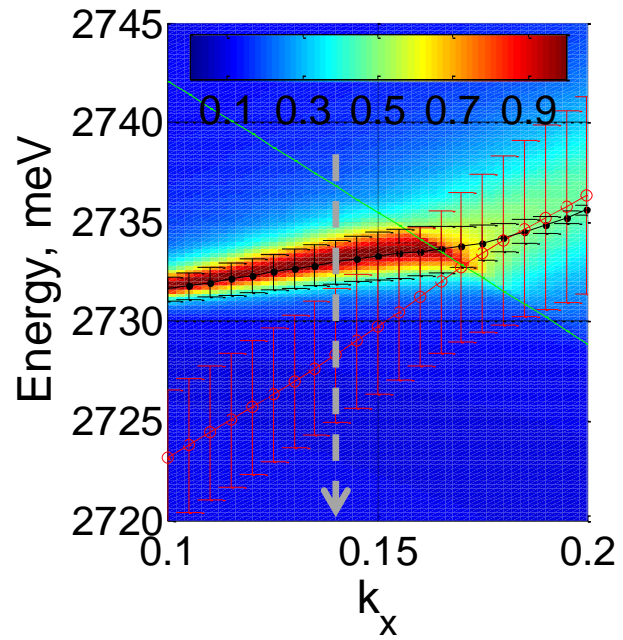


Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction

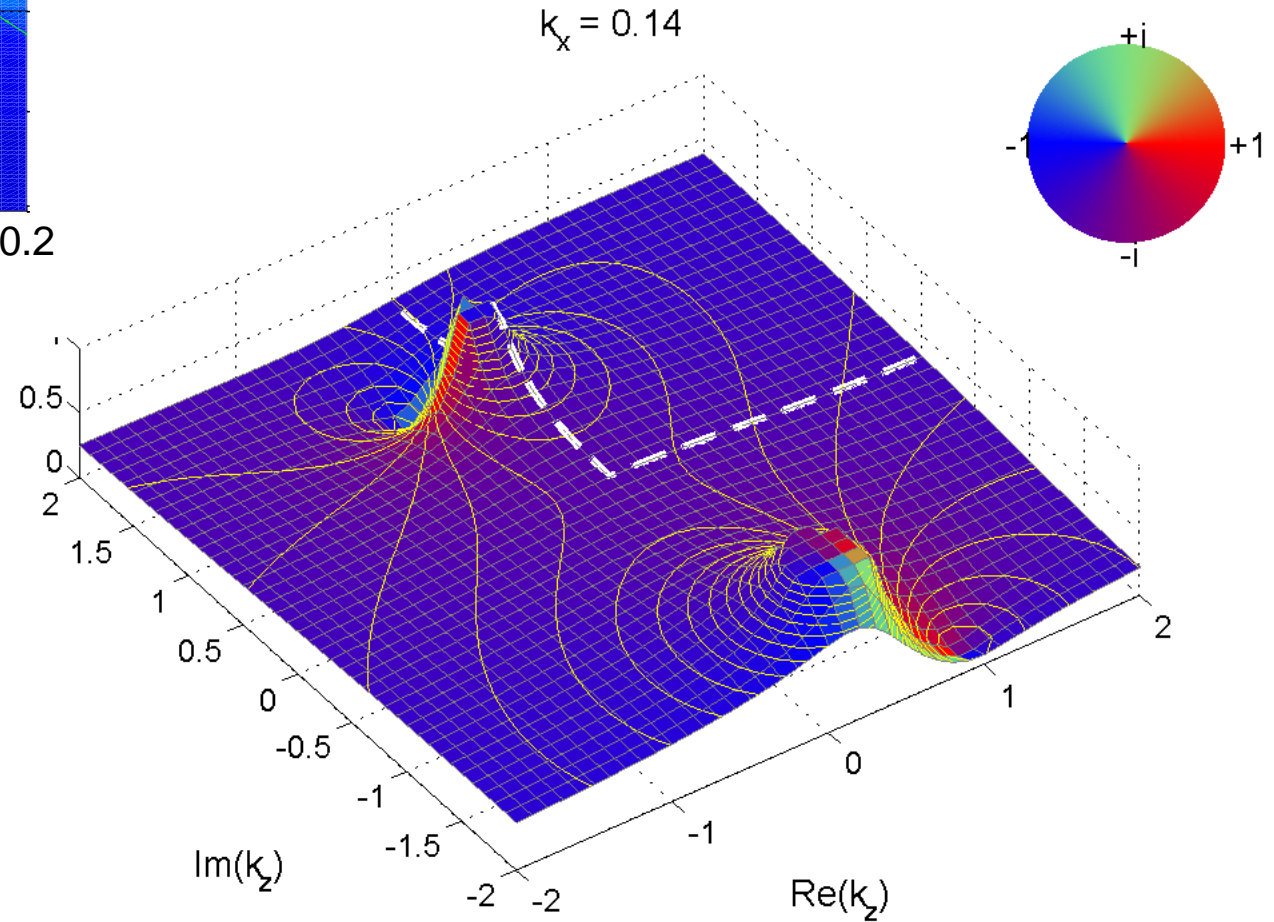


**Resonantly enhanced Wood-Rayleigh anomaly**

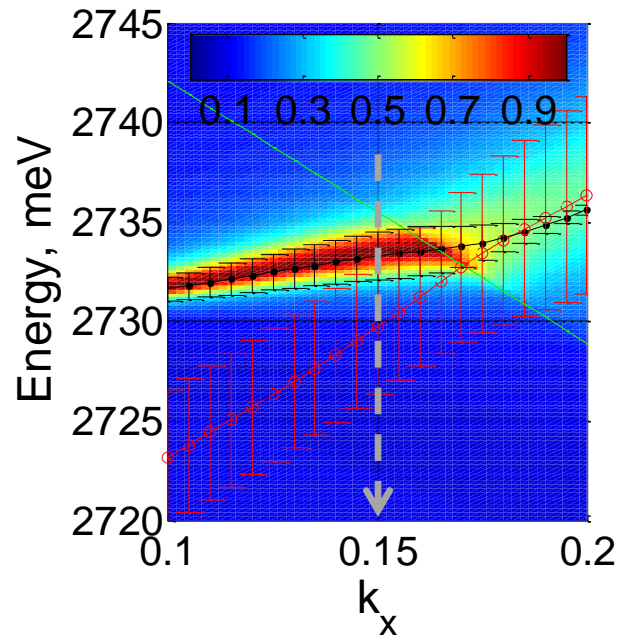




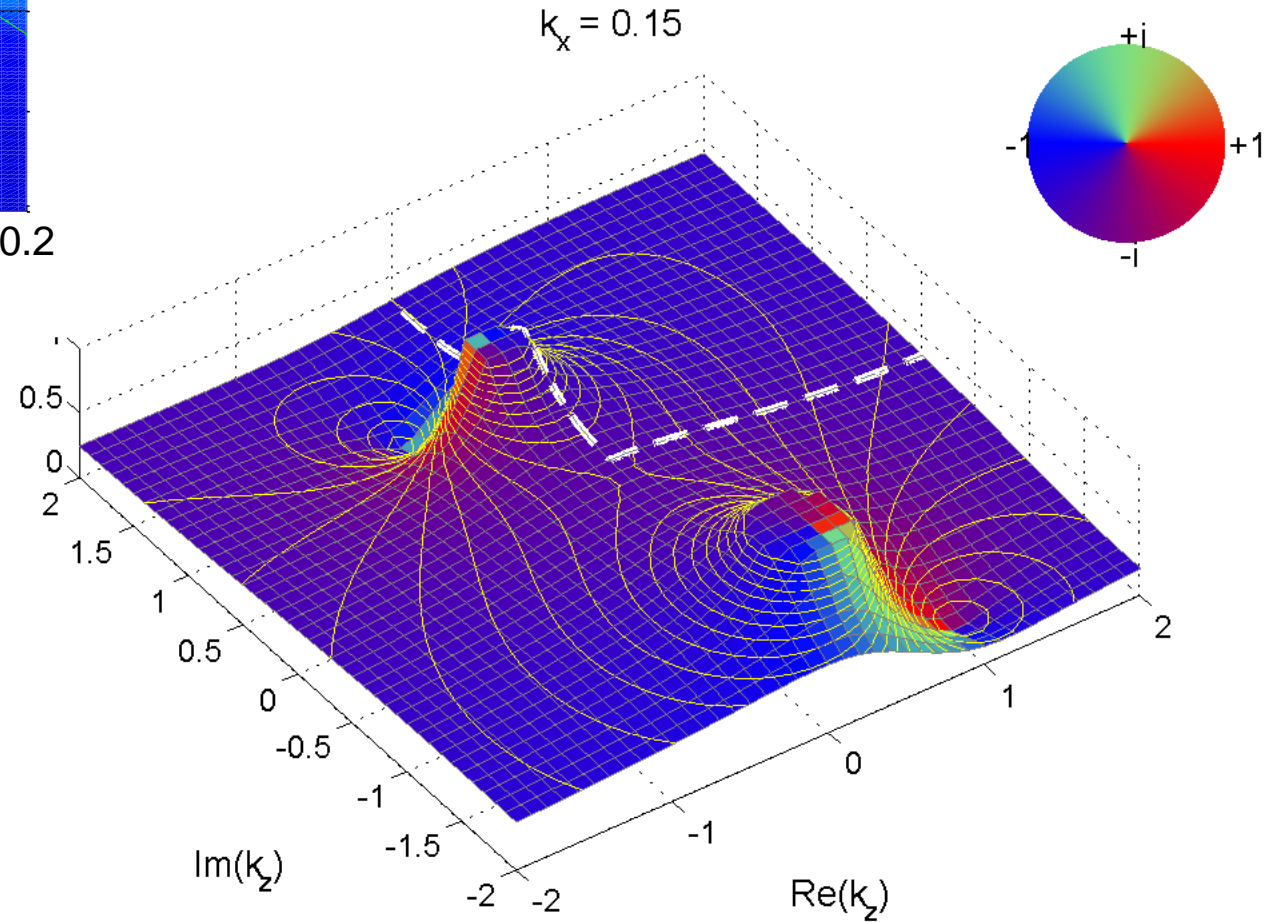
Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction



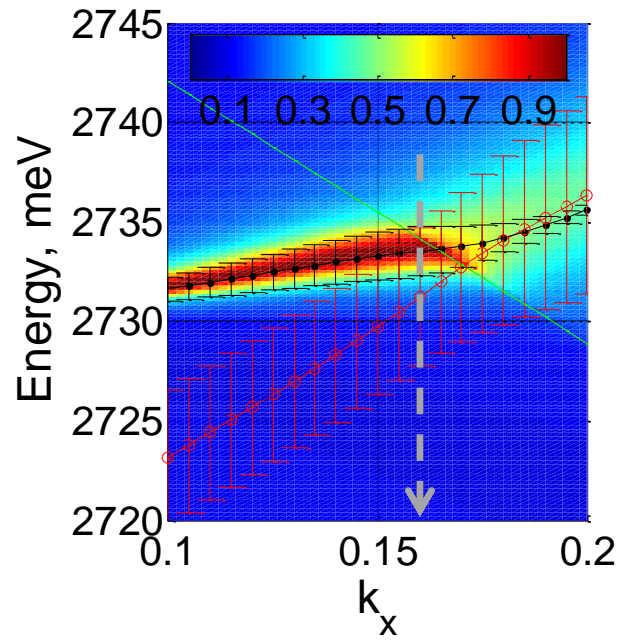
**Resonantly enhanced Wood-Rayleigh anomaly**



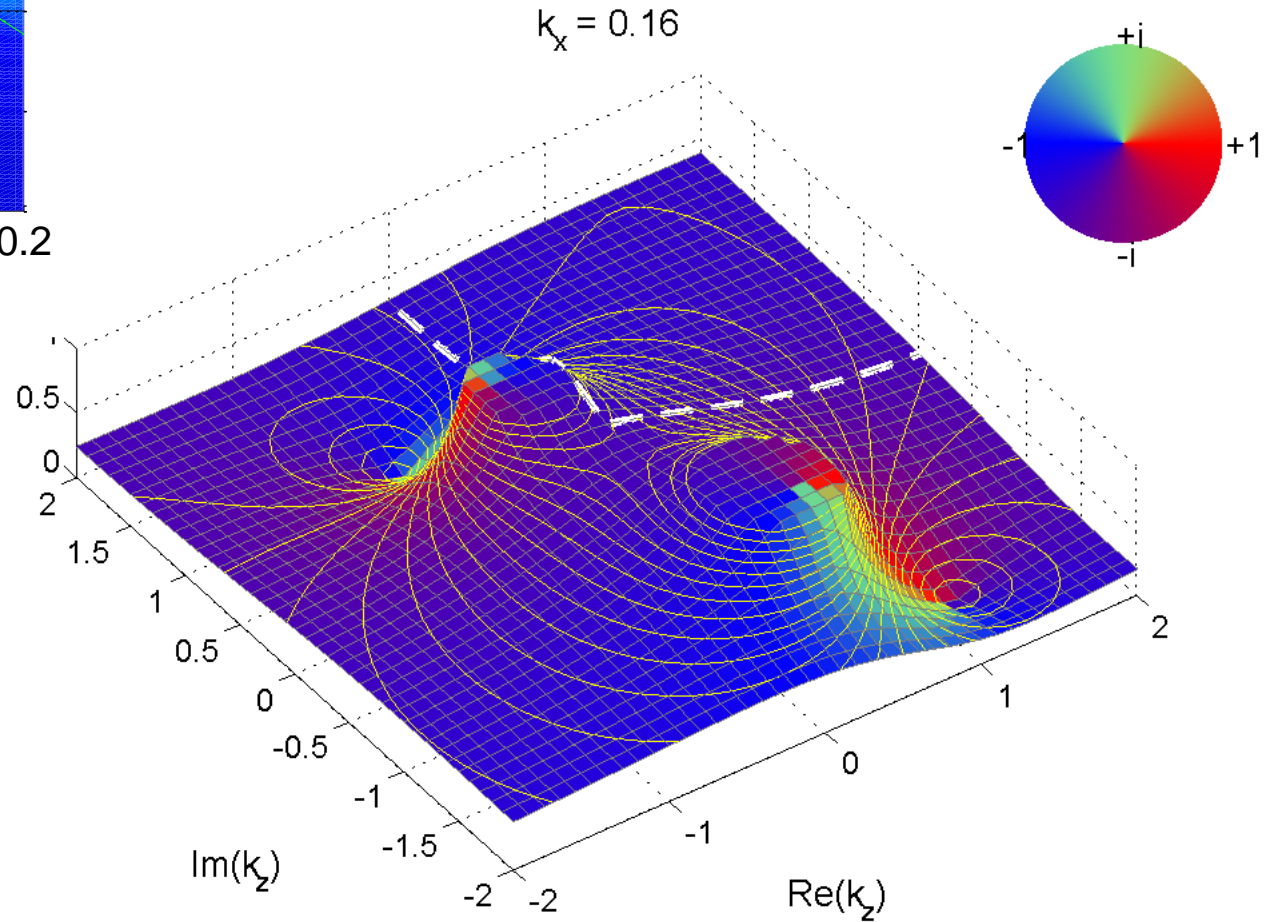
Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction



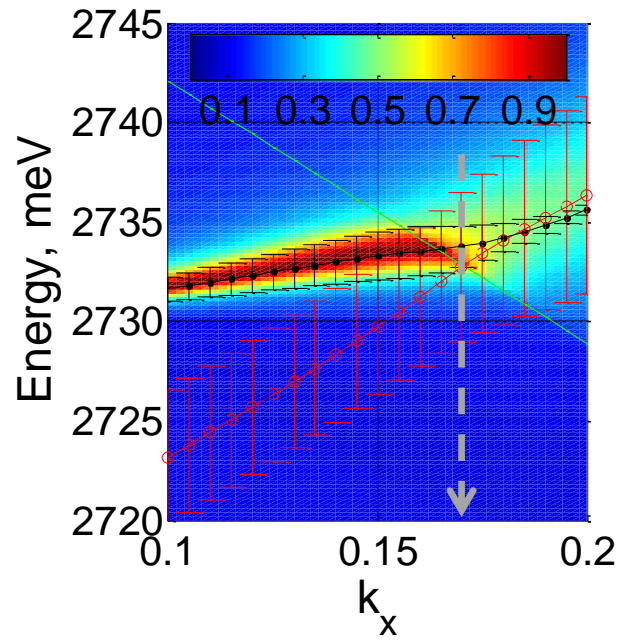
**Resonantly enhanced Wood-Rayleigh anomaly**



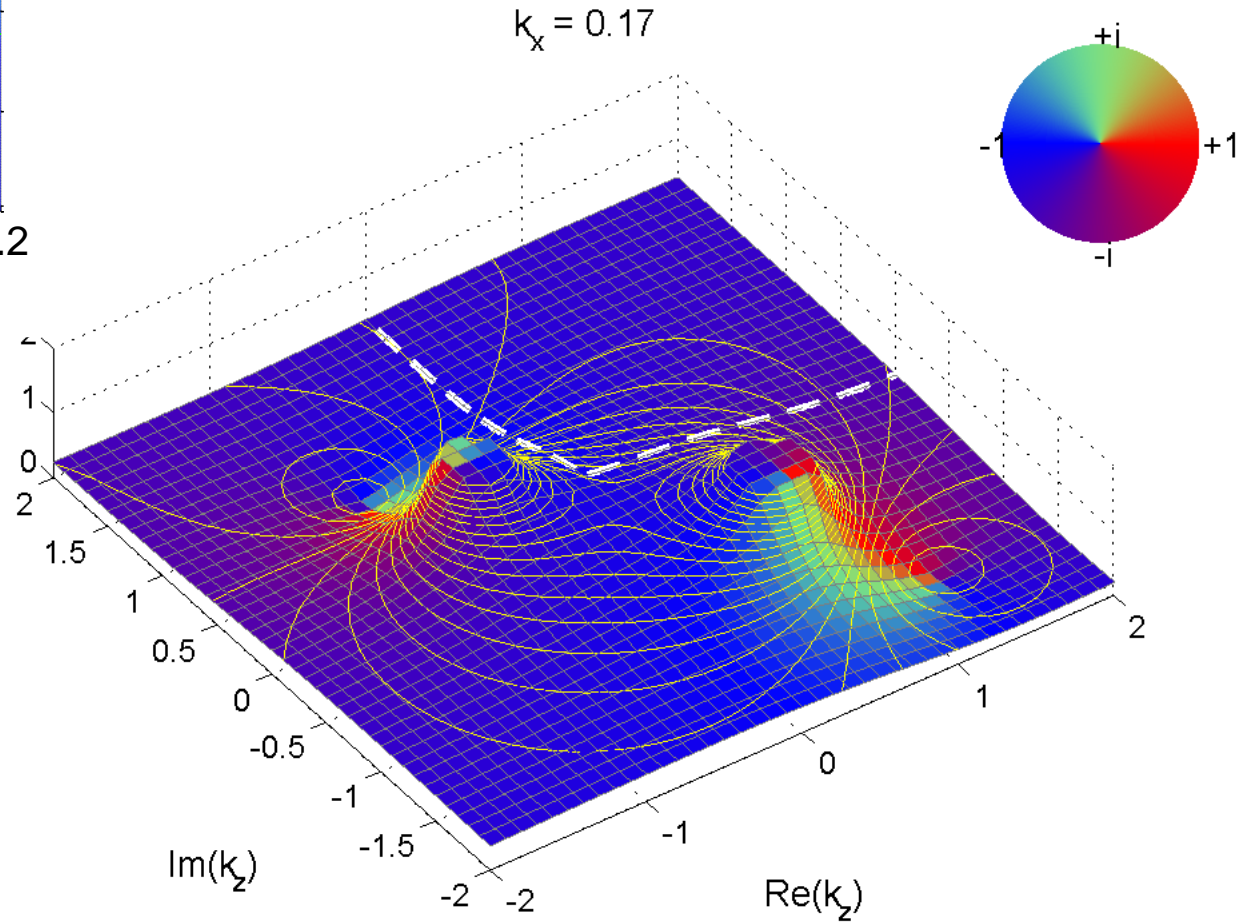
Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction



**Resonantly enhanced Wood-Rayleigh anomaly**

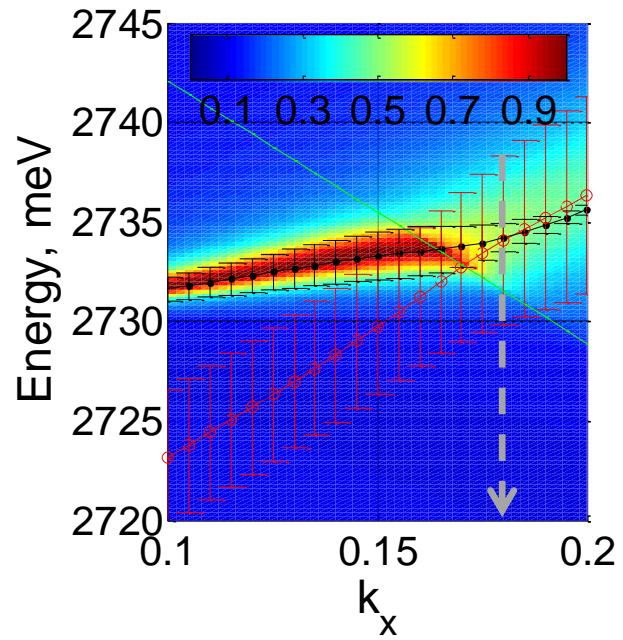


Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction

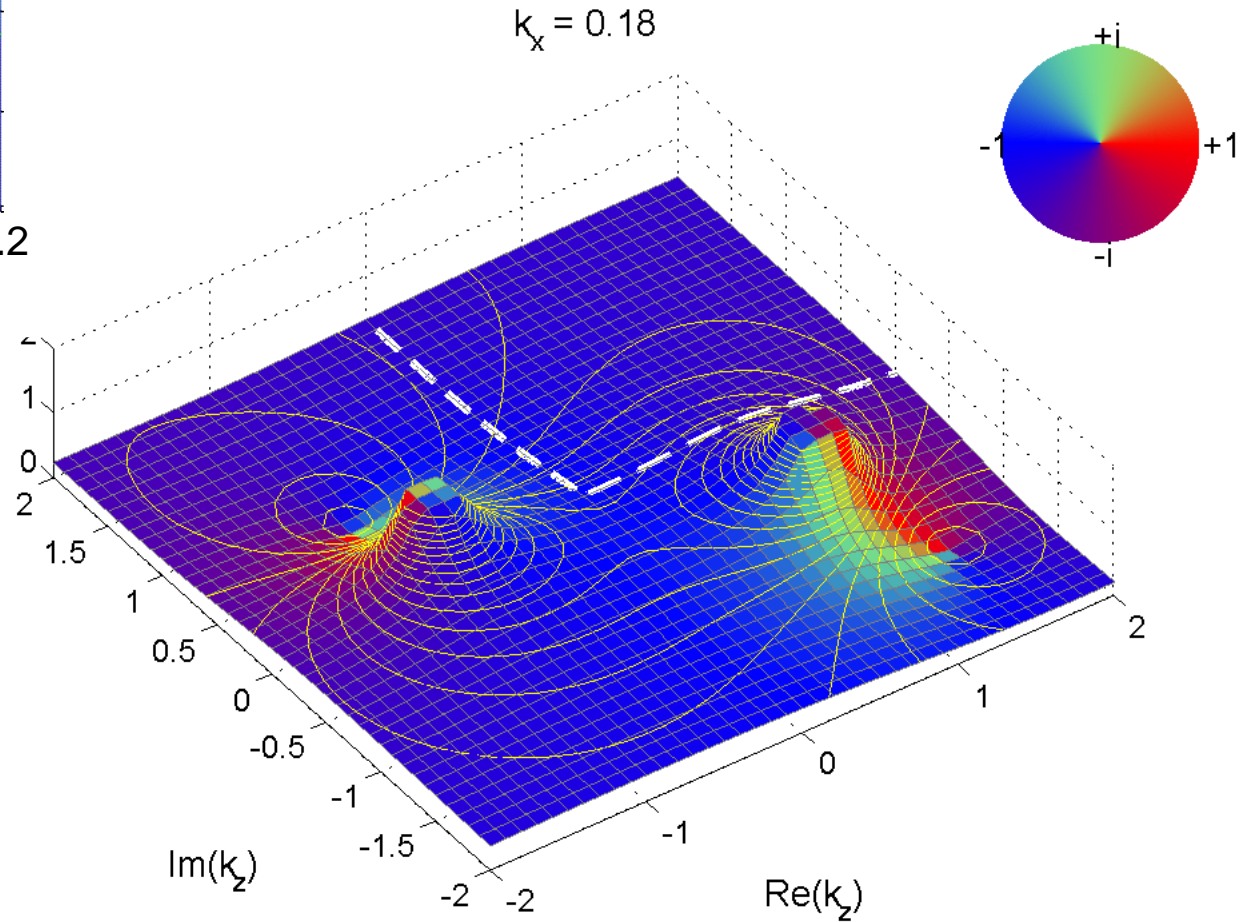


**Resonantly enhanced Wood-Rayleigh anomaly**



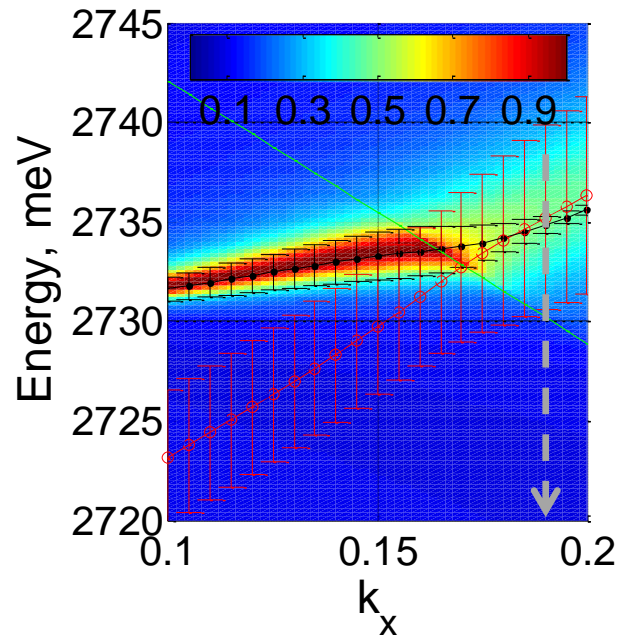


Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction

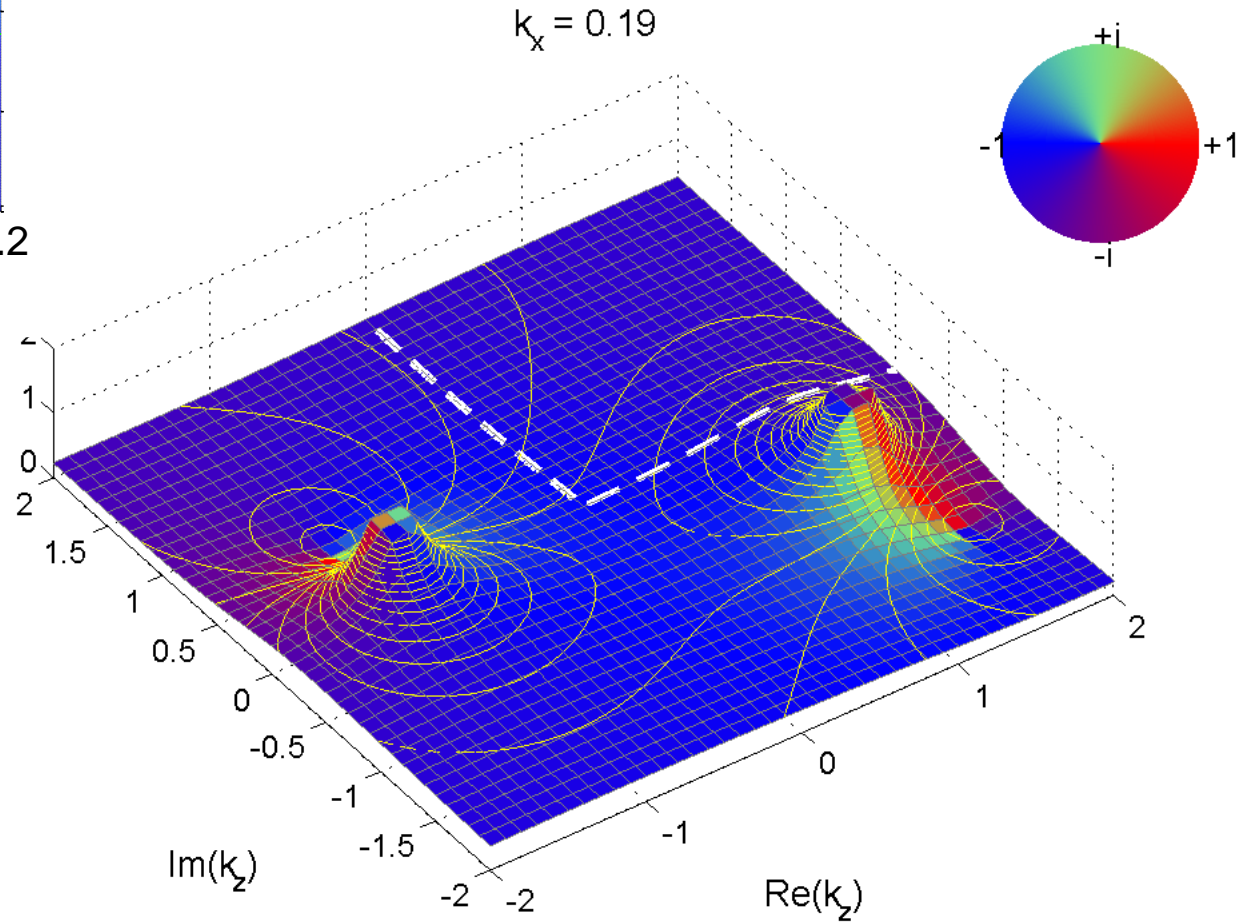


**Resonantly enhanced Wood-Rayleigh anomaly**

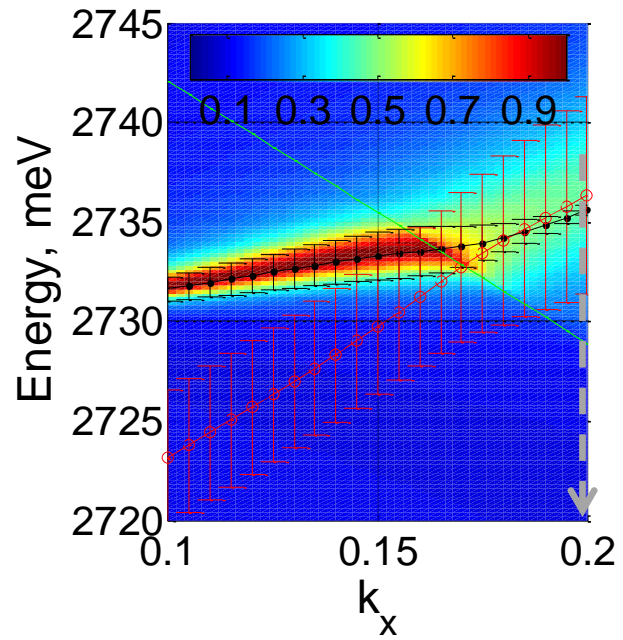




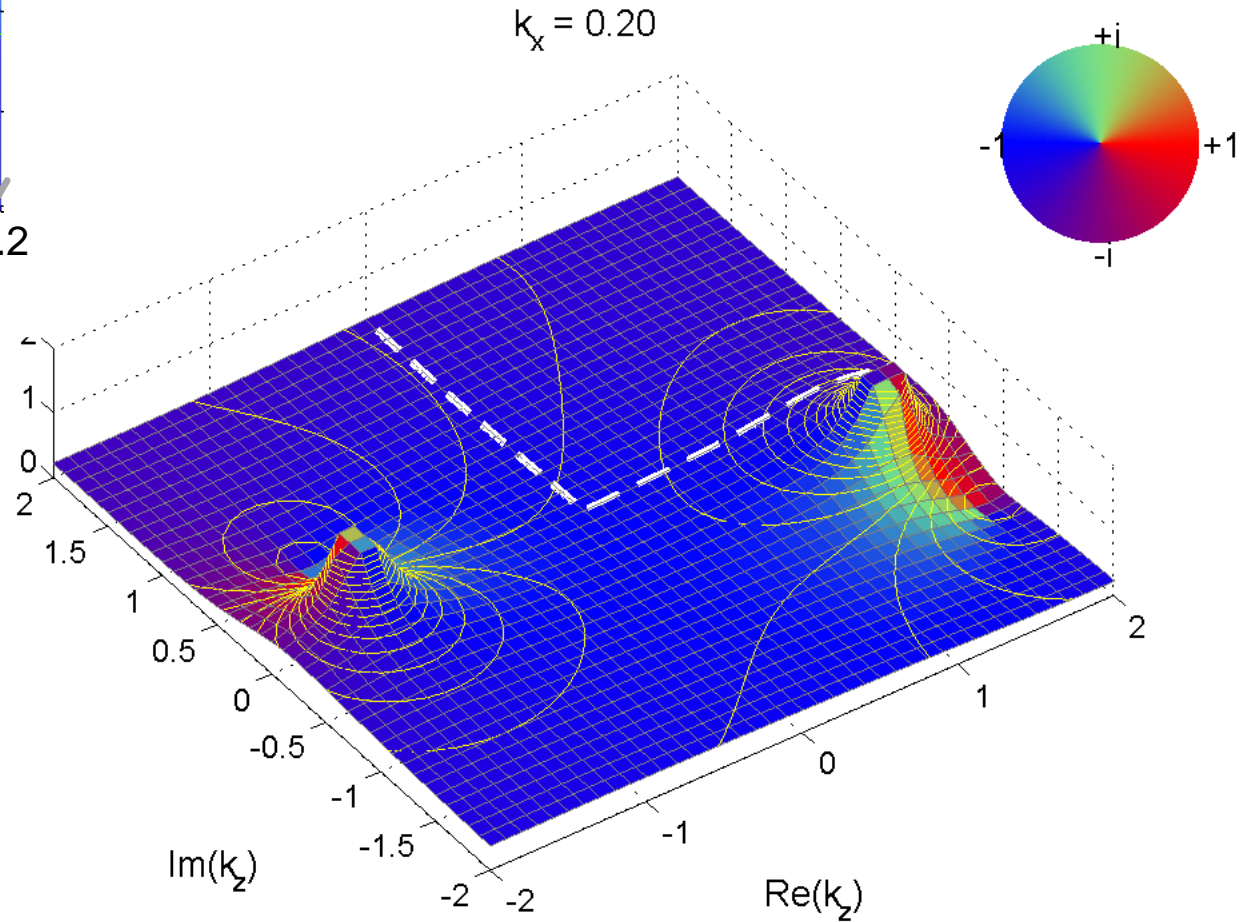
Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction



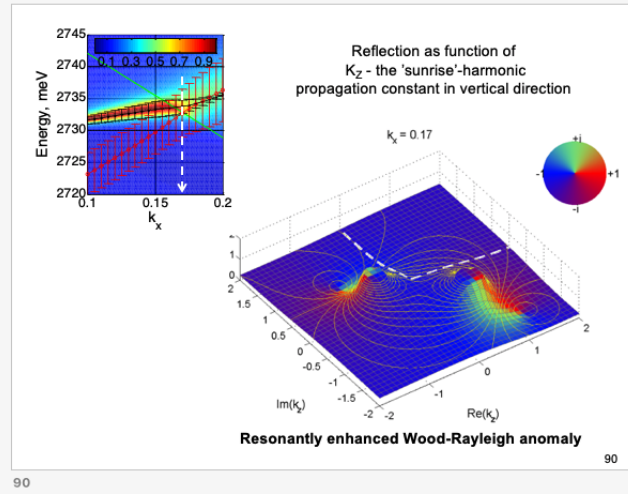
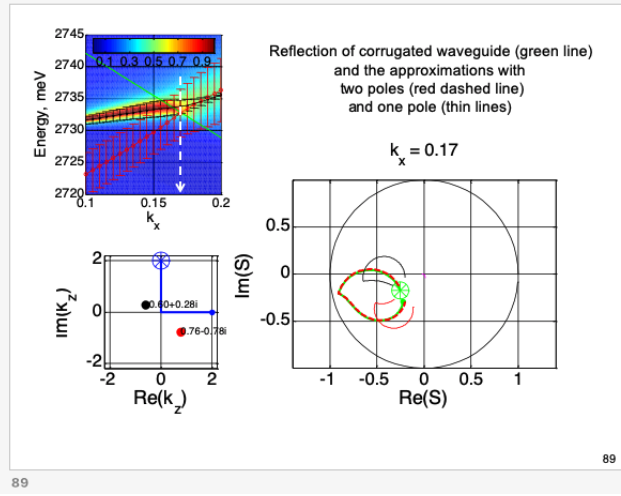
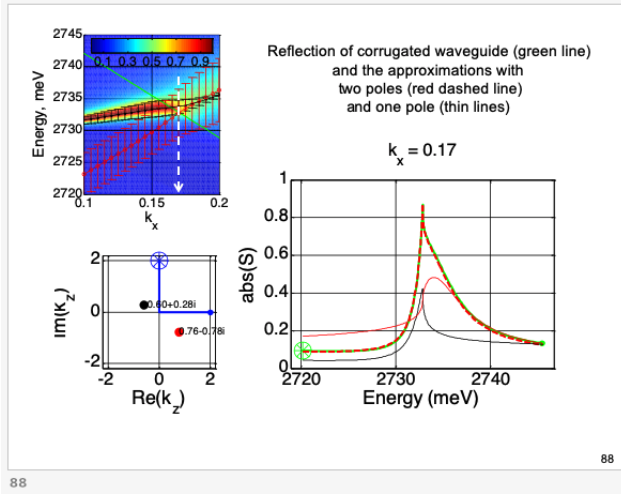
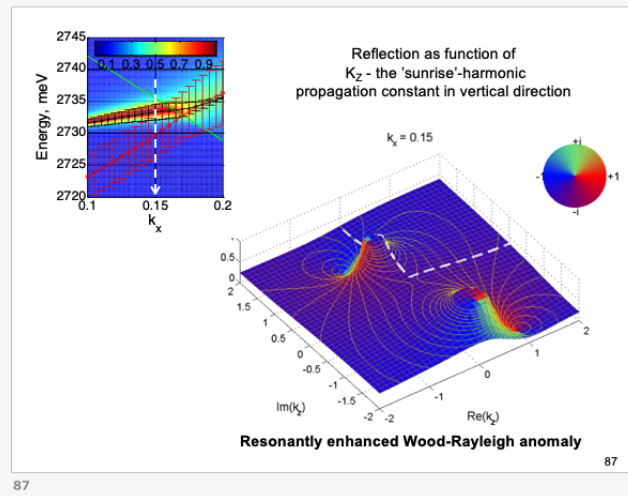
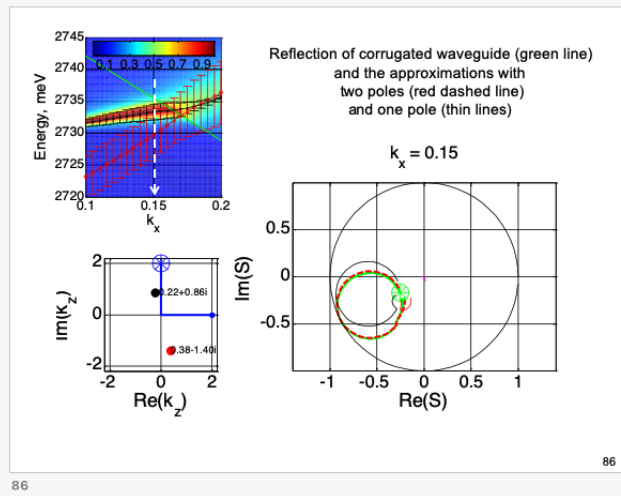
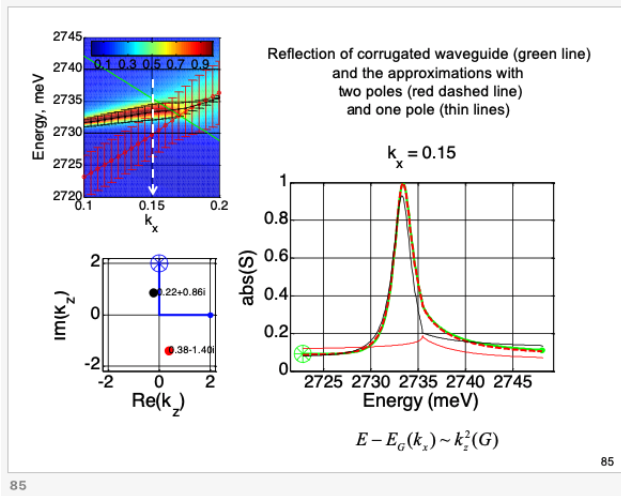
**Resonantly enhanced Wood-Rayleigh anomaly**



Reflection as function of  $k_z$  - the 'sunrise'-harmonic propagation constant in vertical direction

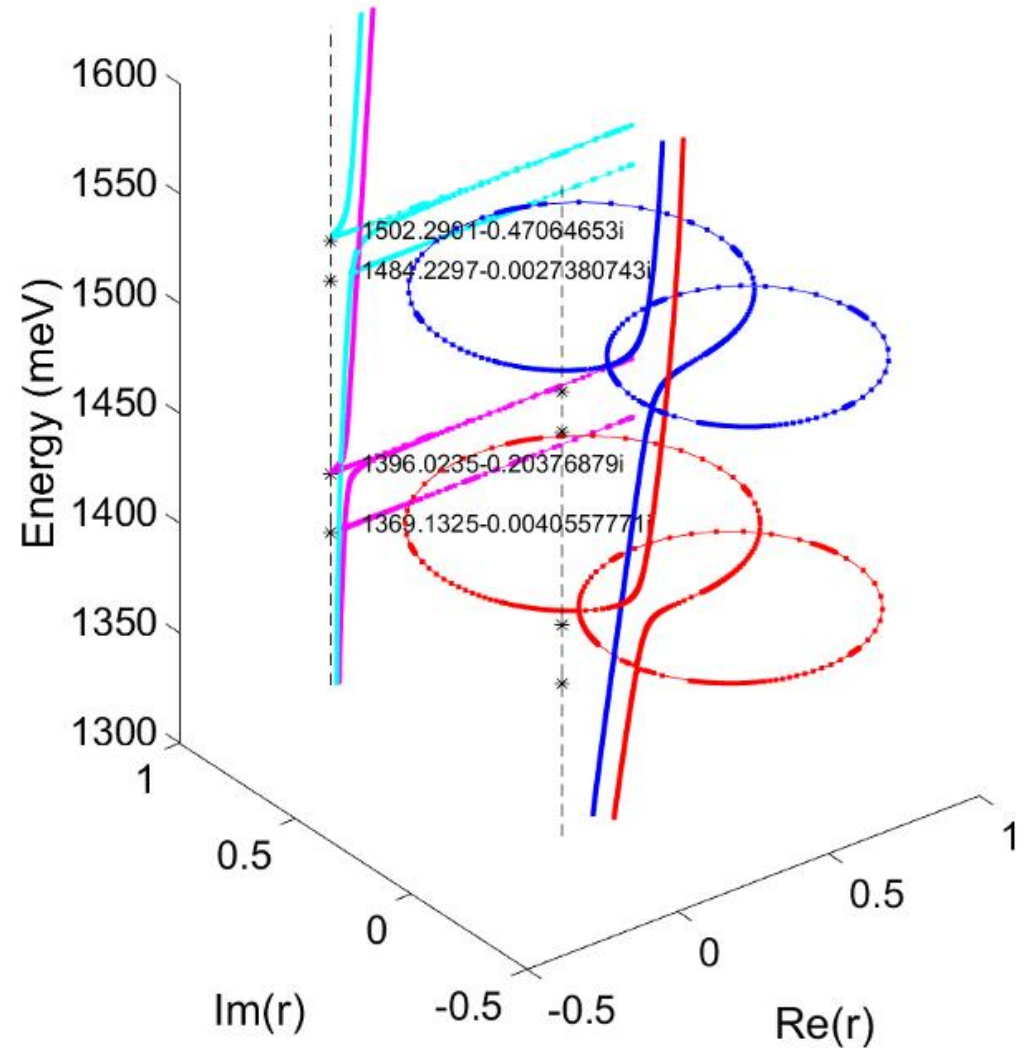
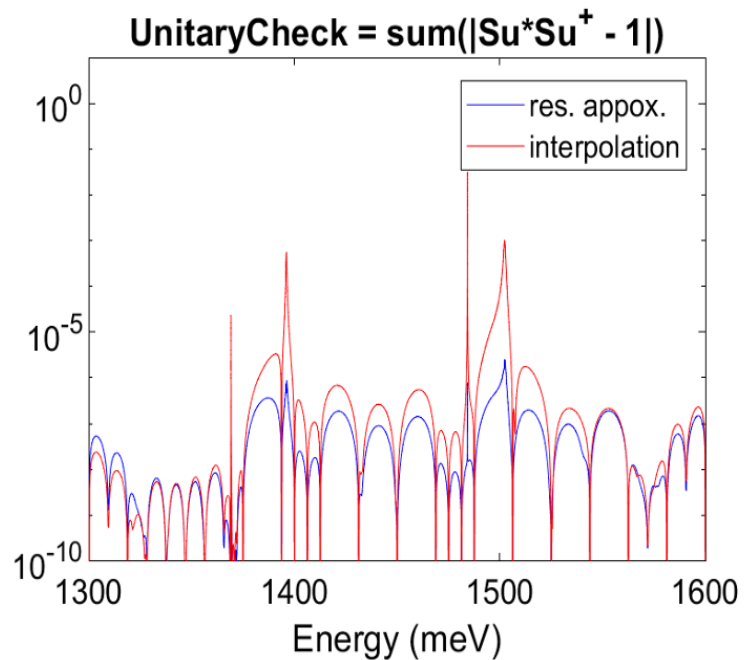


**Resonantly enhanced Wood-Rayleigh anomaly**



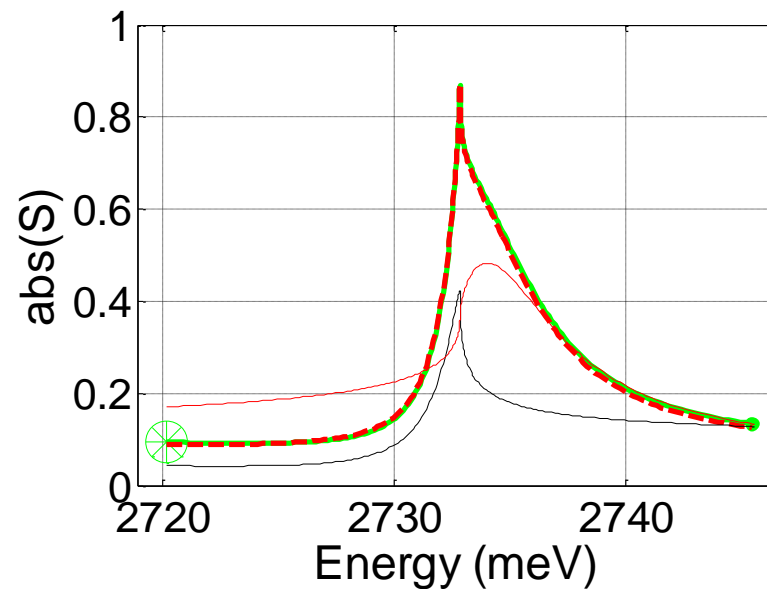
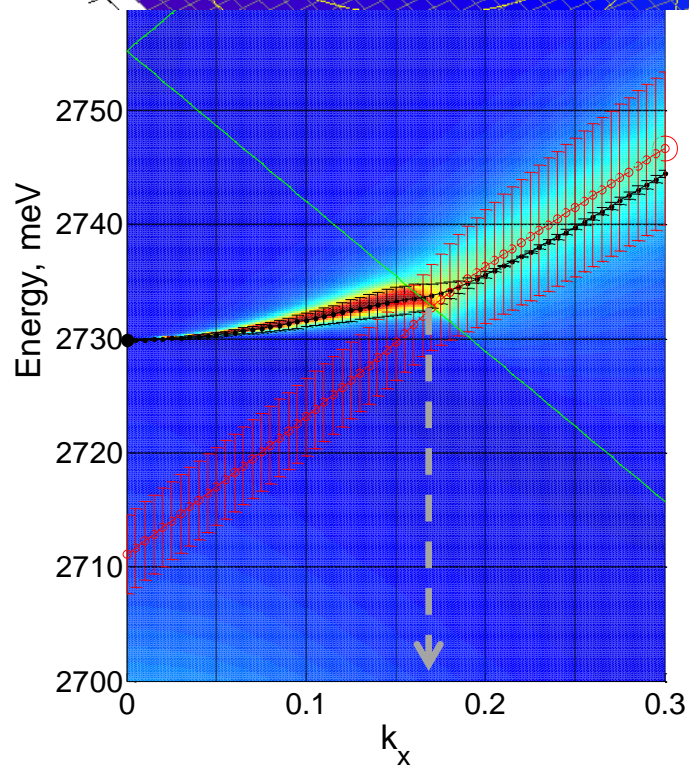
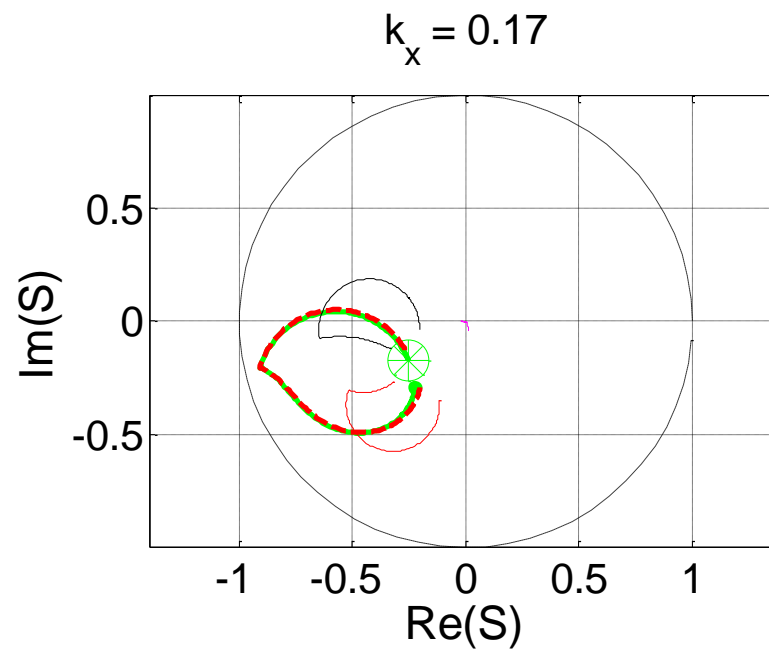
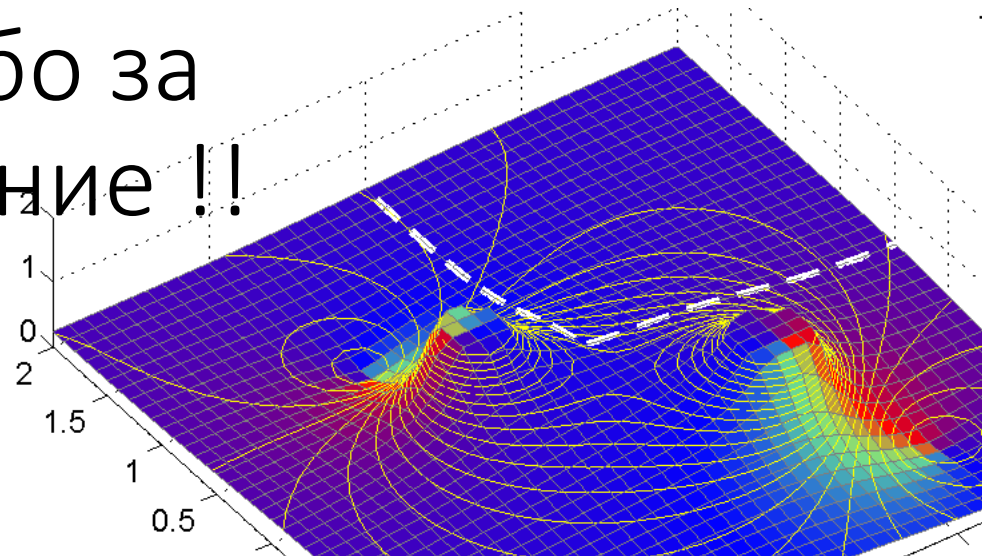
# Выводы -1

интерполяция парциальных матриц  
рассеяния и линейризация резонансной  
фазы позволяет достаточно точно  
описывать резонансы в составных системах



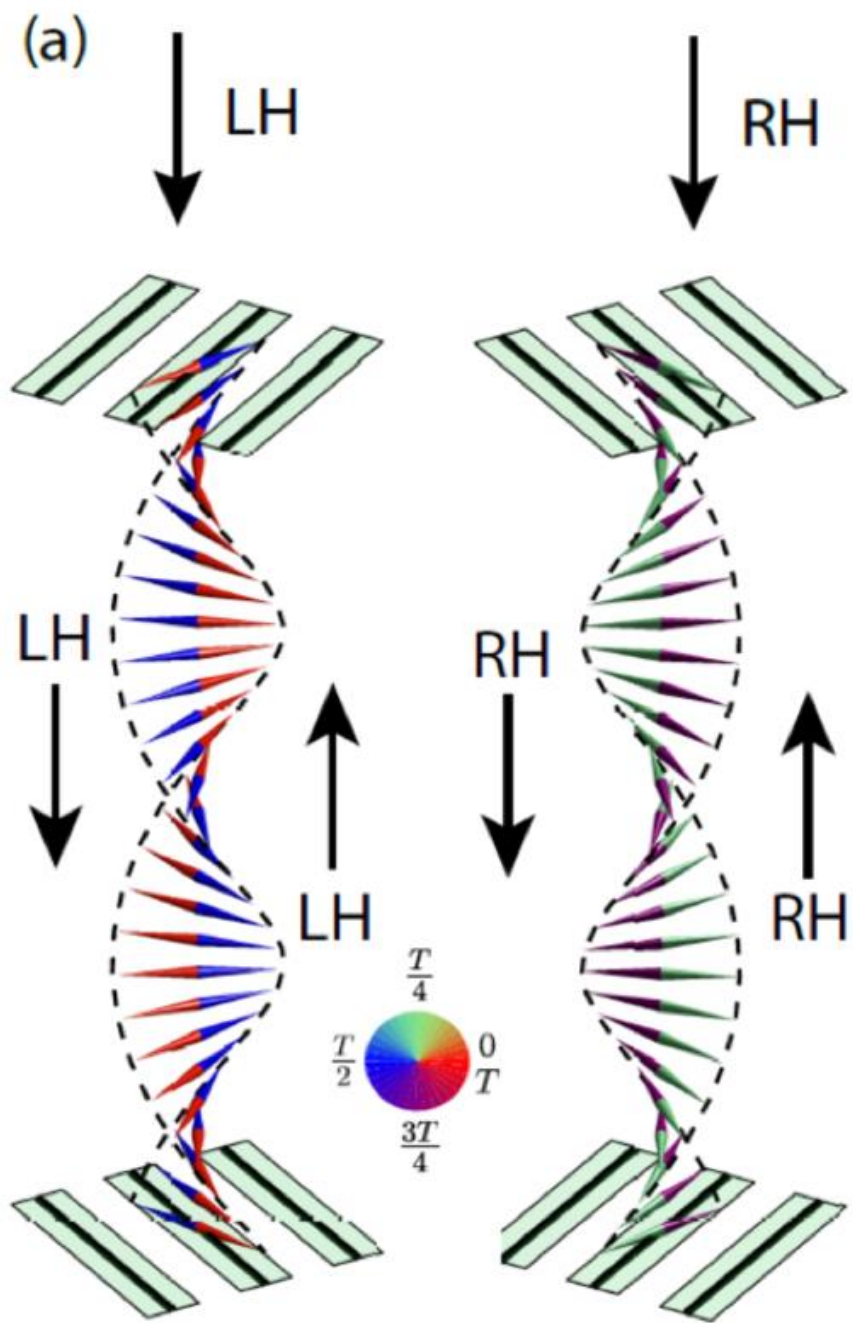


Выводы - 2:  
Спасибо за  
внимание !!





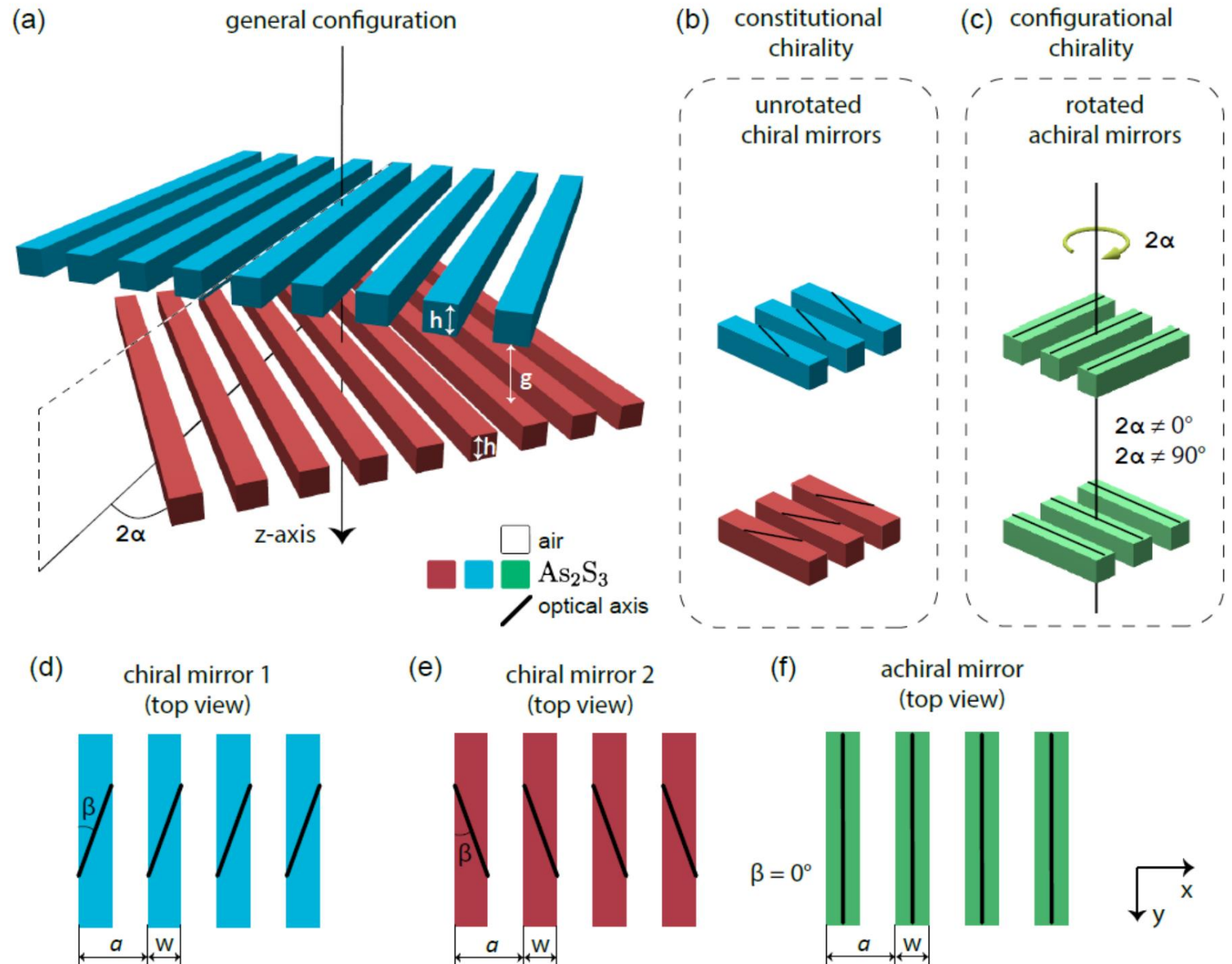
# 4. Как скрутить свет винтом?



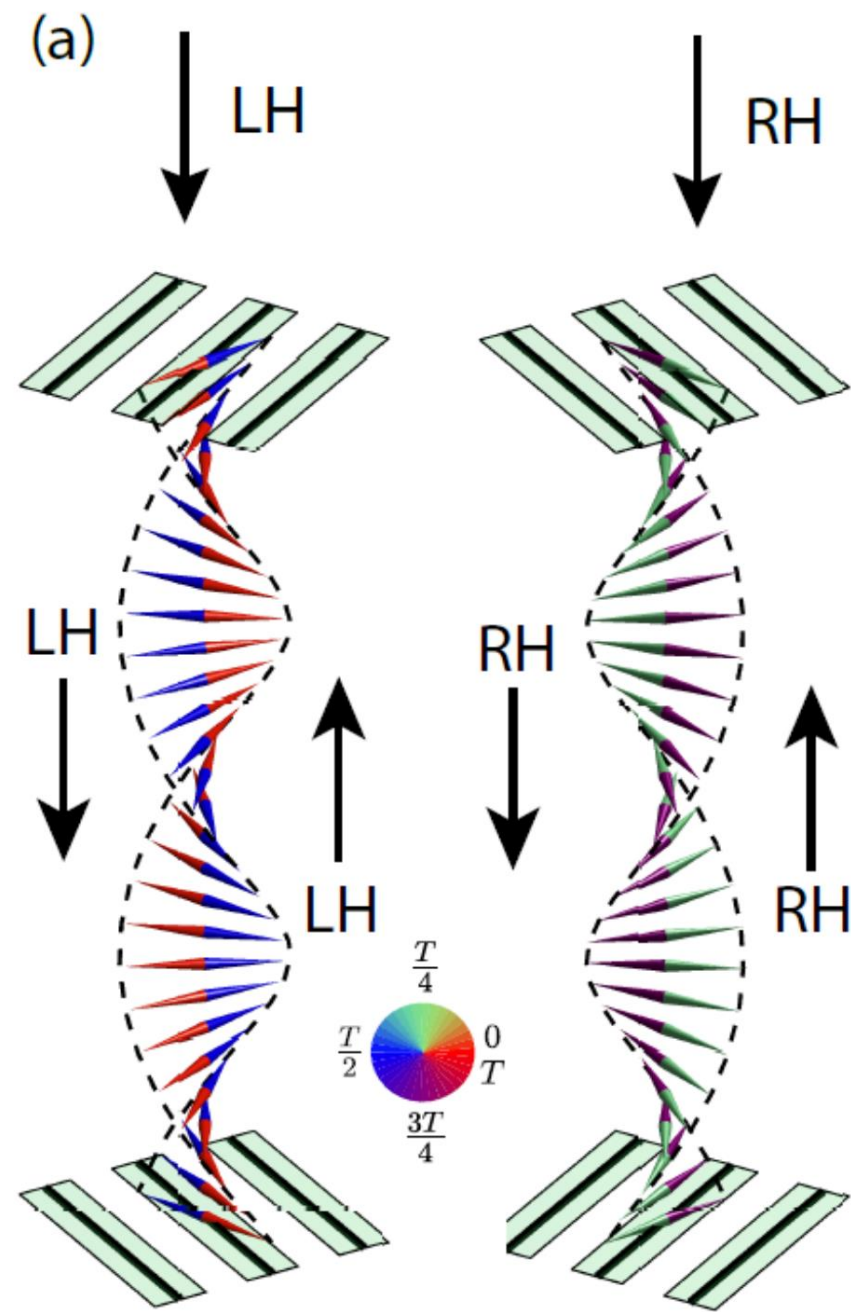
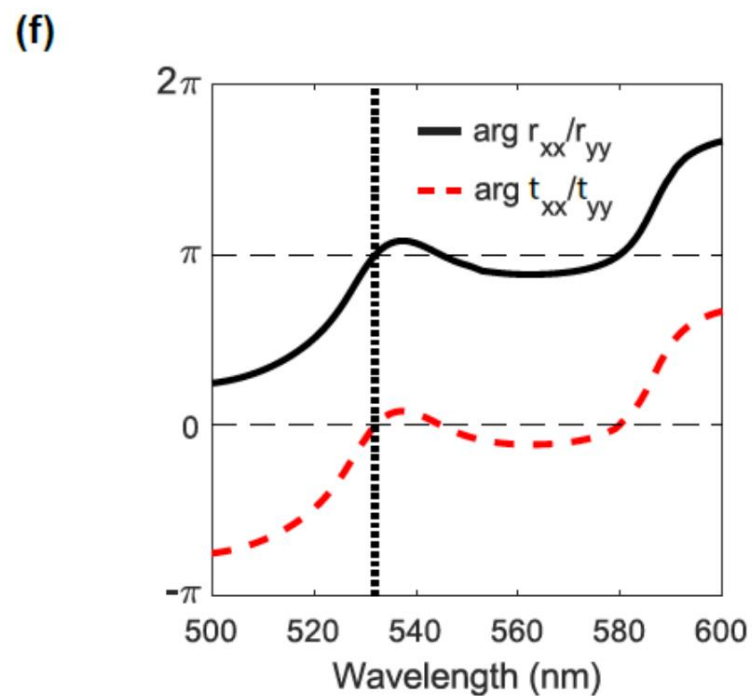
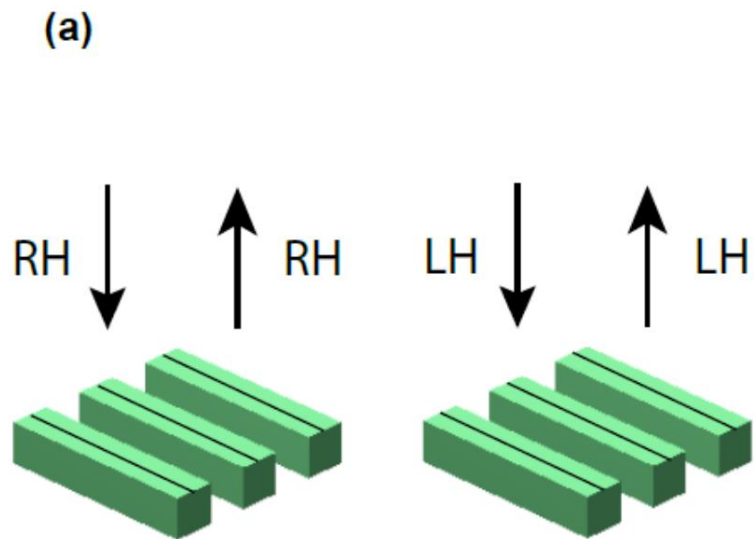
# "Chiral light in twisted Fabry-Pérot cavities"

Sergey A. Dyakov, Natalia S. Salakhova,  
Alexey V. Ignatov, Ilya M. Fradkin,  
Vitaly P. Panov, Jang-Kun Song,  
Nikolay A. Gippius

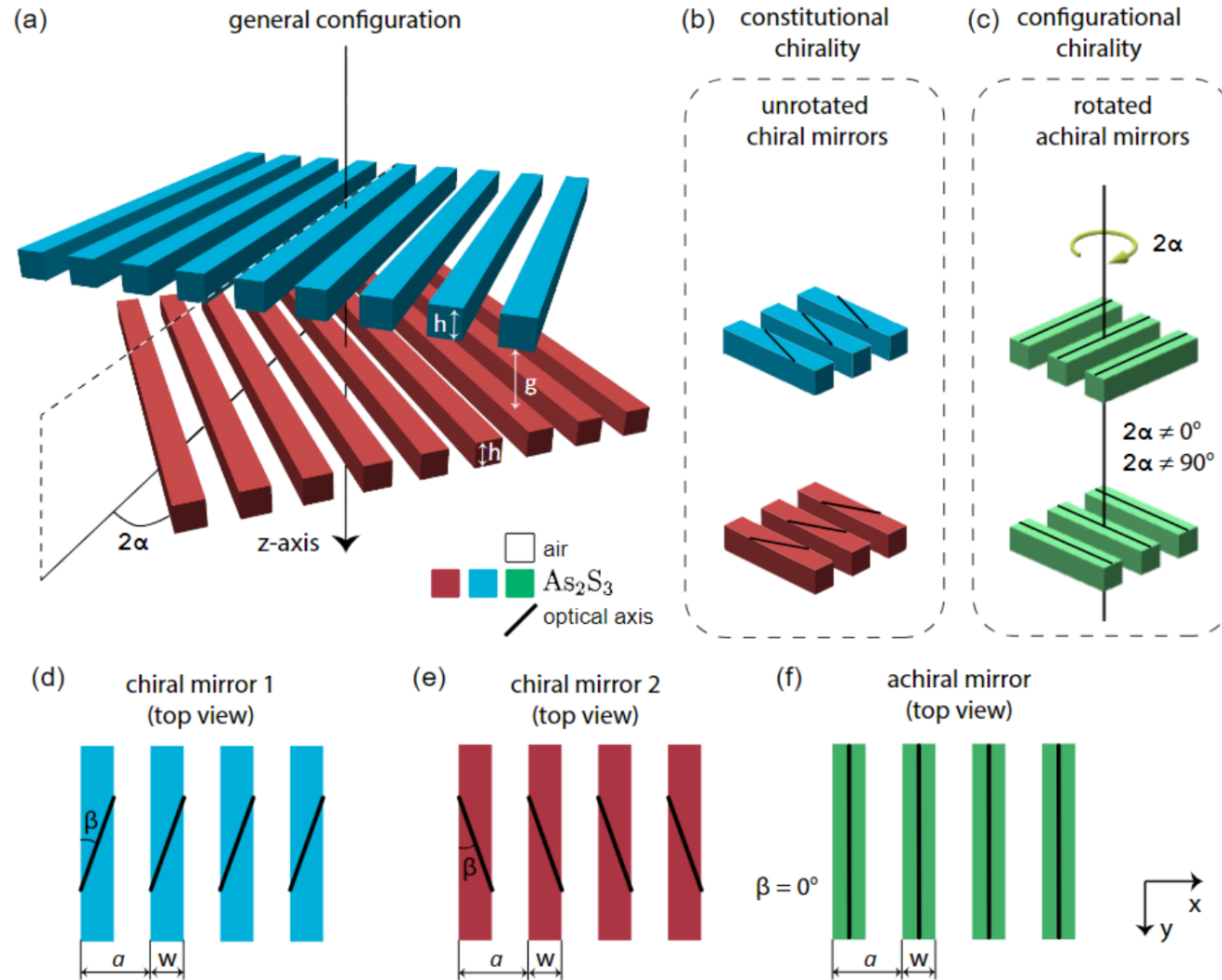
Accepted to  
Advanced Optical Materials.



Cavity with configurational chirality, constructed from non-chiral mirrors twisted into a chiral configuration



# Conclusions



By exploiting the anisotropy of  $\text{As}_2\text{S}_3$ , we have designed cavities with both constitutional and configurational chiralities.

For both types of cavities, we simulated the field distribution of left-handed and right-handed incident waves within the region between the mirrors. At resonant gap sizes, we observed a linearly polarized standing wave with a polarization direction twisted in a helical shape, resulting from the interference between counter-propagating circularly polarized waves of the same handedness.

These chiral Fabry-Pérot cavities can be adjusted to match the technologically available distance between the mirrors by appropriately tuning their twist angle.